

Satellite Attitude Dynamics and Control
Prof. Manoranjan Sinha
Department of Aerospace Engineering
Indian Institute of Technology, Kharagpur

Lecture – 36
Gravity Gradient Satellite (Contd.)

Welcome to the 36th lecture, today we are going to discuss about the Gravity Gradient Satellite what we have been continuing with what we will concentrate on the elliptic orbit. So, if the gravity gradient satellite or satellite which can be gravity gradient has stabilized, if we are working on its motion equation so, how the things we will develop that we are going to look into today.

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Lecture-36
Gravity-gradient Satellite

Pitching motion of gravity-gradient satellite in elliptic orbit:
Assumption: no roll and yaw motion ($\theta_1=0, \theta_3=0$) derivative terms also zero.

$$I_2 \dot{\omega}_2 - \omega_3 \omega_1 (I_3 - I_1) = -3\omega_0^2 (I_3 - I_1) c_{33} c_{13}$$

$$v_{c_{33} c_{13}} = \omega_0^2 \sin^2 \theta = \omega_0^2 \sin^2 \theta$$

$\omega_2 = \dot{\theta}_2 - \omega_0$ [when ω_0 is not a constant]
↑
 Small angular motion / large motion (argument)

$\omega_0 = \text{a constant}$
 i.e. in a circular orbit.
 $\dot{\omega}_0 = 0$ for circular orbit - but not for the elliptical orbit

In fact $\omega_0 \rightarrow$ for elliptical orbit $\rightarrow \omega_0 = \frac{h}{r^3 a}$
 $\omega_0 \rightarrow$ is a constant in a circular orbit because r is a constant

So, we assume here no roll and yaw motion. So, already we have observed that for the circular orbit roll and yaw, this decoupled from the pitching motion. So, let us write that if there is no roll and yaw motion means the theta 1 this will be 0 and theta 3 will be 0. And therefore, the related derivative terms will also be 0 derivative terms also 0.

So, just we have concerned with the pitch motion dynamics and we know that for the motion about the second axis that we can write as $I_2 \dot{\omega}_2 - \omega_3 \omega_1 (I_3 - I_1) = -3\omega_0^2 (I_3 - I_1) c_{33} c_{13}$, ok. This is the way we

have described, but here this we have written for the case where ω_0 , this is a constant means it is a moving in a circular orbit.

Now, already we have observed that C_{33} into C_{13} , this quantity is coming from the third column of the direction cosine matrix. So, we need to put or insert this term. So, this term basically we have approximated as $\cos \theta$ times $\sin \theta$. So, I will put here the θ^2 and θ^2 term. Otherwise, we will have to go back into the old set of lectures and from there find out all this quantities which were involved.

And ω_2 , we have written as $\dot{\theta} - \omega_0$ and this we have done for small angular motion. So, from this place ω_2 dot, this will be $\ddot{\theta} - \dot{\omega}_0$. Then ω_0 is not a constant. So, in that case this will not vanish. For the circular orbit $\dot{\omega}_0$ equal to 0 for circular orbit, but not for the elliptical orbit.

In fact, ω_0 this for elliptical orbit, we have to write it properly. See, here in this case ω_0 for the circular orbit. It has got its notation from μ_0 by r^3 . In the circular orbit, because or this was simply we have written μ , ok. So, in the circular orbit r is constant and therefore, ω_0 turned out to be a constant. So, ω_0 is a constant in a circular orbit, because r is a constant. However, in the elliptical orbit this r^3 is varying and therefore, the things will change, ok. So, we are today going to discuss about all this things.

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So, this θ^2 , θ^2 what we are writing here if we replace this in terms of $\cos \theta$ and $\sin \theta$, it will be much convenient to work with, because this θ^2 will be able to get rid off, ok.

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$I_2 \omega_2$ absolute angular velocity
 $\theta_1 = \theta_3 = 0$ $\theta_2 = \theta$
 $c_{23} = \sin \theta_1 \cos \theta_2 = 0$
 $c_{33} = \cos \theta_1 \cos \theta_2 = \cos \theta$
 $c_{13} = -\sin \theta_2 = -\sin \theta$
 angular velocity of the satellite $\rightarrow \dot{\theta}$
 $\omega_2 = \dot{\omega} = \dot{\theta} - \dot{\phi} \Rightarrow \dot{\omega} = \ddot{\theta} - \ddot{\phi} = \ddot{\theta} - \ddot{\omega}_0$
 $\dot{\phi} \rightarrow$ in central force motion (inverse square gravitational field)
 $r^2 \dot{\phi} = h = \text{angular momentum/mass}$
 $\dot{\phi} = \frac{h}{r^2}$ $\theta_1, \theta_2, \theta_3 \rightarrow$ measured wrt orbital frame.
 $\dot{\phi} \rightarrow$ negative e_2
 $\dot{\omega} =$ angular velocity direction of the orbital axis + angular velocity of the satellite wrt. the orbital axis

