

Computational Fluid Dynamics

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Solution of Navier-Stokes equations

Lecture No. # 20

Pressure equation method for the solution of NS equations

We can see that there are inter cases solution of navier stokes equations. And its not easy to get a solution in a very straight forward way. We have to do something more in order to squeeze out the solution using the conventional C F D method for their generic transport equation. We have seen two methods, the artificial compressibility method and the stream function ψ method for the solution of all the four equations that appear in the navier stroke's equation. Both the methods have the limitations. Now we will look at the one method which has no such limitations. It can be applied for time dependent flows in order to time evaluation of u, v, w as a function of x, y, z, t . And it will also be used for three dimensional flows. So it is a method which is applicable for fully three dimensional transient flows. So the method that we are referring to is what is known as pressure equation method.

We know that we can solve the momentum equations provided; we know what the pressure is. Now we do not know what the pressure is we have get it from the continuity equation, and for incompressible flows pressure is not there as a variable in the continuity equation. So that is where we have tried to eliminate pressure in the stream function ψ method or we have introduced ψ the pressure in the form of artificial compressibility in the other method. Now what we do is that, now we will try to solve the momentum equations along with the Poisson equation for pressure that we have derived earlier. Now we will put it in the general three dimensional context and then we will see how we can evaluate both of this step by step in such a way that we get a velocity and pressure field which together obey the navier stokes equations without the attendant limitations of either the artificial compressibility method or the stream functional method. So the idea is that we write down the momentum equations, we derive the equation for pressure and then we try to come up with the algorithm but, that which we can implement the solution of this two equation together. So that is what we will do.

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Pressure Equation

(A) $\frac{\partial u_i}{\partial t} + \frac{\partial (u_i u_j)}{\partial x_j} = -\frac{1}{\rho} \frac{\partial p}{\partial x_i} + \nu \frac{\partial^2 u_i}{\partial x_j^2}$

(B) $\frac{\partial u_i}{\partial x_i} = 0$

$\frac{\partial (A)}{\partial x_i} \Rightarrow \frac{\partial}{\partial x_i} \left(\frac{\partial p}{\partial x_i} \right) = -\frac{\partial}{\partial x_i} \left[\frac{\partial (u_i u_j)}{\partial x_j} \right]$

