

Introduction to Finite Volume Methods - II
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Lecture - 14
Convection term discretization – VI (Private)

So, welcome back to the lecture series of Finite Volume.

(Refer Slide Time: 00:21)

Convection term discretization

T.E.: $\frac{1}{16} \Delta x^3 \phi_c'' - \frac{3}{128} \Delta x^4 \phi_c'''' + \dots = \sim O(\Delta x^3) \Rightarrow$ 3rd order accurate

Stability: $RHS|_c = -\left(\frac{3}{4}\phi_c - \frac{1}{8}\phi_w + \frac{3}{8}\phi_e\right) \|m_c, 0\| + \left(\frac{3}{4}\phi_e - \frac{\phi_{EE} + 3\phi_c}{8}\right) \| -m_c, 0\|$
 $- \left(\frac{3}{4}\phi_c - \frac{\phi_e + 3\phi_w}{8}\right) \|m_w, 0\| + \left(\frac{3}{4}\phi_w - \frac{\phi_{ww} + 3\phi_c}{8}\right) \| -m_w, 0\|$

$\frac{\partial(RHS|_c)}{\partial \phi_c} = -\frac{3}{8} \|m_c, 0\| - \frac{3}{8} \|m_w, 0\| - \frac{3}{8} (m_c + m_w) = < 0$

For uniform vel. \Rightarrow stable
 For Non-uniform \Rightarrow Does not guarantee the boundedness

FROMM scheme: linear profile between the for U & D nodes

And, this quick scheme if you analyze there are truncation error which would turn out to be $\frac{1}{16} \Delta x^3 \phi_c'' - \frac{3}{128} \Delta x^4 \phi_c'''' + \dots$ and so on. So, which essentially give you some sort of an 3rd order accurate, this is 3rd order accurate scheme. So, as we move along the hierarchical direction 2nd order upwind was 2nd order accurate quick becomes 3rd order accurate. Now, the third point one has to check its stability. So, the stability if you look at the right hand side convection term it should be $-\left(\frac{3}{4}\phi_c - \frac{1}{8}\phi_w + \frac{3}{8}\phi_e\right) \|m_c, 0\| + \left(\frac{3}{4}\phi_e - \frac{\phi_{EE} + 3\phi_c}{8}\right) \| -m_c, 0\| - \left(\frac{3}{4}\phi_c - \frac{\phi_e + 3\phi_w}{8}\right) \|m_w, 0\| + \left(\frac{3}{4}\phi_w - \frac{\phi_{ww} + 3\phi_c}{8}\right) \| -m_w, 0\|$ then $-\frac{3}{8} \|m_c, 0\| - \frac{3}{8} \|m_w, 0\| - \frac{3}{8} (m_c + m_w) < 0$.

So, to get the stability limit check you get this right hand side term with respect to ϕ_c which will turn out $-\frac{3}{8} m_c - \frac{3}{8} m_w - \frac{3}{8} (m_c + m_w) < 0$. So, this shows that it is a stability I mean stable scheme. So, if you have some sort of a uniform velocity this will be always stable;

however, this does not guarantee the solution boundedness specially in case of non uniform velocity.

So, one has to for uniform velocity always stable, for non-uniform velocity does not guarantee the boundedness. So, that is the point one has to note that how it goes. Now moving ahead that is another scheme which is called FROMM scheme. So, which is uses the linear profile between the for upwind and downwind nodes connecting the faces.