So, to start with we have C_{23} , we have written as $\sin \theta_1 \cos \theta_2$, and because the θ_1 , here we are choosing to be θ_1 and θ_3 to be 0. So, this quantity vanishes C_{33} equal to $C_{\theta_1} \cos \theta_2$. So, this we are writing as $\cos \theta_2$, because θ_1 this quantity is 0. So, therefore, C_{θ_1} will become 1. So, better we should write here this quantity to be equal to 1, this is equal to 1 and θ_2 , we are writing as θ . So, this we can write as $\cos \theta$ and C_{13} equal to $-\sin \theta_2$ minus $\sin \theta$.

So, we utilise this information and put it in our equation of motion. Then therefore, I_2 times pick up this now what will be the value of $\ddot{\omega}$; if there is no other motion involved. So, only this motion is involved. So, even for this particular part this is as well valid for the large motion, large angular motion also large motion angular, because if here is the earth is there and this angle let us write this as the ϕ , ok.

And, then we have taken this direction to be the e_3 direction orbital and suppose the satellite is going in the elliptical orbit. So, the velocity vector will be along this direction, but in our case e_1 will be along this and e_2 will be directly going inside the paper.

So, you can see the positive direction of the rotation of positive direction of rotation about the e_2 axis about this particular axis it is in this sense, ok. So, it is opposite to the ϕ direction, right where this is I am drawing the figure here separately, this is e_3

$e_o 1$ and this is $e_o 2$ which is going into the paper here in this place, in this place it is a going into the paper.

So, now we will look here in this motion, ok. So, this motion is in the, this particular ϕ . So, the changes which are taking place for this ϕ which we will write as the ϕ dot, ok. So, this is along the negative $e_o 2$ direction, ok. And therefore, because no other motion is involved so, you can check it directly, we can write here the angular velocity of the satellite to be motion with respect to the.

So, ω so, this will be the angular velocity of the orbital axis plus angular velocity of the satellite with respect to the orbital axis. So, angular velocity of the satellite, we are indicating by θ satellite. So, this is pitching motion basically satellite is indicated by θ . So, therefore, you can see that θ and this for angular velocity of the satellite is indicated by θ dot. So, ϕ dot and θ dot they are opposite to each other, ok.

So, the ω then becomes ω can be written like ω equal to this is only in one direction. And therefore, we can write this as θ dot minus ϕ dot, ok. And therefore, the ω dot it can be written as θ double dot minus ϕ double dot which in the on the earlier page we have written this as ω_0 dot,.

Now, this ϕ dot, you will have to be go back and look into my lecture on space flight mechanics and how this ϕ is derived. So, basically in a, in central force motion of inverse square field gravitational, we can write here $r^2 \phi$ dot this equal to h . And this is called the angular momentum per unit mass; this is angular momentum per unit mass ok. So, ϕ dot at any instant it is a related to related by this expression. So, h is your angular momentum per unit mass.

So, the set of equation which is required to indicate this ω_0 dot here, is nothing, but your this quantity ω_0 dot is equal to ϕ double dot. So, it will directly appear from this place, ok. So, you can look back into that lecture it is perhaps in the beginnings from the 6th, 7th lecture you will find it, it is a part of the central force motion,.

So, what I wanted to tell here that this ω the angular velocity of the satellite, it will be indicated like this. So, here in this case, this is your ω_2 equal to ω , what we are writing here ω_2 . Why we are writing ω_2 ? This is I_2 times ω_2 is the

here the absolute angular velocity absolute, angular velocity, this is the absolute angular velocity.

So, this notation should be little bit clear, see here omega 2, we have written omega 3 we have written. So, all these are absolute angular velocity and this then we have converted in terms of theta 2. So, theta 1 theta 2 theta 3 ah, let us write here theta 1 theta 2 theta 3, this we are measured with respect to the orbital frame, ok. And therefore, it appears in this format so, this is measured with the respect to the orbital frame. And this is the motion of the orbital frame itself, ok. Now we are ready to face the problem, ok.

