

Thermodynamics (Classical) for Biological Systems

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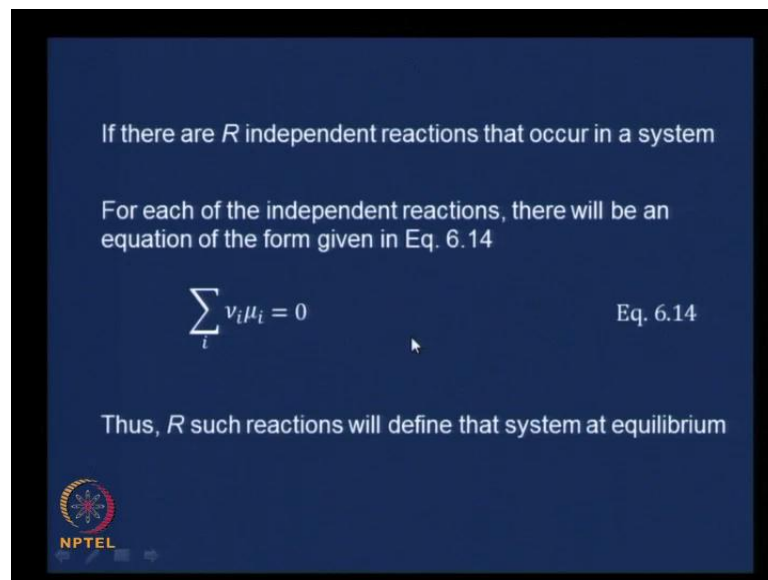
Module No. # 06

Reaction Equilibria

Lecture No. # 35

Phase rule for Reacting Biosystems, Equilibrium constants

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If there are R independent reactions that occur in a system

For each of the independent reactions, there will be an equation of the form given in Eq. 6.14

$$\sum_i \nu_i \mu_i = 0 \quad \text{Eq. 6.14}$$

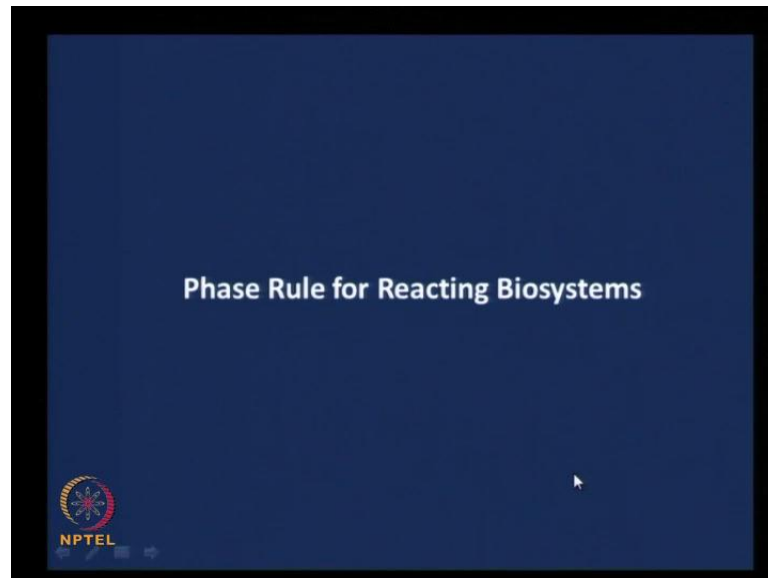
Thus, R such reactions will define that system at equilibrium



Welcome! to this class.

