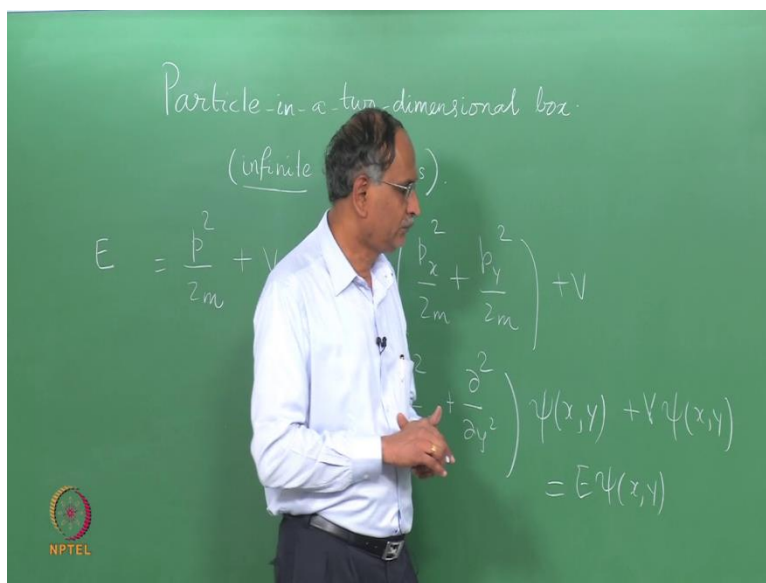


**Introductory Quantum Mechanics and Spectroscopy**  
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**Lecture – 4**  
**Part I**  
**Probabilities-in-a-two-dimensional box (Infinite barriers)**

Welcome back to the lecture. The earlier lecture talked about, in the earlier lecture I talked about the particle in one-dimensional box and in the current one, let us discuss the particle in a two dimensional, two dimensional model or two degrees of freedom model. The particles position coordinates are given by 2 x and y, two coordinates in a plane, orthogonal to each other and then, we discuss the quantum problem.

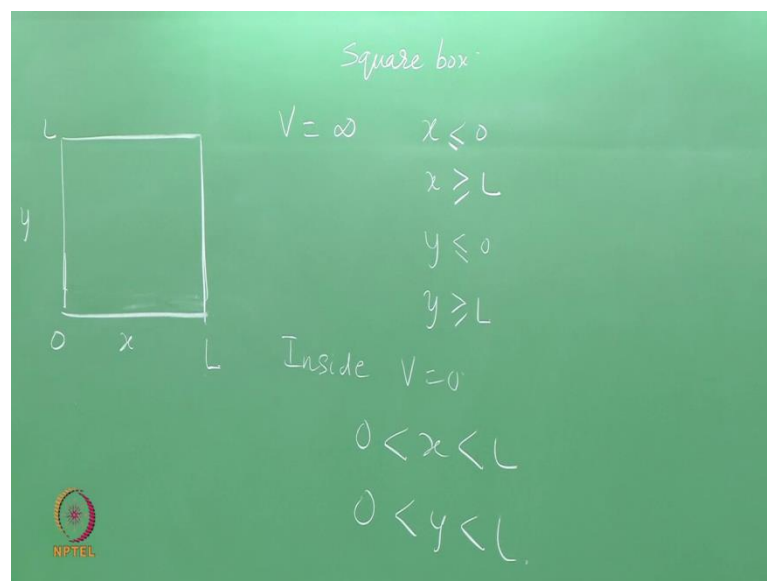
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The barriers are infinite, therefore if you remember the problem  $p^2$  by  $2m$  plus  $V$ , which is the energy term, gets changed to or it is rewritten as  $p_x^2$  by  $2m$  plus  $p_y^2$  by  $2m$  plus  $V$ . And the  $p_x$  is replaced in quantum mechanics by the minus  $\hbar^2$  square by the term minus  $\hbar^2$  square by  $2m$ , the partial derivative now because we have the wave function as a function of two coordinates  $x$  and  $y$  and the momentum in the  $x$  direction is given by the partial derivative. And this is the square of the momentum, so you have minus  $\hbar^2$  square by  $2m$  times  $\frac{\partial^2}{\partial x^2}$  plus  $\frac{\partial^2}{\partial y^2}$  times  $\psi(x,y)$  plus  $V\psi(x,y) = E\psi(x,y)$ .