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$$I_2(\ddot{\theta} - \dot{\phi}) - \omega_3 \omega_1 (I_3 - I_1) = -\frac{3M}{r^3} (I_3 - I_1) C_{33} C_{13} = +\frac{3M}{r^3} (I_3 - I_1) \sin\theta \cos\theta$$

$$C_{33} C_{13} = \cos\theta (-\sin\theta) = -\sin\theta \cos\theta$$

$$\ddot{\theta} - \dot{\phi} - \frac{\omega_3 \omega_1}{I_2} (I_3 - I_1) = -\frac{3M}{r^3} \frac{(I_3 - I_1) \sin\theta \cos\theta}{I_2}$$
 for the case when $\omega_3 = \omega_1 = 0$ [no yawing/rolling motion involved]

$$\ddot{\theta} - \dot{\phi} - \frac{3M}{r^3} \frac{(I_3 - I_1) \sin\theta \cos\theta}{I_2} = 0$$
 ?pitching motion equation in elliptic orbit
 no approximation assumed here

ω_2 | ω_1
 $\omega_3 \rightarrow$ component of the absolute angular velocity along the third body axis

So, we have I 2 times omega 2 here in this case we will pick up from this place theta dot minus phi dot. So, this becomes theta double dot minus phi double dot omega 3 times omega 1 I 3 minus I 1 this equal to, on the right hand side we have minus 3 mu by r cube times I 3 minus I 1, C 3 1 and C 3 3 and C 1 3, this 3 is that for this 3 appears here in this place. For this one this one is appearing here then this is the 3 3 appearing here. For the case we are ah, we are considering the orbital third axis toward the centre of the earth,.

And this thing, so, already we have discussed in detail. So, there is no need of further going into this ok. So, C 3 3 all this things we have to write here C 3 3 times C 1 3 then becomes C 3 3 is here cos theta and C 1 3 is minus sin theta. So, therefore, this is minus

sin theta into cos theta. So, here we can write plus 3 mu by r cube times I 3 minus I 1 sin theta into cos theta.

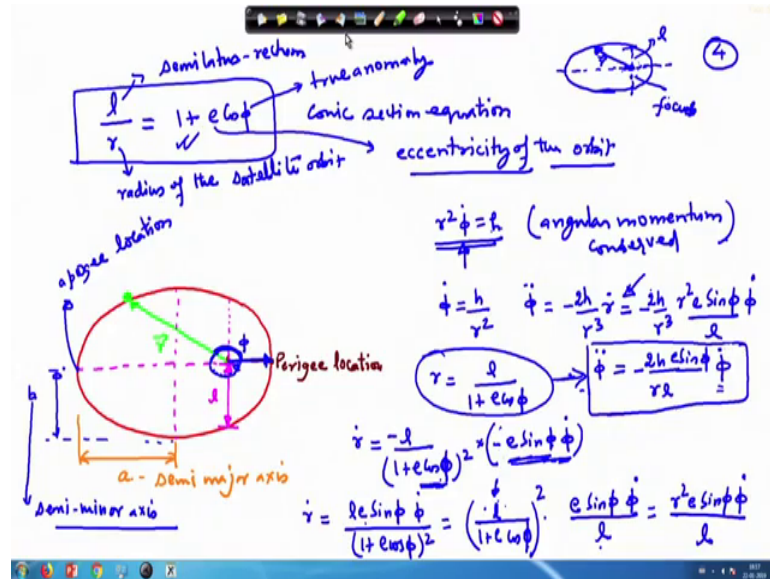
Ah what we can do I 2, we can take it on the right hand side. So, we will simplify the case here itself, theta double dot minus phi double dot is equal to plus 3 mu by r cube other equations have already been eliminated, ok. Now, if we are considering the large motion, ok. So, omega 3 and omega 1, if you look into the omega 3 and omega 1 motion. So, omega 3 is nothing, but your related to the, this is the absolute angular velocity, absolute angular velocity along the or this is the component of the component of the absolute angular velocity along the third axis body axis third along the third body axis.

So, the quantity omega 3, because we are considering that is no motion along the see what we have done that if this is my satellite, it is moving in the orbit. And there is only motion involved which is about this axis, which you are indicating by theta 2, the motion theta 1 which is one direction and three direction along this there is no motion involved. And therefore, this will drop out and the special case where I 3 equal to I 1 also this will drop out.

So, for this particular case so, for the for the case when omega 3 is equal to omega 1, this is equal to 0, ok. So, now, yawing slash rolling motion, this motion is not involved. So, in that case your equation get simplified otherwise if you take it the equation will be very complicated itself, we will see that in the elliptical orbit this turns out to be a nightmare. So, this is plus. So, this minus 3 mu by r cube I 3 minus I 1, here also we need to divided by I 2 divided by I 2 sin theta times cos theta, this equal to 0.