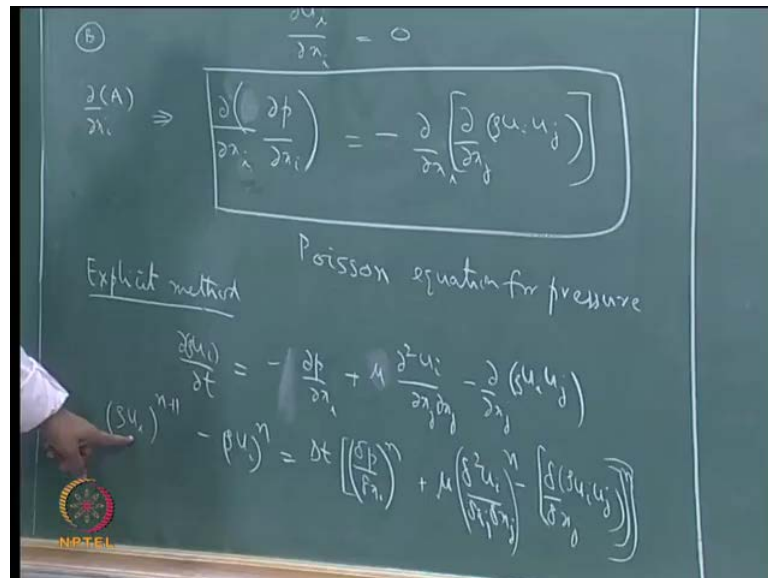
Poisson equation for pressure

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We will call this as pressure equation approach. For simplicity we will use the index method to write on the momentum equations like the X J U I U J. Understood that there is summation convention employed here, 1 by root make it high vise I plus mew del square u i by. We have left out gravity it's not a big problem. So we will just use these equations. These are the equations that are three momentum equations and we have Mew I by D X I equal to 0 as the continent equation. So these are the equation that we have solved and this is the generic pressure transparent equation we can solve this provided we have the pressure and what we tried to do is that we take the divergence of this. That is we take we call this as equation A and equation B here. Then we take D by D X i of A given that each time some i is appearing here so that we effecting taking the divergence. And we can show that in this particular case and we take the divergence for incompleation flow with cost and proper place and with that continent equation, we can get an expression for.. Essentially this is nothing that the aplsmium of the pressure. So this is equal to minus 2 row. Let we see it is correct. This must be, I wrote in this.. Let we just verify. If this is the case, yes. So this is what Poisson equation for pressure. And this as we have mentioned earlier we have written down the two dimensional form of this in the stream functional wrath city method. And this can be solved provided the right hand side of the Poisson equation is known. So one idea would be solve this within assumed pressure and get the velocities, put the velocities here, get the pressure and the get use this pressure to go back and calculate this and then get a velocity field, we can do this kind of equation. So that is one way for solving this. Easier, probably way of

implementing this for an explicit, when we consider an explicit its solution moisture equation for a time and accurate method is something that can be done in a simpler way like this.

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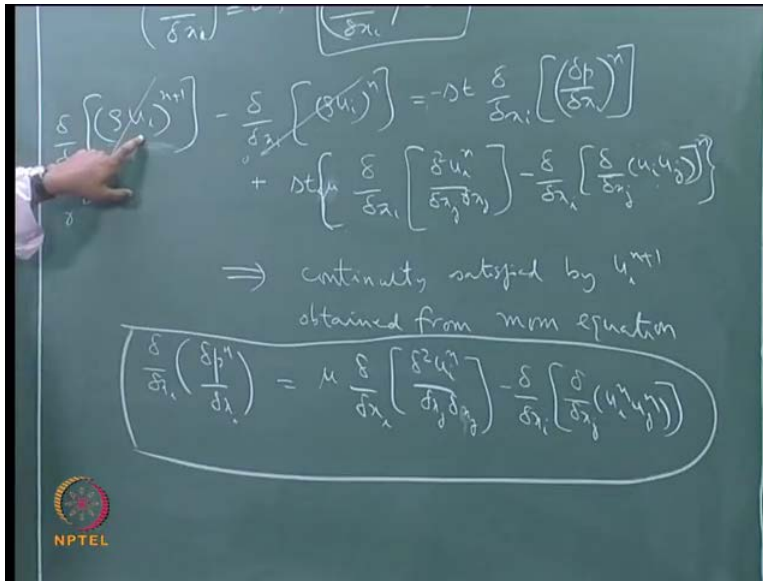


Let's consider explicit methods. Now we take this consider, we know that we have the advection time and we have diffusion time and all these things. And we can rewrite this as minus and, so we have the left hand side term and then we have right hand side terms involving these special derivatives. Out of this, we have... if we did not have this then it will be straight forward through evaluate this because, that has only the velocity components. Now what we say is that this is determined explicitly so we can write down from this. For the sake of nothing without missing any generality we can take density here, we can make this dynamic viscosity and we can put density here. And all the way I am dealing with incompressible flow we have to remember that the incompressibility that is coming from the pressure versus velocity relation for example, as an encapsulate in the Bernoulli equation. We can have other reasons for density variation, which is not part of encapsulation pressure for example, you can have air at, to say at one atmospheric pressure and it may be that you have a reaction which is going which a for example, oxygen is in observed and carbon dioxide is released. So you can have the composition of the air the gases medium is changing from location to location and as a composition changes density changes. So that that kind of variation is not part of the incompressible flow consumption. So that is why one reason to put a row inside here