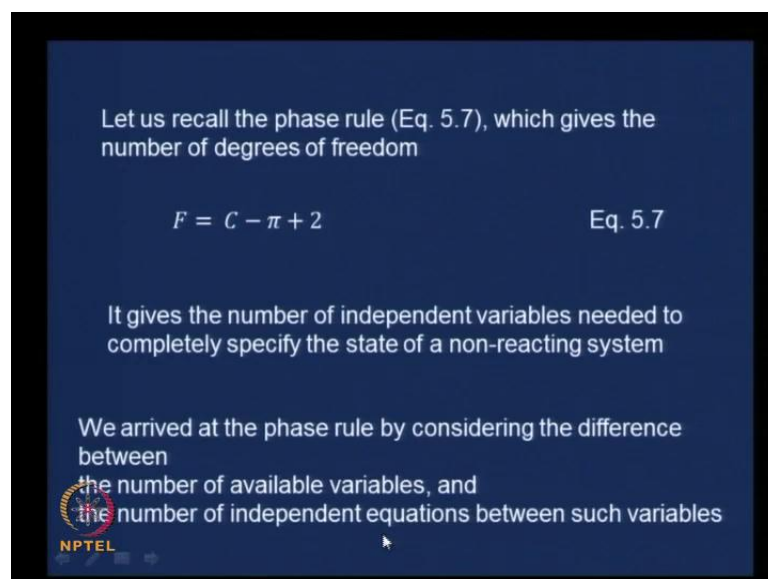
In the last class, in this module on reaction equilibrium, we saw one of the fundamental relationships for reaction equilibrium, and that happens to be for an independent reaction the sum over i $\nu_i \mu_i$, this is a stoichiometric co-efficient, μ_i is a chemical potential of that species i , the product of that, and sum over all i that equals 0. We also said that if there are R independent reactions that occur in our system, then we have one such equation for each of the R independent reactions. Therefore, R such reactions will define that system at equilibrium. This is one of the very significant results, equation 6.14. It is very nice and compact, and it is good to remember this. This will automatically become a part of the memory; the recall is very easy for this equation.

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What we are going to do next is to look at phase rule for reacting biosystems. The way we are going to do that is through some sort of an interactive method. We have already derived the phase rule for non-reacting systems in the previous module. What I am going to do is, give you some thoughts again on the background to that, and I will give you all the data; and then let you come up with the phase rule for reacting biosystems. I will give you time for that in this class.

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If you think about, or if we recall the phase rule, equation 5.7, and that gave the number

of degrees of freedom. In other words, the number of degrees of freedom is the number of independent variables that are required to complete the description of the system. We saw that the phase rule for the non-reacting system was F the number of degrees of freedom is C minus π plus 2; C , the number of components in each phase of the system, π the number of phases. This was our equation 5.7 in the previous module. Let us recall, that we arrived at the phase rule by considering the difference between the number of available variables, and the number of independent equations between such variables.

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In the case of a system with reactions, at equilibrium
the number of variables is the same, $\pi(C+1)$
 π : the number of phases
 C : the number of components

The number of independent equations between the variables:
 $(\pi - 1)$ equalities of temperature
 $(\pi - 1)$ equalities of pressure
 $C(\pi - 1)$ equalities of chemical potential

In addition to the above, which were also applicable for a non-reacting system, for a reacting system with R independent reactions

R conditions of equilibrium of the form $\sum_i v_i \mu_i = 0$

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In the case of a system with reactions at equilibrium, the number of variables is the same, π into C plus 1. There are no additional variables that are there. As mentioned earlier, π is the number of phases, and C is the number of components. The number of independent equations between the variables is π minus 1 equalities of temperature ... recall that T_α equals T_β and so on. So, if there are π phases, you could consider that one composite equation as π minus 1 equations, or equalities of temperature. In a similar fashion there are π minus 1 equalities of pressure that corresponds to equation 5.2, and π minus 1 equalities or equations of chemical potential for each component. We can write one for each component and therefore, the total number of such equalities of the chemical potential is C into π minus 1.

So, all this remains the same even for a system with reactions at equilibrium. The only additional thing is that for a reacting system with R independent reactions, we have R

conditions of equilibrium of the form, $\sum \nu_i \mu_i = 0$. I am going to leave you with this background. This is more than enough background to derive the phase rule for reacting biosystems. Please, go ahead and do it. You can take about 15 minutes since you doing it for the first time. It will give you a lot more confidence in the material itself. Go ahead and do it, and I will present the solution or the phase rule for reacting biosystems after about 15 minutes. Go ahead please.

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So, we all know the background; the number of variables in the reacting system at equilibrium, the number of independent equations between the variables – all this is same as the case of the system without reactions. The only additional thing being these R conditions of equilibrium of the form $\sum \nu_i \mu_i = 0$.