This is the pitching motion equation and elliptic orbit and remember that that we have not done any approximation here. So, no approximation, no approximation assumed here. So, we will explore this particular part, write in a particular way and then develop this equation.

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So, let us start with as I told you that go and refer back to refer to our lecture on space flight mechanics. So, this equation, I am writing here, it is a called the conic section, equation conic this is called the semi latus rectum. And r is; obviously, the radius of the satellite orbit and this θ I will change it to ϕ instead of using θ I would notation, I will be using here ϕ . So, this called the true anomaly and this e , it is the eccentricity, eccentricity of the orbit.

So, we are dealing with the elliptical orbit, elliptical orbit. How does it look like, something like this? Here this is the focus, this is the focus. So, your earth is situated here and because already we have discussed that this will be considered as a point mass. So, the whole if it is considered as if basically what we are, we have assumed that if earth is assumed to be circular and of uniform density. Then we can assume the earth mass to be concentrated at a point at this focus here in this case. And then your satellite is moving in this orbit.

So, this part is your from these distance to these distance, this is your focus. So, this part is l what we have written here, ok. And then if I draw a vector like this may be little bigger figure I will draw here. So, this is the focus I as discussed earlier and this is called the Periapsis or Perigee location. So, distance from here to here, this is the shortest distance. And then we have the radius vectors so, the satellite is located here in this place this angle we show this as ϕ . So, this we have written as the true anomaly,. And of

course, your l this is the distance vertical to this. So, either up or down this distance is your l from this place to this place, this is l . Then the length from this place to this place is called the semi major axis, ok.

Similarly, the distance from this point, this distance from here to here is indicated by b and this is called semi minor axis. This point is your apo axis or the apogee location, because of this discussion given in the space flight mechanics lecture. So, here this is the conic section equation. So, from there we can get the what will be the rate at which this ϕ changes. So, this is a simple differentiation and moreover, because the angular momentum is conserved in such kind of orbit. So, this equation is also applicable angular under central force motion so, this is always valid.

Now we are ready to take it take it to the next stage. So, $\dot{\phi}$ from here we get as r square and $\ddot{\phi}$. So, the rate at which this angle is changing, it is directly available to you, ok. As r changes along this direction, you can see the this is the smallest r possible from this place to this place, this is the smallest possible r as we move along the orbit. So, we go from this place to this place. So, r is varying, from here to here this r will keep varying.

. So, $\ddot{\phi}$ from this place this is $-\frac{h^2}{r^3} \dot{r}$. Now \dot{r} we have to insert here in this place to get the complete solution r equal to $l / (1 + e \cos \phi)$. Differentiate this, if we differentiate this we get here \dot{r} equal to $l / (1 + e \cos \phi)^2$, which minus $\sin \phi$ here and then differentiating this quantity.

So, this will be $e \sin \phi$ times $\dot{\phi}$ and plus $1 - \sin \phi$ here from minus $\sin \phi$ here. So, $e \sin \phi$ for this, we have written here e and then differentiation of $\cos \phi$ with a $\sin \phi$ and thereafter the $\dot{\phi}$ and $\cos \phi$ differentiation then minus sign also appears simultaneously,

So, therefore, \dot{r} which is appearing here in this place ok. So, this becomes $l e \sin \phi$ times $\dot{\phi}$ divided by $(1 + e \cos \phi)^2$. Then again utilise this particular equation so, we can write this as, utilise this. So, we can write this as $l / (1 + e \cos \phi)^2$, and then rest of the things on this side $e \sin \phi$ times $\dot{\phi}$ divided by l , $e \sin \phi$ times $\dot{\phi}$ we have taken here l , because the l^2 term is getting introduced here. So, therefore, we have divided here by l . So, this can be written as $r^2 e \sin \phi$ times $\dot{\phi}$ divided by l . So in this equation then minus h^2 / r^3 and

$r \dot{\phi}$ is $r^2 e \sin \phi$ times $\dot{\phi}$ divided by l , ok. So, minus h by $r l$ times $e \sin \phi$. So, this is your $\ddot{\phi}$, this is $\dot{\phi}$ here. Some more simplification, we can do to this equation. So, I will have to go to the next page.