because that takes care of those kinds of situations where this **two** variation is not arising from the compressibility consideration but, from other consideration such as comprehension changes. So we can put it down without lose of generality without affecting the assumptions that we are made in the terms of incompressible flow at the time step limitations and all that. So we put it like this and we can make use of discretization or we can say that, μu in plus 1 minus μu in m is equal to Δt times. We can call this as Δt by Δt in an explicit method this is evaluated at n times using some discretization so that we are replacing the partiality reaction with Δt to indicate that is discretized part we are not specifying what the discretization is that is immaterial to this as of made plus e times Δx by Δx . So this is a discretization in an explicit way of the momentum equation so that we can get row I at n plus 1 for laws for the pressure gradient at the e times and the discretization of this all evaluated at n times. So this method can be used provided we know. Now what we are to do that because pressure and velocity are not known readily we have to make sure that we put the correct pressure gradient in evaluating this.

So we want to make sure that is the pressure gradient is evaluated or the pressure is evaluated properly at the n th time step, so that using this discretization for the momentum equation will give us a correct value of velocity at n plus 1. So we want to make sure that in solving this using this express method to go from μu in I to μu in I plus 1, we put the correct pressure gradient here at n . So the evaluation of the pressure is now post in a different way it post in such a way that when we put the pressure gradient n th times step in this discretized equation, we get the correct value of μu in m plus y . So what is meant by correct evaluation so that velocity is satisfying the momentum equation the discretized momentum equation. But, it also has to satisfy the discretized continuity equation at m plus 1. So we say that the pressure field the pressure evaluation at m plus and such that the evaluated velocity of m plus 1 must satisfy the continuity equation.

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Now what is continen

equation it is $\frac{d}{dx} u$ by $\frac{d}{dx} x$ it is equal to 0. So in a discretized form is equal to 0. So the n plus the n is equal to 0 and similarly, n plus 1 is equal to the 0 so this must also be equal to 0. So if we can impose this condition in evaluating the pressure gradient the n , then we can make sure that we have that the correct velocity field here. **Right**. Now we have to get an expression for this. So what we do is we take the divergence of this. Just as we have taken divergences of $\frac{d}{dx}$ of this divergence of discretized equation and therefore, we can say that, $\frac{d}{dx} u$ at n plus 1 that is the divergence of the first term, minus $\frac{d}{dx} u$ by $\frac{d}{dx} x$ of $\frac{d}{dx} u$ at n is equal to $\Delta t \frac{d}{dx} \left[\frac{d}{dx} u \right]$, that is the this one plus Δt and we have μ here coming for this. And then we have $\frac{d}{dx} \left[\frac{d}{dx} (u \cdot u) \right]$ of this, if now take it inside. Let me just go after this and then we have solving this one. Let's see this says square brackets square brackets, so just we will put this as it is. So this is the divergence of this discretized equation and we have to put the correct indices this is n here and here and this is n and n , this one is evaluated essentially by making n values for this. Now here, we can impose the condition that this is 0 for **comparative** so at this time must be 0 and this is the condition we want to apply here. So we say that for the velocity field at $\frac{d}{dx} u$ at n plus 1 to satisfy this then that the divergence of equation that we are representing for a discretizing equation must be such that this term is 0 so this term is any way 0 because the velocity at the n th time should be satisfied continuity. Otherwise we can also make corrected this. So we have this condition here. So provided, the pressure gradient at n plus one is such that divergence of the pressure gradient is equal to the divergences of the viscous term and the advection term here. Then in which case will have continuity equation **continuity equation** being satisfied by $\frac{d}{dx} u$ at n plus 1 obtained from **from** the momentum equation. So the argument that we are saying is: we have the