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Thus the total number of available equations

$$(\pi - 1)(C + 2) + R$$

Therefore, the number of degrees of freedom for a reacting system is

$$F = \pi(C + 1) - \{(\pi - 1)(C + 2) + R\}$$

which can be simplified to

$$F = (C - R) - \pi + 2 \quad \text{Eq. 6.15}$$

The number of degrees of freedom reduces by the number of reactions, compared to a non-reacting system

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So, how did we go about deriving the phase rule earlier? We looked at the total number of available equations which is $\pi - 1$ into $C + 2$, as for the case without reactions plus this R. The number of variables minus the number of available equations, available independent equations between those variables, actually, gives a degrees of freedom that we saw. Therefore, F the number of degrees of freedom is π into $C + 1$, the number of variables, minus the number of equations between them, independent equations between them. If we rearrange this, essentially some algebra, we would get $C - R - \pi + 2$.

So, as you can realize whatever was C earlier can be replaced by C minus R . So, C minus R minus π plus 2 – hopefully you got this in your own efforts – gives the number of degrees of freedom in a system with R reactions. Let us call this equation 6.15. ... As can be seen, the number of degrees of freedom reduces by the number of reactions, compared to a non-reacting system.

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Let us move forward. Now, let us look at equilibrium constants. I am sure you know what equilibrium constants are. You must have done this in your 11-th standard, 12-th standard chemistry, and probably in some other courses in the engineering discipline or in the biotechnology course before you took this course. So, you already know what equilibrium constants are, what I am going to give you here is a formalism from the fundamentals to get at equilibrium constants. So, you will know where it all rests on, and so it will be much easier for you to have an overall picture in terms of some fundamental thermodynamic quantities.

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Let us first consider a perfect gas mixture


As we saw in Module 4, Eq. 4.1, the expression for chemical potential of species, i , in such a case is

$$\mu_i = \mu_i^0 + RT \ln p_i \quad \text{Eq. 4.1}$$

Substitution of Eq. 4.1 in Eq. 6.14

$$\sum_i \nu_i \mu_i = 0 \quad \text{Eq. 6.14}$$

yields

$$\sum_i \nu_i (\mu_i^0 + RT \ln p_i) = 0$$


To do that, let us first consider a perfect gas mixture. As we saw in module 4, especially equation 4.1, the expression for chemical potential of species i , in the case of a perfect gas mixture was, μ_i equals μ_i^0 plus $RT \ln p_i$, which is the partial pressure of component i in the mixture. If you substitute equation 4.1 in equation 6.14 – equation 6.14 is the condition for reaction equilibrium $\sum \nu_i \mu_i = 0$. Now, we are going to substitute μ_i with $\mu_i^0 + RT \ln p_i$. If we do that then we could write $\sum \nu_i \mu_i$ into $\sum \nu_i (\mu_i^0 + RT \ln p_i) = 0$. This is this condition here.

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which can be written as


$$\sum_i \nu_i \mu_i^0 + RT \sum_i \ln p_i^{\nu_i} = 0$$

or

$$\sum_i \nu_i \mu_i^0 = -RT \ln \prod_i p_i^{\nu_i} \quad \text{Eq. 6.16}$$

Let us define

$$\prod_i p_i^{\nu_i} \equiv K_p \quad \text{Eq. 6.17}$$

 K_p : the equilibrium constant based on partial pressure

Which can, of course, be written as we just multiplying the terms and combining them appropriately using the definitions of log and the definitions of log of a to the power b and so on. We are going to do that now. So, if you do that sum over i $\nu_i \mu_i$ naught plus RT , is a constant here, T is a constant here ... sum over i \ln of p_i power ν_i , if you recall, this would have been $\nu_i \ln p_i$ which is, \ln of p_i power ν_i . You can work this out by looking at the properties of the log function. And this is summed over all i ; this equals 0. Or you can write this as $\nu_i \mu_i$ naught sum over i equals minus RT – this is some of log, which is nothing but log of the product; that we all know.

Product of ... as represented by this big pi here ... product of p_i power ν_i . Let us call this equation 6.16. Let us define this particular quantity, product over i p_i power ν_i the stoichiometric co-efficient, the partial pressure time raised to the power of this stoichiometric co-efficient, for each of the species involved, as K_p . You would recall K_p , the equilibrium constant with respect to partial pressure. That, essentially comes from this. So, product over i p_i power ν_i is defined as K_p ; equation 6.17.