(Refer Slide Time: 04:35)

Convection term discretization

$\phi(x) = \phi_U + \kappa_1(x - x_C)$
at x_D & $x_U \Rightarrow \phi_D, \phi_U$

$\phi(x) = \phi_U + \frac{\phi_D - \phi_U}{(x_D - x_U)}(x - x_U)$

Uniform, 1D, C-D: $\phi_C = \frac{\phi_D + \phi_U}{2}$

$\phi_f = \phi_U + \frac{\phi_D - \phi_U}{x_D - x_U}(x_f - x_U)$
 $= \phi_C + \frac{x_f - x_U}{x_D - x_U}(\phi_D - \phi_U)$

$\phi_f = \phi_C + \frac{\phi_D - \phi_U}{4}$

$m_e \phi_e = (\phi_C - \frac{\phi_U + \phi_E}{4}) \|m_{e,0}\| - (\phi_E - \frac{\phi_E + \phi_C}{4}) \| -m_{e,0} \|$

$\phi_w m_w = (\phi_C - \frac{\phi_E + \phi_U}{4}) \|m_{w,0}\| - (\phi_U - \frac{\phi_U + \phi_C}{4}) \| -m_{w,0} \|$

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So, how would that look like, if one has to put the stencil again this is our C this is the phase, these are the downstream; this is downstream and this is upstream U U across this cell phases assuming velocity in this direction. So, what it uses it uses some sort of an profile calculation between these points this is phi D phi U and in between you have phi C and phi f. So, this is how it takes into account and the function definition is K naught plus K 1 x minus x C and using the value at x D and x U which is corresponding to phi D and phi U what you get the phi x equals to phi U plus phi D minus phi U x D minus x U x minus x U that is the function that you get.

Now again one can carry out the similar exercise for this particular scheme and what you end up getting is that. Now, if you have this upwind node phi C which you can calculate as phi U plus phi D minus phi U divided by x D minus x U into x C minus x U. So, that is how you get. Now, again the assumption remains uniform grid 1 D convection

diffusion system then phi C becomes phi D plus phi U by 2 which is the centroid value becomes the average between this and this.

Now, with that phi f would be phi U plus phi D minus phi U x D minus x U x f minus x U which one can write phi C plus x f minus x C divided by x D minus x U multiplied with phi D minus phi U. So, you use this information here and you get this and for the uniform grid it reduces to phi f equals to for uniform system it reduces to phi C plus phi D minus phi U divided by 4. So, that is what you get for uniform system.

So, to put this thing information back in the discretized equation you calculate the mass fluxes I mean m dot e fluxes at the phases which will be phi C minus phi W plus phi E by 4 multiplied with m dot e 0 minus phi E phi E E by 4 plus phi C by 4 multiplied with m dot e 0. Similarly you get phi W mw which is phi C minus phi E by 4 phi W by 4 multiplied with m dot w 0 minus phi W phi WW by 4 phi C by 4 minus m w 0. So, this we are doing it for the uniform system.

(Refer Slide Time: 09:29)

Convection term discretization

Discretized eq.:

$$\left\{ \begin{aligned} & (\phi_C - \frac{\phi_W + \phi_E}{2}) \|m_{e,0}\| - (\phi_E - \frac{\phi_{EE} + \phi_C}{2}) \| -m_{e,0} \| \\ & + (\phi_C - \frac{\phi_E + \phi_W}{2}) \|m_{w,0}\| - (\phi_W - \frac{\phi_{WW} + \phi_C}{2}) \| -m_{w,0} \| \end{aligned} \right.$$

$$- \left[\left(\frac{\rho \dot{V}}{\Delta t} \right)_C (\phi_E - \phi_C) - \left(\frac{\rho \dot{V}}{\Delta t} \right)_W (\phi_C - \phi_W) \right] = 0$$

$$a_C \phi_C + a_E \phi_E + a_W \phi_W + a_{EW} \phi_{EW} + a_{EE} \phi_{EE} = 0$$

$$a_E = F_{EW} F_E = -\Gamma_e \frac{S_e}{\Delta x_e} + \frac{1}{4} \|m_{e,0}\| - \| -m_{e,0} \| - \frac{1}{4} \|m_{w,0}\|$$

$$a_W = F_{WW} F_W = -\Gamma_w \frac{S_w}{\Delta x_w} + \frac{1}{4} \|m_{w,0}\| - \| -m_{w,0} \| - \frac{1}{4} \|m_{e,0}\|$$