momentum equation and this allows us to calculate the velocity at n plus 1 by discretizing it appropriately in a way that we choose and for this equation to satisfy the continuity equation the computed velocity at m plus 1 to satisfy the continuity equation, we have to evaluate the pressure gradient the pressure at m in such a way that this satisfies this equation. So it satisfies this, if I say that this, that is a minimize that if this is equal to, so if we evaluate pressure from this and use that pressure to evaluate the pressure gradient here, which is then wise to evaluate the new u m plus 1 then we have the guarantee that momentum equation is being satisfied by new u m plus 1 and also the continuity is being satisfied by the velocity field that is presented. So we are enforcing the condition of satisfaction of continuity by choosing to evaluate the pressure gradient such that satisfies the continuity. In that sense we have chosen a pressure field p i j k at n plus yet n value the previous step value to the evaluated from the previous values of the velocity field. See on the right hand side we have the other velocity and velocity that is appearing and velocity is at the n th time step, which are known neither by **ensile** conditions are somehow. So if you know the velocity field then we can make use of known velocity field to get a pressure field at n th time step, which when substituted into the explicit form of the momentum equation, will **engrave** calculate a velocity field at n plus 1 **it** satisfies continuity equation. So this is the basic pressure equation method in the explicit way..

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The chalkboard shows the following mathematical derivation:

$$\left(\frac{\delta u_i}{\delta x_i}\right)^n = 0, \quad \left(\frac{\delta u_i}{\delta x_i}\right)^{n+1} = 0 \quad (E)$$

$$\frac{\delta}{\delta x_i} \left[\left(\frac{\delta u_i}{\delta x_i}\right)^{n+1} \right] - \frac{\delta}{\delta x_i} \left[\left(\frac{\delta u_i}{\delta x_i}\right)^n \right] = -\Delta t \frac{\delta}{\delta x_i} \left[\left(\frac{\delta p}{\delta x_i}\right)^n \right] \quad (F)$$

$$+ \Delta t \mu \left[\frac{\delta}{\delta x_i} \left[\frac{\delta^2 u_i}{\delta x_j \delta x_j} \right] - \frac{\delta}{\delta x_i} \left[\frac{\delta}{\delta x_j} (u_i u_j) \right] \right]$$

⇒ continuity satisfied by u_i^{n+1} obtained from mom equation

$$\frac{\delta}{\delta x_i} \left(\frac{\delta p}{\delta x_i} \right)^n = \mu \frac{\delta}{\delta x_i} \left[\frac{\delta^2 u_i}{\delta x_j \delta x_j} \right] - \frac{\delta}{\delta x_i} \left[\frac{\delta}{\delta x_j} (u_i u_j) \right]$$

So we start with u_i so that can u_i at $n=0$. This is the initial time that is part of the specification of the initial conditions so that is u_i . So once this is known then use that to evaluate the right hand side of if you call u_i .

So we have equation A here and then we have equation B and equation C and we can see that this equation and this equation are similar this term will go to 0 and we have that one. So this u_i u_j of $u_i \times u_j$ and that is exactly what we have here. In a discretized form there maybe some left over them so this time is expected to be very small but, we are enforcing that one create any problem. So we have equation C here and this is, let's call this as equation D and this is our equation E and this is equation F and finally, we have equation G here.

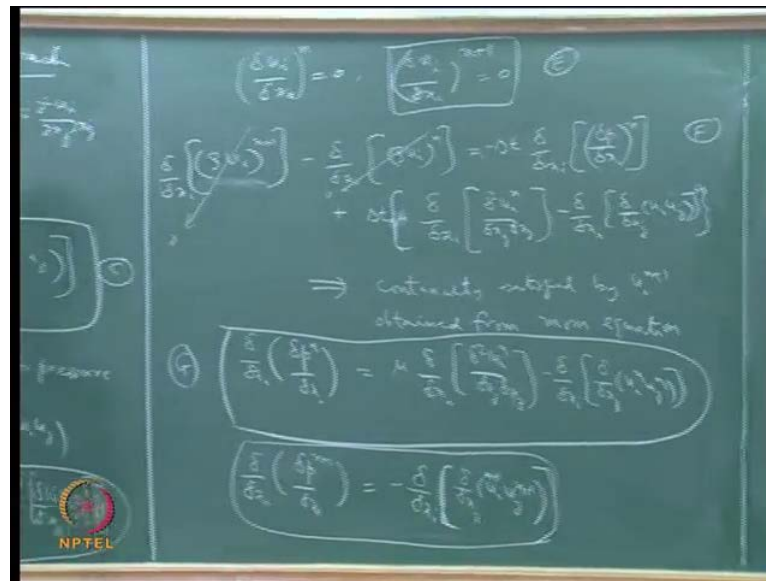
So we are making use of the pressure gradient obtained by equation G to substitute in this, in order to that velocity field u_i from known velocities. So if known velocities at n th time step, evaluate the pressure in n th time step, such that it satisfies the continuity equation and then using the pressure evaluated from this Poisson equation we evaluate the pressure gradient here and then we substitute in this and we calculate the u_i at $n+1$. So that is the method that we can use.