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$$\ddot{\phi} = \frac{-2he \sin \phi}{r l} \dot{\phi}$$

$$= \frac{-2he \sin \phi \mu}{r h^2} \dot{\phi} = \frac{-2\mu e \sin \phi}{r h} \dot{\phi}$$

$h^2 = \mu l$ ←

⑤

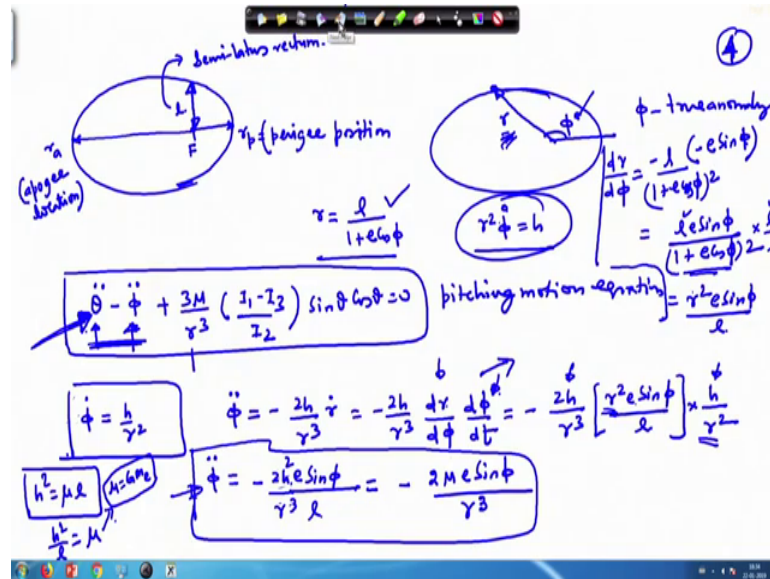
So, we have $\ddot{\phi}$ minus $h e \sin \phi$ divided by r times l and then $\dot{\phi}$. Now, h^2 this quantity is related to l by this relationship and this is derived in the space flight mechanics lecture. So, we will utilise this relationship directly here and eliminate this l , ok. So, this becomes $h e \sin \phi$ divided by r and l is equal to $h^2 \mu$ goes upside. So, this is h^2 divided by μ ah. So, μ will be upside and then we of course, we have $\dot{\phi}$ here in this place minus μ times $e \sin \phi$ divided by $r h$ $\dot{\phi}$.

And here we have missed out a term see if while h by r^2 we are writing here. So, there are 2 factor will appear here in this place. So, this 2 we have missed out everywhere, ok. So, we have to write here this 2 factor [FL].

Student: [FL].

[FL] repeat [FL], ok.

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So, what we have been discussing about the quantity or so, this is the focus and this is the shortest distance which we write as r_p and we can call as the this is related to the perigee position perigee position. And this part here from the distance from this place to this place, this is your focus F. So, this is r_a and this is your apogee location and the length from this place to this place this is the vertical distance from here to here, this is your l which we have written as semi latus rectum, ok.

So, in this figure we are measuring ϕ angle which is the true anomaly ϕ is your true anomaly this is r , ok. And for this conic section equation is given by r equal to l by $1 + e \cos \phi$. And this particle just like we can take the case of the sun earth system. So, in that the angular momentum is also conserved and it is written as $r^2 \dot{\phi}$, this equal to h . So, what we are interested is we will write the equation what we have written here. So, that we do not have to refer back again and again. So, $\ddot{\theta} - \ddot{\phi} + \frac{3M}{r^3} (I_1 - I_2) \sin \theta \cos \theta$, this is the equation we have got for the pitching moment, pitching motion, pitching motion equation, ok.

So, this part we have to resolve and write in a particular way so, that it is a convenient for analytical analysis, ok. So, we use this relationship l if r say the r is equal to l by $1 + e \cos \phi$, we are interested in finding out this quantity, this quantity is already with respect to the orbital frame, your changes are taking place this $\ddot{\theta}$, $\ddot{\phi}$, $\dot{\phi}$, $\dot{\theta}$,

we are interested in finding, ok. So, $\ddot{\phi}$ so, this is a quantity h by r^2 . So, $\ddot{\phi}$ double dot this quantity, we can write as $-2h$ by r^3 times \dot{r} , ok. And this can be further written as r is changing with respect to ϕ . So, we write it, it rather in this way.

We could have equally gone another way first finding out \dot{r} and then working out. Right now, if we differentiate, because you can see that this r is varying with respect to this ϕ , as this ϕ changing r will also change. And therefore, r is a function of ϕ and ϕ is a function of t . So, we are utilizing this information to work out here, ok.