$$a_{EE} = F_{EE} F_{EE} = \frac{1}{4} \| -m_{e,0} \| ; a_{WW} = F_{WW} F_{WW} = \frac{1}{4} \| -m_{w,0} \|$$

$$a_C = \sum_{f \in \text{faces}(C)} F_{mf} \gamma_f = \Gamma_e \frac{S_e}{\Delta x_e} + \Gamma_w \frac{S_w}{\Delta x_w} + \|m_{e,0}\| - \frac{1}{4} \| -m_{e,0} \| + \|m_{w,0}\| - \frac{1}{4} \| -m_{w,0} \|$$

$$= - (a_E + a_W + a_{EW} + a_{EE}) + (m_{e,0} + m_{w,0})$$

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And then the converted discretized equation would be looking like phi C minus phi W by 4 plus phi E by 4 multiplied with m dot e 0 minus phi E phi EE by 4 plus phi C by 4 multiplied with m dot e 0. And then you get phi C phi E by 4 phi W by 4 multiplied with m dot w 0 minus phi W minus phi WW by 4 plus phi C by 4 multiplied with minus m dot w and then we have the gamma term for the diffusion terms. So, this portion contributes

the convection and then this is the diffusion term by Δx at w which is ϕ_C minus ϕ_W equals to 0.

So, eventually the equation will look like the compact equation will look like a $C \phi_C$ a $E \phi_E$ plus a $W \phi_W$ a $WW \phi_{WW}$ plus a $EE \phi_{EE}$ which is 0. So, this I mean the fact which remains same in the final volume once you get a discretized equation they look exactly similar for different different scheme. So, if you have a one discretization in place one can actually just using different coefficient you have a different scheme in space minus 1 by 4 $m \cdot e$ 0 minus $m \cdot e$ 0 1 by 4 $m \cdot w$ 0.

Then this flux F_w minus $\gamma_w S_w$ by Δx w $m \cdot w$ 0 minus, minus 1 by 4 minus $m \cdot e$ 0. Now a E which will become flux F_e 1 by 4 minus $m \cdot e$ 0 a WW which will be flux f_{small} ww 1 by 4 $m \cdot w$ 0. And finally, a C which is over the phases flux C_f which will get you $\gamma_e S_e$ plus γ_w plus $m \cdot e$ 0 minus 1 by 4 $m \cdot e$ 0 $m \cdot w$ 0 minus 1 by 4 $m \cdot w$ 0 which will again remain as the summation of a E a W a WW EE plus $m \cdot e$ $m \cdot w$.

So, these are the only places the calculation of the coefficients, were you see the differences and if you look at the truncation error for this Fromm scheme.

(Refer Slide Time: 14:11)


Convection term discretization

T.E. = $O(\Delta x^2) \rightarrow$ 2nd order accurate

Stability: $RHS|_{Cw} = -\left(\phi_C - \frac{\phi_W}{4} + \frac{\phi_E}{4}\right) \|m_e, 0\| + \left(\phi_E - \frac{\phi_{EE}}{4} + \frac{\phi_C}{4}\right) \|m_e, 0\|$
 $-\left(\phi_C - \frac{\phi_E}{4} + \frac{\phi_W}{4}\right) \|m_w, 0\| + \left(\phi_W - \frac{\phi_{WW}}{4} + \frac{\phi_C}{4}\right) \|m_w, 0\|$

$\frac{\partial (RHS|_C)}{\partial \phi_C} = -\frac{3}{4} \|m_e, 0\| - \frac{3}{4} \|m_w, 0\| - \frac{1}{4} (m_e + m_w)$