So use equation G to get p and use equation D to get u_i at $n+1$. So those are the three ways, with this by you can go to use equation G to get p at $n+1$. From this is that n we make the use of velocity field you can find out. So now you can go to the $n+1$ equation D to get u_i at $n+2$ and then equation G here to get p at $n+2$ and solve and so the evaluate a pressure here for a more velocity field from a main velocity field, subject to the condition that the evaluated pressure p put in the momentum equation, will give us a velocity at $n+1$ step which satisfies the continuity equation. So this requires us to solve a Poisson equation for pressure in which all the right hand terms are known in the explicit way and once we have that we have to evaluate the gradient in a consistent way. So the discretization on this terms has to be done in a proper way that we don't get a instability and lack of consistency and all that. Then we have this equations.

And here this is an explicit equation you have three equation and you can evaluate row u row v row w and then $n+1$ and then use those thing and come back and evaluate p at $n+1$ and solve. So this is an explicit method and this way we can look at how u v w

and p evolve as a function of x , y , z and t . And we note here this is over all method in this we are not made an assumptions about either two dimension or **or** compressibility or any such thing. We are just evaluating them in such way that we are solving equation A and equation B. Where are we solving equation A in the discretized form in equation d to get ρ^{n+1} and how are we solving equation B, Here we are putting it in a discretized form here in order to get pressure field. So the pressure field is seen here as a continuity enforcing variable. In incompressible flow pressure has more significant, whether it is highlight bar or one bar or point one bar it is the same. The flow only reacts to pressure gradients, pressure variation from one point to another point this spatial velocity gradient. It does not react with the absolute value of the variant. It the pressure has low influence in the properties except outside calculation done over here. So you may have to evaluate the correct value of ρ and **through** at a particular pressure. But during the flow because of the pressure variance we do not have a variation of ρ of density and that is understanding in this So that density and pressure relation is taken out and pressure comes as a **fictitious** variable in enforcing continuity and that is what we have done here. So we have evaluated the pressure in such a way that continuity equation is satisfied for the velocity field at $n+1$. So this is a pressure equation approach, pressure equation method for the solution of the Navier Stokes equation. We have iterative explicit method and **explicit method** especially for coupled equation is obviously restricted to very small types **strips**. So you can have a we would have a explicit method if we want to have a implicit method then, here we will have a $n+1$ and $n+1$ and $n+1$ so that makes a you have $n+1$ here. This is the still there this is $n+1$ and then this will be $n+1$ this will be $n+1$. So we have to solve this in the we can write down the corresponding implicit version we will leave this for a time being.