So, therefore, $\ddot{\phi}$ this becomes $-2h$ by r^3 and \dot{r} we can get from this place. So, $\frac{dr}{d\phi}$ we can write as l by $1 + e \cos \phi$ whole square and then on the up side $-e \sin \phi$.

So, this becomes $l e \sin \phi$ divided by $1 + e \cos \phi$ whole square. And we can also summarize this as by multiplying this by in the numerator and denominator by 1 and 1 we multiply it by 1 and 1 in the numerator and denominator. So, this becomes so, l^2 square term this term and this term becomes l^2 square. So, l^2 square by this square term that we can write as r^2 , ok. And rest other things $e \sin \phi$ divided by this 1 .

So, this is the quantity which interfere in this place, this is $2h$ by r^3 and then r^2 $e \sin \phi$ divided by 1 and this quantity is $\dot{\phi}$ this is $\frac{d\phi}{dt}$, ok. This quantity is nothing, but $\dot{\phi}$. So, we will pick up it from this place. So, that becomes h by r^2 . And therefore, $\ddot{\phi}$, we can write this as $-2h$ by r^3 and r^2 r^2 it will cancel out, we are left by $2h e \sin \phi$ and here we will write as h^2 , this h and this h makes it h^2 . So, $2h^2 e \sin \phi$ r^2 this r^2 and this r^2 they will cancel out,. So, we are left with r^3 and divided by 1

And the relationship between h and l is so, h^2 divided by l this is nothing, but μ , which is the gravitational constant of the earth μ , we have written as g times m earth, or it may be any planet. So, this is $-2\mu e \sin \phi$ divided by r^3 this is $\ddot{\phi}$, ok.

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The whiteboard shows the following derivations:

$$\dot{\theta} = \frac{d\theta}{dt} = \left(\frac{d\theta}{d\phi} \right) \frac{d\phi}{dt} = \theta' \dot{\phi} = \theta' \frac{h}{r^2}$$

$$\ddot{\theta} = \frac{d}{dt} \left(\frac{d\theta}{d\phi} \right) = \frac{d}{dt} (\theta' \dot{\phi}) \quad \left| \begin{array}{l} \frac{d}{dt} (\theta') = \frac{d}{dt} \left(\frac{d\theta}{d\phi} \right) \\ = \frac{d}{d\phi} \left(\frac{d\theta}{d\phi} \right) \cdot \frac{d\phi}{dt} \\ = \frac{d^2\theta}{d\phi^2} \dot{\phi} \end{array} \right.$$

$$= \frac{d}{dt} (\theta') \dot{\phi} + \theta' \ddot{\phi}$$

$$= \frac{d^2\theta}{d\phi^2} \dot{\phi}^2 + \theta' \ddot{\phi}$$

$$= \theta'' \dot{\phi}^2 + \theta' \ddot{\phi}$$

At the bottom, a boxed equation is shown:

$$\ddot{\theta} = \theta'' \left(\frac{h}{r^2} \right)^2 + \theta' \left(\frac{-2\mu e \sin\phi}{r^3} \right) = \left(\frac{h^2}{r^4} \right) \theta'' - \frac{2\mu e \sin\phi}{r^3} \theta'$$

Now, we can proceed and work out this particular part. So, everything we need to express in terms of phi. So, theta dot again this part we can is nothing, but d theta by d t. So, we can write this as d theta by d phi times d phi by d t

And this quantity, we are going to write as theta prime means it is a derivative of theta with respect to phi. So, this is the quantity we are writing here and this part becomes phi dot obviously this quantity is known to us. So, this is theta prime times h by r square, ok. Similarly, theta double dot is the quantity which we require. So, theta double dot it is present here, and this quantity we require. So, theta double dot, we need to differentiate this again. So, this will be d by d t d theta by d t theta prime times 0.1.

So, if once we differentiate it d by d t plus theta prime phi double dot. So, phi double dot this quantity already we have worked out, theta prime this is just the derivative of theta with respect to phi which is written here. So, this quantity we need to work out. So, theta prime d by d t, we need to work out d theta by, ok. So, we have to write it in the same way d by d phi this theta is there. So, first we differentiate it with respect to d phi and then we write it like this d phi by d t d square theta by this becomes d phi square times phi dot.

So, if we insert here in this equation so, this will become d by d t d prime d square theta divided by d phi square. And then phi dot square plus theta prime plus phi double prime, ok. And this quantity, then we are going to write as theta double prime, phi double dot.