For ϕ_C vcl. = stable
 For ϕ_W vcl. = Does not guarantee stability all the time


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Ashoke De 43

The truncation error would be order of Δx square which also provides you 2nd order accurate system and the stability criteria the stability if you look at the convection term


the right hand side. So, that will have these components ϕ_C minus ϕ_W by 4 plus ϕ_E by 4 $m \cdot e$ plus ϕ_E minus ϕ_{EE} by 4 ϕ_C by 4 multiplied with minus $m \cdot e$ 0 minus ϕ_C ϕ_E by 4 ϕ_W by 4 $m \cdot w$ 0 plus ϕ_W minus. So, if you take the derivative of this term with respect to ϕ_C which will become minus 3 by 4 $m \cdot e$ 0 minus 3 by 4 $m \cdot w$ 0 minus 1 by 4 $m \cdot e$ plus $m \cdot w$.

So, like the quick scheme for constant velocity this guy is always stable and for varying velocity this is I mean it does not guarantee stability all the time it may not. So, but if you have constant on uniform velocity this will always provide you back the stability. So, if you look at the comparison of all this profiles.

(Refer Slide Time: 17:01)

Convection term discretization

<p><u>Comparison</u></p> <p><u>C, D</u></p>	}	<ul style="list-style-type: none"> SOU $(-3/2)$ UD (-1) → false diffusion FROMM $(-3/4)$ QUICK $(-3/8)$ CD (0)
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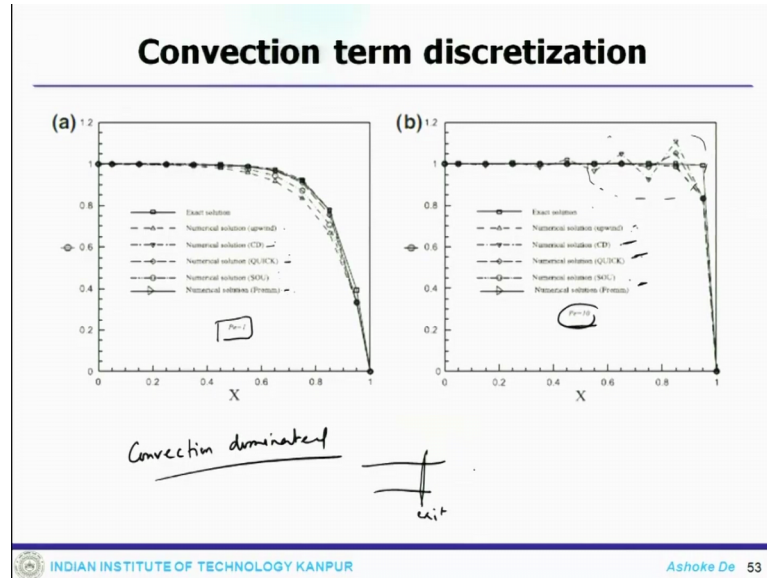

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Ashoke De 44

So, basically some sort of an comparison if you look at various schemes what you see is that with the most negative coefficient is associated with 2nd order upwind scheme that is minus 3 by. And then you have upwind scheme which is minus 1 then FROMM scheme F R O M M FROMM scheme with minus 3 by 4 and then finally, the quick which is minus 3 by 8 and C D is coefficient 0, which is sort of a neutral stability. So, the self corrective action is the sum of both convection and diffusion contribution.

So, the self corrective measures comes from both the convection and diffusion contribution and the false diffusion essentially produced by the upwind scheme which adds to. So, upwind scheme actually the false diffusion or numerical diffusion which

adds to the stability even though the coefficients is minus 1 and it is the most stable scheme among this all that we have discussed.

(Refer Slide Time: 18:33)



So, if you put them together and plot then you see this curve for varying Peclet number this is Peclet number 1 and this is Peclet number 10 and if you plot all the scheme against the exact solution or the analytical solution. So, you get upwind C D quick 2nd order form similarly exact numerical upwind scheme C D scheme quick scheme 2nd order upwind and Fromm scheme.

So, once you look at all the cases for Peclet number 1 which is at the low Peclet number. So, you see that C D Fromm and quick this guy this guy and this guy they I mean some, I mean basically the Peclet number 1 case they show some sort of an regals here I mean the high Peclet number case, but in the low Peclet number case they do not show anything.