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So we have this is equation G that has to be solved along with D one in which will have $n + 1$ here, $n + 1$ here and $n + 1$ here. So in order to solve this equation in order to solve this equation we need to know pressure at $n + 1$ and in order to solve pressure at $n + 1$ we need to know u I pressure at $n + 1$. So that makes you difficult. So it becomes an implicit solution. So that means that as you go from time t equal to n new $n + 1$ first of all we have **monelarity**. So you have the new n is some of the term appear here and then we have to stop at this value that is from the previous separation and then use that here using the n star and using the n star you will get a p star and then you supply the p star and you can make it and take some of these to the right hand side and then you will be able to get a new value of velocity and use the new value of the velocity here to reevaluate to the right hand side that will be a new value of pressure here. And again put the pressure and the new value of velocity set we have got to get another value of velocity. So in that way we have to solve **retalatively** G I and the correspond D I in order to get velocity field and the pressure field which satisfies this equation and this equation. When we do that, then we have an implicit method which probably allows us to take much longer time steps than explicit method. But, even one can see that which is also more complicated because at each time steps you have to solve this equation and this equation several times. Firstly with u_n and p star and then we substitute here to get you $i + 1$ star and then we put that value here and get this and then we have to keep on doing this **assignment** several times until we get the conversions. And so there is more work to be done if I use implicit method on this. But, the advantage that we gain is that you can

take longer time steps than what is possible with the explicit method. How much longer and whether the additional work that you have to do for the solution of the equation is worth it depends on a number of factors including what methods we use for the solution of this and this. Because this is a Poisson equation so this will be converted into $\nabla^2 \phi = b$ as equal to b several solutions of these this A here remains the same because it is only the right hand side that is changed this. In such a case when you have to solve $A\phi = b$ several times in which A is a matrix then maybe we can do more efficiently by using some like early decomposed method or somewhere once we convert A into some other diagonal matrix or triangular matrix then the solution future to take solution of this will become easier so we can reduce that convolutional time required and repeated solution by taking specific advantage of the nature of this equation and what we are all $\nabla^2 \phi = b$.

If you look at this equation here in an implicit way if we are treating this full implicit way we have to bring this to the left hand side and also bring this also to the left hand side and this goes to the right hand side. So when you look at one equation to other equation this term won't change this term will change. And both this term will change so because of this, the matrix corresponding to this will be changing to equation to equation. But, the matrix corresponding to this will not change from equation to equation because their left hand side is fixed but, only the right hand side will change.

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The chalkboard contains the following derivations:

$$\left(\frac{\partial u}{\partial x}\right)^2 = 0, \quad \left(\frac{\partial u}{\partial x}\right)^2 = 0 \quad (1)$$

$$\frac{\partial}{\partial x} \left[\left(\frac{\partial u}{\partial x} \right)^2 \right] - \frac{\partial}{\partial x} \left[\left(\frac{\partial u}{\partial x} \right)^2 \right] = -\frac{\partial}{\partial x} \left[\left(\frac{\partial u}{\partial x} \right)^2 \right] \quad (2)$$

$$+ \frac{\partial}{\partial x} \left[\frac{\partial}{\partial x} \left(\frac{\partial u}{\partial x} \right)^2 \right] - \frac{\partial}{\partial x} \left[\frac{\partial}{\partial x} \left(\frac{\partial u}{\partial x} \right)^2 \right]$$

\Rightarrow continuity equation obtained from momentum equation

$$\frac{\partial}{\partial x} \left(\frac{\partial u}{\partial x} \right) = \mu \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial x} \right) - \frac{\partial}{\partial x} \left(\frac{\partial}{\partial x} \left(\frac{\partial u}{\partial x} \right) \right)$$

$$\frac{\partial}{\partial x} \left(\frac{\partial u}{\partial x} \right) = -\frac{\partial}{\partial x} \left(\frac{\partial}{\partial x} \left(\frac{\partial u}{\partial x} \right) \right)$$