So, now look into this all these quantities are known to us. So, theta double prime and phi dot equal to h by r square. So, this becomes whole square and theta prime and phi double dot phi double dot, we have just now derived this is minus 2 mu e sin phi minus 2 mu e sin phi divided by r cube.

So, we have got the quantity here h square h square by r to the power four theta double prime minus 2 mu e sin phi divided by r cube and theta prime. So, this is our theta double dot. So, by doing so of we can express how your motion will look like as phi keeps varying, instead of t we are expressing the whole thing in terms of phi.

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Handwritten equations on a whiteboard:

$$\ddot{\theta} = -\theta' \frac{2M}{r^3} e \sin \phi + \frac{M \dot{\phi}}{r^4} \theta'$$

$$\ddot{\theta} = -\theta' \frac{2M}{r^3} e \sin \phi + \frac{M(1+e \cos \phi)}{r^3} \theta''$$

Key relation: $\frac{l}{r} = 1 + e \cos \phi$ (6)

$$\frac{I_1 - I_3}{I_2} \omega_z \dot{\theta} = 0$$

$$\ddot{\theta} - \frac{\dot{\theta}^2}{r} + \frac{3M}{r^3} \left(\frac{I_1 - I_3}{I_2} \right) \sin \theta \cos \theta = 0$$

$$-\theta' \frac{2M}{r^3} e \sin \phi + \frac{M(1+e \cos \phi)}{r^3} \theta'' + \frac{2M e \sin \phi}{r^3} + \frac{3M k_2 \sin \theta \cos \theta}{r^3} = 0$$

$$-2\mu e \sin \phi \theta' + \mu(1+e \cos \phi) \theta'' + 2\mu e \sin \phi + 3\mu k_2 \sin \theta \cos \theta = 0$$

Now, we write the equation therefore, theta double dot equal to minus theta prime 2 mu y the previous equation we have written e sin phi plus mu by r cube and go back and this is the whole thing we are writing. So, this particular part, we have to work out h square by r to the power 4.

So, h square by r to the power 4, this is mu times l. This particular part, this h square by r to the power 4 is equal to mu times l and divided by r to the power 4. So, here let us go one more step by r to the power 4 and then theta double prime, we can eliminate one term here theta prime 2 mu by r cube l by r if this equals to 1 plus e cos phi. So, use this relation so, this becomes 1 plus e cos phi and r cube term is present here.

So, $\ddot{\theta}$, so, $\ddot{\theta}$ is expressed like this. And then you can complete the whole equation and it will so, happen that because the r^3 is present here, ok. And other terms also they contain r^3 so, the this can be eliminated. Now we pick up this equation $\ddot{\theta} - \ddot{\phi}$, and everything we have already worked out. So, $\ddot{\theta} - \ddot{\phi} + \frac{3\mu}{r^3} + \frac{3\mu}{r^3} \frac{1 - \epsilon^2}{2} \frac{1 - \epsilon^2}{2} \sin\theta \cos\theta$, $\sin\theta \cos\theta$, this equal to 0. So, insert all this quantities here. So, this is your $-\ddot{\theta} + \frac{2\mu}{r^3} \epsilon \sin\phi + \frac{\epsilon \cos\phi}{r^3} \ddot{\theta} - \ddot{\phi}$.

So, this quantity $\ddot{\phi}$, we have worked out here in this place. This is $-\frac{2\mu}{r^3} \epsilon \sin\phi + \frac{2\mu}{r^3} \epsilon \sin\phi$ divided by r^3 . So, this is your $\ddot{\phi}$ term, ok. And there after the other terms, we need to put here $\frac{3\mu}{r^3}$. And let us write this term as usual $\frac{1 - \epsilon^2}{2}$, this we have written as k . So, we will write this as k , and then rest of the term $\sin\theta \cos\theta$, so this equal to 0. So, you can see that in the denominator everywhere this r^3 term is appearing and because r is non zero. So, there is no problem, we can take it on the right hand side, eliminate it from this part. Now this gets reduced to $-\ddot{\theta} + \frac{2\mu}{r^3} \epsilon \sin\phi + \frac{\epsilon \cos\phi}{r^3} \ddot{\theta} - \ddot{\phi} + \frac{3\mu}{r^3} k \sin\theta \cos\theta$ this equals to 0.

So, we can see that this is the second derivative term. So, we can write here in the beginning $\ddot{\theta} + \frac{\epsilon \cos\phi}{r^3} \ddot{\theta}$. And then anywhere the $\dot{\phi}$ term, we have to include, ok. So, this is your $\dot{\phi}$ term and $\frac{2\mu}{r^3} \epsilon \sin\phi$ this particular term, this one. So, we utilise it here and here you see, $\frac{2\mu}{r^3} \epsilon \sin\phi$ term is appearing.