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Thus, we can write Eq. 6.16 as

$$-RT \ln K_p = \sum_i \nu_i \mu_i^0 \quad \text{Eq. 6.18}$$

To illustrate the above, since we are familiar with equilibrium constants from earlier courses, let us consider the reaction mixture given by Eq. 6.1, as a perfect gas mixture

$$2A + B \rightarrow 3C$$

Then, Eq. 6.18 can be written as

$$-RT \ln \left(\frac{p_C^3}{p_A^2 p_B} \right) = 3\mu_C^0 - 2\mu_A^0 - \mu_B^0 \quad \text{Eq. 6.19}$$


And therefore, we can write 6.16 as minus $RT \ln K_p$ equals sum over i $\nu_i \mu_i$ naught. Therefore, we have a means of getting K_p in terms of the μ_i naught values, the standard chemical potentials of i species multiplied by the appropriate stoichiometric constants and summed over all i would lead to minus $RT \ln K_p$. Let us call this equation 6.18. To illustrate the above ... to get it into familiar territory ... since we are familiar

with equilibrium constants from earlier courses, let us consider the reaction mixture given by equation 6.1, which 2A plus B giving you 3C. Let us consider that for a minute to be a perfect gas mixture.

In such a case, this is the equation ... equation 6.18 for this particular reaction can be written as minus RT ln K p. What was ln K p? ... Ln of the product of p i power nu i. You know that nu i is positive for the products and negative for the reactants. And therefore, if you multiply that all together it be p C cubed into p A to the power of minus 2 into p B to the power of minus 1, which can be written as p C cubed divided by p A squared p B ... and of course, the ln comes from here. This equals 3 mu C naught minus 2 mu A naught minus mu B naught. Sum over i nu i mu i naught 3 mu C naught minus 2 mu A naught minus mu B naught. We will call this equation 6.19. Therefore, this must be very familiar to you ... in terms of the equilibrium constant. How do find that if given the partial pressures is like this, and you can relate it to the standard chemical potentials for that particular reaction.

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For a **perfect gas mixture**, other equilibrium constants based on concentration or mole fraction can be arrived at

For the *i* th component

$$p_i = \frac{n_i}{V} RT = c_i RT \quad \text{Eq. 6.20}$$

Substituting Eq. 6.20 in Eq. 6.17, and transposing the equation:

$$K_p = \prod_i (c_i RT)^{\nu_i} = (RT)^{\sum \nu_i} \prod_i c_i^{\nu_i} \quad \text{Eq. 6.21}$$



Let us stay with the perfect gas mixture for a little while longer. As you can recall, there are many other equilibrium constants. The earlier one was based on partial pressure. You could have an equilibrium constant based on concentration, based on mole fraction, and so on. They can be arrived at as follows. For the *i*-th component the partial pressure can be written as, n_i by $V RT$. This is from the ideal gas equation. Perfect gas mixture

... so, mixture of ideal gases and so on. So, n_i by $V RT$ is partial pressure. ... n_i by V ... is nothing but the definition of the concentration, the molar concentration. Therefore, $c_i RT$. Let us call this equation 6.20 in this context.

If we substitute equation 6.20 in equation 6.17 ... that will become very clear right now ... and transposing, we will get K_p equals p_i into $c_i RT$, the whole power ν_i . Can you do this bit of algebra right now? Just substitute 6.20 in 6.17 and get this expression. Take about 5 minutes to do that which can further be written as this is just you know the algebra part the log and raising it to powers and so on. You need to get comfortable with this. So, I will give you about 5 minutes to do this. Go ahead please.

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You should have arrived at K_p equals product over i $c_i RT$ the whole power ν_i which can of course, be written as RT to the power of sum over ν_i times the product of c_i power ν_i . Let us call this equation 6.21. This is the definition of K_p .