But, the schemes at the high Peclet number case show some regals and amongst these is the least accurate is the solution by the 1st order upwind scheme. Now and the solution of this guy is slightly more accurate than 2nd order upwind and; obviously, it is going to be accurate than the 1st order upwind also. Now when you look at this guy this guy and this guys C D quick and Fromm they show some regals in this locations of the solution.

So, here the solution which is obtained by the 1st order upwind and the 2nd order upwind their profile being the most accurate close to the exact solution. And now coming back to the Fromm and quick scheme on the other hand they show some ripples, but with some small amplitude and the reason for this ripples is that they imposed boundary recondition at the exit of the domain. So, since the problem is convection dominated it is affected by the. So, these are for convection dominated flow the solutions at the downstream is partially affected by the upstream calculations or upstream information.

So, solution which are actually at the exit of the domain which are specified this is the exit of the domain that have some impact on this ripples and all these things, but. However, the 2nd order upwind scheme is expected to give rise to some oscillation in the presence of high gradient in the domain if you have a shock wave or like that. So, that give some idea on some sort of an comparison of this.

(Refer Slide Time: 22:03)

Convection term discretization

Functional Relationship for Uniform & Non-Uniform grids

Curvilinear co-ordinate

Uniform grid

Scheme	Uniform	Non-Uniform
upwind	$\phi_f = \phi_c$	$\phi_f = \phi_c$
downwind	$\phi_f = \phi_D$	$\phi_f = \phi_D$
CD	$\phi_f = 0.5(\phi_c + \phi_D)$	$\phi_f = \phi_c + \frac{(\phi_D - \phi_c)(\xi_f - \xi_c)}{(\xi_D - \xi_c)}$
SOU	$\phi_f = \frac{3}{4}\phi_c - \phi_U/2$	$\phi_f = \phi_c + \frac{(\phi_c - \phi_U)(\xi_f - \xi_c)}{(\xi_c - \xi_U)}$
QUICK	$\phi_f = \frac{3}{4}\phi_c - \phi_U/8 + \phi_D/8$	$\phi_f = \phi_U + \text{term I} + \text{term II}$
FROMM	$\phi_f = \phi_c + \frac{\xi}{4}(\phi_D - \phi_U)$	$\phi_f = \phi_c + \left(\frac{\xi_f - \xi_c}{\xi_D - \xi_U}\right)(\phi_D - \phi_U)$

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Now, moving ahead you can actually find some functional relationship for uniform and non uniform grids. So, far our discussion was restricted to uniform grid. So now, one can find out some correlation or the mapping between uniform and non uniform grid and how you can do that essentially let us see if you have some sort of an element C here which will be connected with a phase this is the phase f and the velocity here will be in this direction then there is a cell D then there would be cell D D downstream the

upstream there is U upstream U . So, this goes in a ξ direction. So, sort of a curvy linear coordinate system.

So, typically my x y system or the uniform x y system the cells would be sitting like that this is the cell down of that up of that. So, this is where your uniform grid system or orthogonal grid system and this is in curvilinear grid system. So, one important point which is here the functional relationship for uniform grid remain exactly same and for the non uniform also independent of whether the Cartesian or the curvilinear grid system is used. For non uniform grid the independent variable x which is appearing in the functional form should be replaced by ξ .

So, what essentially it says the functional form would not get change it only the notation like uniform grid used get a functional form of x ; in this case it should be function of ξ which will actually represent the system with the discussion. Now, if you get all this functional relationship for different scheme then it should be put together let us say this is for uniform system this is non uniform system, this is non uniform system then these are the schemes that we have so far have discussed.