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So the solution of this is probably can be done more efficiently for repeated calculation here we may have to do **we may have to** recall quick the corresponding coefficient matrix in this. But, having said that in a this equation will probably be much more difficult to converge as then equation then this equation convection deflection type equation, so one has to really try put and see whether the convection that over head that is required for an implicit method is really worth the additional large time step that is more possible in a implicit method as compared to explicit method. What we like to point out is that explicit method, when we solve in order to go from n to $n + 1$ we solve equation $G1$ when we say equation G poisson equation for known velocity of some known velocity is here. So **this is..** this matrix equation we solved by direct method or indirect method but, we solve it to get p_n once we do that and go to equation d here in the explicitly from like this. And the solution is very straight forward we have three equation and this all are them are explicitly we can the equation is very straight forward so the solution that effort that is needed to go for n th times to $n + 1$ step where from u and p_n will get u_{n+1} and p_{n+1} . We need to solve one Poisson question and three explicit question the conventional effort for this is very negligible compare to this where as in the implicit method we have to solve this as a u_{n+1} equal to a_{n+1} is equal to b_{n+1} and we have to solve p_{n+1} is equal to b we have to solve this several times we have to solve this several times and in the process of solving a does not change from iteration to iteration but, A one will change because we have some non linearity's and coupling terms which appear in this. So this equation will require update the value of them every time so some methods which give us an efficient solutions for repeated solution of a so it cannot be use full for this. So the overall effort that is implement implicit method is much more than in the case of explicit method but, the navier stokes equations are coupled equation in the sense that you want solve the explicit equation we also have to know v and w . And we also have this pressure business that is coming into picture so the overall solution is quit not reveal. So that means that time step may be possible in an explicit method maybe quite small. So this every advantage is .going to an implicit method but, one has to see what kind of compute program skills have to keep the best out of methods for this solution of these type equations and this type of equation. Once we do that an implicit method preferable to an explicit method so this over all pressure equation method is solving the momentum equation like this. And the continental equation is used to deriving a pressure equation, a pressure which will enable the velocity of $n + 1$ step time to be satisfied with the known values of u and p . So that

is a pressure equation method, it can be applied for the time accurate solution of u , v , w at when we say time accurate that knows where we want to get the correct variation of u , v , w with time. So that kind of time accurate things it can be used for three dimension and we are not in working any kind of impossibility effect that comes into the specification that time step limitation here. so the time step limitation is associated with the scalar convection deflection equation that are component from the Courant number it is a component from the deflection parametric which together determine the time step limitation in the case of an explicit method now in a case of implicit method because we are solving this. With constant and specified coefficient there should not be any time step limitation. But, because of u , v , w terms where we need to know the velocity of other component in the solution of x momentum or y momentum they that can be something possibility of instability so we can use large time steps but, not very large time steps so this method is possible and its quite used.

There are other method well discuss one more method which has become almost the standard method for the solution of incompressible equation, especially in problems may effect in process industry that plus inside that is it those kind of things that becomes very very popular in which essentially the same idea is used. In the same idea that is we have a discretized equation and we have continuity equation that pressure does not required, that we reformulate the continuity equation in such a way that correct pressure field is obtained and by correct pressure field we mean that pressure field which was substituted to the momentum equation will enable us to get the correct velocity field or the correct velocity field which will satisfy the continuity equation. So that kind of approach is in this and where we are solving not directly for the pressure. But we are solving for the pressure correction and what we mean by a pressure correction is that with a guess pressure field we can get a velocity field from the momentum equation. So that velocity field need not necessarily satisfies the continuity equation because in solving this with a guess pressure field and for example, an implicit method here we are not enforced continuity we are not used the continuity equation condition at any stage. So having got an idea just estimated whether this pressure field and this pressure field we do not want to correction to the guesses in such a way that we get an improved pressure which will give us more correct solution to which will **which will** satisfy the continuity equation.

So from disturbance we gain between the exact velocity field which satisfies the continuity equation and the estimated velocity field and the momentum equation with a guessed pressure field, so that **destinies** is used to derive a pressure correction which will be used to update or improve the guess pressure field that to get a new value so we will be essentially be solving pressure correction and that correction is incorporated in successive updates of in corporative velocity and must the same way what I am talking about here. And that method is proposed in a late sixty s early seventy s by Pathe and Correns Poltaens and developed in college in London and that has become very widely popular and we will discuss that method and its prove very robust for a longest number of equation and so we look at that method later.