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If we define


$$\prod_i c_i^{\nu_i} \equiv K_c \quad \text{Eq. 6.22}$$

Then, we can write Eq. 6.21 as

$$K_c = K_p (RT)^{-\sum \nu_i} \quad \text{Eq. 6.23}$$

Further, recognizing that

$$p_i = P y_i \quad \text{Eq. 6.24}$$



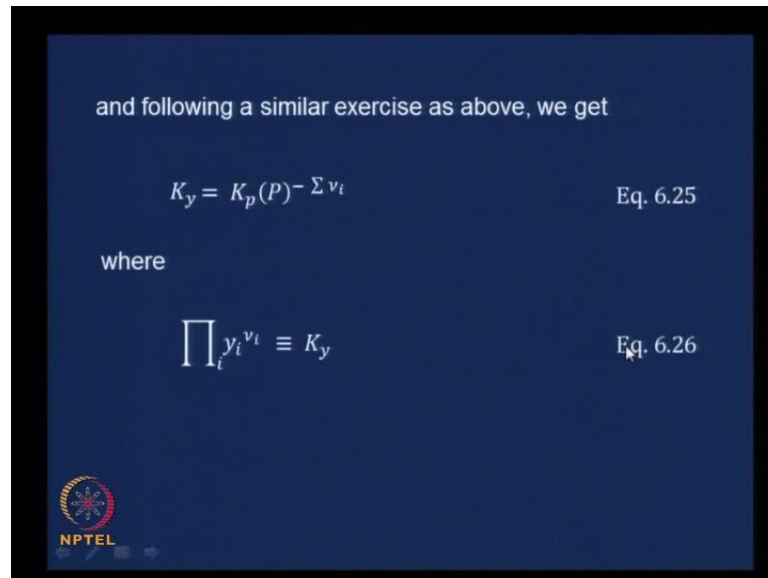
If we call this as K_c , product over i c_i power ν_i as K_c . Then, it is easy to see that K_c is nothing but K_p into RT raised to the power of minus sum over all these stoichiometric constants, ν_i . Let us call this equation 6.23. Just not this, there is probably one more. Let us recognise that p_i is nothing but P into y_i - equation 6.24. Let me leave you with this for another 10 minutes to derive K_y . You know the equilibrium constant based on

the mole fractions y_i . Can you go ahead and do it on the same lines as we did it for this. I will present the results after 10 minutes, go ahead.

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Taking p_i equals the total pressure times the mole fraction y_i .

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and following a similar exercise as above, we get

$$K_y = K_p (P)^{-\sum \nu_i} \quad \text{Eq. 6.25}$$

where

$$\prod_i y_i^{\nu_i} \equiv K_y \quad \text{Eq. 6.26}$$

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If we follow a similar exercise as above ... write ... K_y in terms of K_c and so on, you should get K_y is nothing but K_p into P raised to the power of minus sum over all ν_i . Hopefully, you got this in your own exercise. Where our ... K_y is defined as product over i y_i the mole fraction raise to the power of ν_i , the stoichiometric co-efficient. Let us call that equation 6.26.

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On the other hand, if the gas mixture is imperfect, then the chemical potentials need to be expressed in terms of fugacities as in Eq. 4.2

$$\mu_i = \mu_i^0 + RT \ln \hat{f}_i$$

A similar exercise as for K_p above, leads to

$$-RT \ln K_f = \sum_i \nu_i \mu_i^0 \quad \text{Eq. 6.27}$$

K_f , the equilibrium constant based on fugacity, is defined as

$$\prod_i \hat{f}_i^{\nu_i} \equiv K_f \quad \text{Eq. 6.28}$$


Now, all those equations that we derived so far, the equilibrium constants, were applicable for perfect gas mixtures. If the gas mixture is imperfect which is usually the case, then, the only change that we need to make, since we are working in terms of chemical potentials is that, the chemical potential needs to be expressed in terms of fugacities instead of partial pressures. This ... was given in equation 4.2. Chemical potential in terms of fugacities, μ_i equals μ_i^0 plus $RT \ln \hat{f}_i$. Replace fugacities wherever there were partial pressures earlier. If we do a similar exercise ... I leave this for you to do as a homework.

So, if you do a similar exercise as for K_p we should get minus $RT \ln K_f$. K_f is the equilibrium constant based on fugacities, as equal to sum over i $\nu_i \mu_i^0$ where K_f , the equilibrium constant based on fugacity is defined as, product over i $\hat{f}_i^{\nu_i}$ raised to the power of ν_i ; that is our K_f . Let us call that equation 6.28. ... We are almost out of time. When we meet in the next class, we will take things further. See you then.