So, when you talk about the upwind scheme this will be uniform case ϕ_f is ϕ_C this case also ϕ_f is ϕ_C when you talk about downwind scheme then ϕ_f is ϕ_D this case ϕ_f is also ϕ_D . When you talk about C D this is ϕ_f is 0.5 into ϕ_C plus ϕ_D and this case ϕ_f equals to ϕ_C plus ϕ_D minus ϕ_C by ξ_D minus ξ_C multiplied with ξ_f minus ξ_C . When you go to 2nd order upwind get ϕ_f equals 3 by ϕ_C minus ϕ_U by this case it would be ϕ_f equals to ϕ_C plus ϕ_C minus ϕ_U divided by ξ_C minus ξ_U multiplied with ξ_f minus ξ_C .

Now, quick you get ϕ_f equals to 3 by 4 ϕ_C minus ϕ_U divided by 8 plus ϕ_D into 3 by 8 and this case this is ϕ_U plus you have term I plus term II we will see how they look like. And, if you put Fromm scheme this would be should be ϕ_f equals to ϕ_C plus 1 by 4 ϕ_D minus ϕ_U and this case it is ϕ_C plus ξ_f minus ξ_C divided by ξ_D minus ξ_U multiplied with ϕ_D minus ϕ_U . So, you can see this difference in the functional form of this expression.

(Refer Slide Time: 28:07)

Convection term discretization

$\phi_{U \text{ face}}$: Non-uniform: term I = $\left(\frac{\xi_f - \xi_u}{\xi_D - \xi_u} \right) \left(\frac{\xi_f - \xi_c}{\xi_D - \xi_c} \right) (\phi_D - \phi_U)$
 term II: $\left(\frac{\xi_f - \xi_u}{\xi_c - \xi_u} \right) \left(\frac{\xi_f - \xi_D}{\xi_c - \xi_D} \right) (\phi_c - \phi_U)$

$\xi_u = \xi_{uu} + (\xi_u - \xi_{uu})$
 $\xi_{uu} = \sqrt{(\alpha_{uu} - \alpha_0)^2 + (\beta_{uu} - \beta_0)^2 + (\gamma_{uu} - \gamma_0)^2}$
 $\xi_u - \xi_{uu} = \sqrt{(\alpha_u - \alpha_{uu})^2 + (\beta_u - \beta_{uu})^2 + (\gamma_u - \gamma_{uu})^2}$
 $\xi_1 - \xi_2 = \sqrt{(\alpha_1 - \alpha_2)^2 + (\beta_1 - \beta_2)^2 + (\gamma_1 - \gamma_2)^2}$

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Now, once for the quick the non-uniform system the term I is $\xi_f - \xi_u$ divided by $\xi_D - \xi_u$ multiplied with $\xi_f - \xi_c$ divided by $\xi_D - \xi_c$ with $\phi_D - \phi_U$. And, term II they would look like $\xi_f - \xi_u$ divided by $\xi_c - \xi_u$ multiplied with $\xi_f - \xi_D$ divided by $\xi_c - \xi_D$ with $\phi_c - \phi_U$.

So, where ξ_u is $\xi_{uu} + (\xi_u - \xi_{uu})$ and ξ_{uu} is nothing, but $\sqrt{(\alpha_{uu} - \alpha_0)^2 + (\beta_{uu} - \beta_0)^2 + (\gamma_{uu} - \gamma_0)^2}$. So, in this particular this somewhere centre is sitting. So, plus $\sqrt{(\alpha_u - \alpha_{uu})^2 + (\beta_u - \beta_{uu})^2 + (\gamma_u - \gamma_{uu})^2}$. Now so, if you put them together $\xi_1 - \xi_2$ is essentially $\sqrt{(\alpha_1 - \alpha_2)^2 + (\beta_1 - \beta_2)^2 + (\gamma_1 - \gamma_2)^2}$.

So, this is how you can actually get the functionality between uniform and non-uniform system and the functional form look similar only thing the calculation of the non-uniform system the coordinate system becomes ξ system. So, we will stop here and will continue the discussion in the next lecture.

Thank you.