So, we can write this as $-\ddot{\theta} + \frac{2\mu}{r^3} \epsilon \sin\phi + \frac{\epsilon \cos\phi}{r^3} \ddot{\theta} - \ddot{\phi} + \frac{3\mu}{r^3} k \sin\theta \cos\theta = 0$. So, this is the equation of motion in elliptical orbit for pitching only, equation of motion in elliptic orbit for pitch only. And this equation can be expanded it can worked out. So, we will continue in the next lecture. Thank you for listening.

Student: Start sir.

So, already we have observed that $\ddot{\theta}$ is the quantity which is shown here.

(Refer Slide Time: 57:21).

The whiteboard shows the following derivation:

$$\ddot{\theta} = -\theta' \frac{2\mu}{r^3} e \sin \phi + \frac{\mu(1+e \cos \phi)}{r^3} \theta''$$

$\frac{I_1 - I_2}{I_2} = k_i$ (6)

$$\ddot{\theta} - \ddot{\theta} + \frac{3\mu}{r^3} \frac{(I_1 - I_2)}{I_2} \sin \theta \cos \theta = 0 \quad \text{pitching motion equation}$$

$$-\theta' \frac{2\mu}{r^3} e \sin \phi + \frac{\mu(1+e \cos \phi)}{r^3} \theta'' - \left(\frac{2\mu e \sin \phi}{r^3} \right) + \frac{3\mu}{r^3} k_i \sin \theta \cos \theta = 0$$

$$(1+e \cos \phi) \theta'' + (2e \sin \phi - 2e \sin \phi \theta') + 3k_i \sin \theta \cos \theta = 0$$

$$\rightarrow (1+e \cos \phi) \theta'' + 2e \sin \phi (1 - \theta') + 3k_i \sin \theta \cos \theta = 0$$

If we assume $e = 0$ $\theta'' + 3k_i \sin \theta \cos \theta = 0$

So, theta double dot this is the quantity, minus theta prime 2 mu by r cube minus theta prime 2 mu by r cube and e sin phi plus mu times 1 plus e cos phi divided by r cube and times theta double prime, ok. So, in our equation of motion theta double dot minus phi double dot plus 3 mu by r cube I 1 minus I 3 sin theta times cos theta equal to 0, this is the pitching motion equation. And we need to insert the quantities respective quantities here. So, theta double dot, we can insert from this place 2 mu by r cube e sin phi plus e cos phi divided by r cube.

And then minus theta phi double prime. So, phi double prime, we have derived earlier this is minus 2 mu e sin phi by r cube 2 mu e sin phi divided by r cube which a minus sign. So, this is the quantity phi double dot and then 3 mu by r cube I 1 minus I 3, this is there is the divided by I 2 also. So, I 1 minus I 3 divided by I 2 I 3 divided by I 2, we have written this quantity as k i. So, I will replace this in terms of k i here itself, this is k i and then sin theta times cos theta.

So, if we rearrange this quantity r cube r cube from the denominator, it will get eliminated and we have mu times e cos phi times theta double prime. And moreover, if we see that this mu is also present everywhere. So, this mu also we can eliminate maybe at this stage itself. So, your mu is present here, mu is present here, it is also present here this also present here and this is a non zero quantity. So, therefore, we eliminate here and write it like this.

Then this is plus the quantity which is present here this quantity and this quantity we can combined it together to write $2\mu e \sin \phi$ again, we do not have to write here plus 2 times $e \sin \phi$, ok. And from this place minus 2 times $e \sin \phi$ times θ' and plus 3μ by r cube again this quantity goes here only 3 will remain, ok. So, if μ is eliminated r cube is a eliminated and we get $3k i \sin \theta$ times $\cos \theta$ this equal to 0 $e \cos \phi$ times θ'' plus $2e \sin \phi$, it can be taken outside the bracket this is $1 - \theta'$ plus, as we can observe this is the this is a complicated equation. Now under simplified condition this can be solved.

So, if we assume then this gets simplified to we can see from this place this term will get eliminated and we get here θ'' , ok. This term we can get rid off and we can write this as $3k i \sin \theta$ times $\cos \theta$, this equal to 0 . So, this equation is written here only for the e value, we shall continue with this lecture in the; we will continue with this topic in the next lecture and we are stop in the meanwhile.

Thank you very much.