Correspondingly, for  $p_y^2$ , you have  $-\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial y^2}$ . This is the operator part for the kinetic energy of the Hamiltonian plus and the wave function is a function of  $x$  and  $y$  plus  $V$ , some potential, times  $\psi(x, y)$  is equal to  $E \psi(x, y)$ . This is the two dimensional Schrodinger equation in which you have got the  $\hbar$ , this term plus the  $V$  acting on the  $\psi$  giving you  $E \psi$ . And for the current problem of particle-in-a-two-d box, we consider  $V$  to be infinite for all values of  $x$  other than from  $0$  to  $L$  and all values of  $y$  from  $0$  to some other, say  $a$  or  $L_1$  or  $L_2$ . It does not matter if it is a rectangular box,

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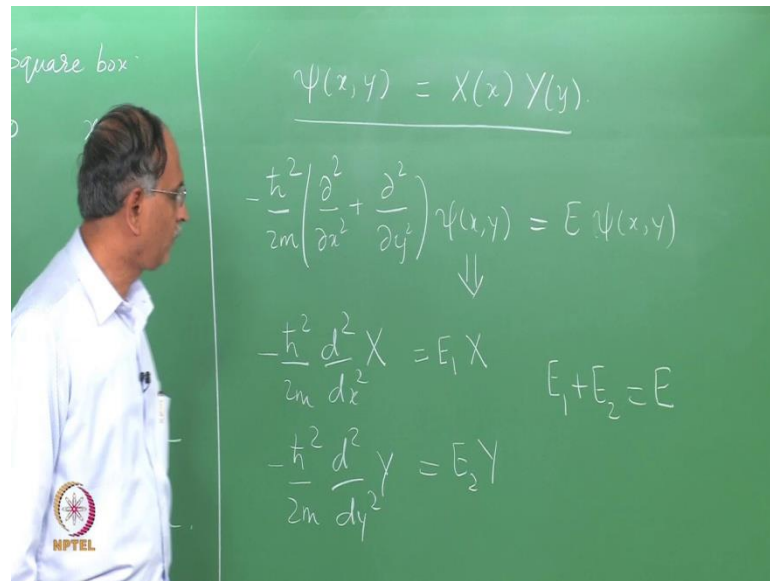


If it is a square box, then essentially you are looking at the, let us see if we can have a square, something like that. So,  $0$  to  $L$  and  $y$ , it is also  $0$  to  $L$ . Only in this region, we are looking at the particle properties and the particles behaviour, and for all others we have  $V$  is infinity for all values of  $x$  less than  $0$  or equal to, and for all values of  $x$  greater than or equal to  $L$ . Likewise, for  $y$  less than or equal to  $0$ ,  $y$  greater than or equal to  $L$ . So, this is the infinite boundaries that you have.

It is not the single dimensional quantity, but it is a surface in a sense, that we protect the particle from escaping this region. And inside  $V$  is  $0$  between  $x$  and  $L$ , between  $y$  and  $L$  and this is a square box.

So, if we do that, obviously, the differential equation simplifies without this term, and you have a derivative square in one direction, a derivative square in another direction and then you have the  $\psi$  of  $x, y$ .

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Such a problem is easily solved by, is written in terms of a product of a function of  $x$  alone and a function of  $y$  alone. With this choice, it is possible to separate this equation minus  $\hbar$  square by  $2m$  dou square by dou  $x$  square plus dou square by dou  $y$  square  $\psi$  of  $(x, y)$  is equal to  $E$  times  $\psi$  of  $(x, y)$  into two equations, namely, minus  $\hbar$  square by  $2m$  d square by d  $x$  square  $X$  is equal to  $E_1$  of  $X$  and minus  $\hbar$  square by  $2m$  d square by d  $y$  square times  $Y$  is equal to  $E_2$  times  $Y$ . But these two constants  $E_1$  and  $E_2$  are constrained by  $E_1 + E_2 = E$ .

The actual separation of this is given in the notes that accompanies this video lecture. Therefore, I would request you to look into that to see how this equation is separated into two one-dimensional equation, one for  $X$  and one for  $Y$  with the constraint, that the energies for the two one-dimensional problems are related to the total energy, has this sum  $E_1 + E_2$ .

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independently of the other. Therefore, the above equation is satisfied only if each term is separately equal to a constant.

$$-\frac{1}{X(x)} \frac{\hbar^2}{2m} \frac{d^2 X(x)}{dx^2} = E_1, \text{ or, } -\frac{\hbar^2}{2m} \frac{d^2 X(x)}{dx^2} = E_1 X(x) \text{ and}$$
$$-\frac{1}{Y(y)} \frac{\hbar^2}{2m} \frac{d^2 Y(y)}{dy^2} = E_2, \text{ or, } -\frac{\hbar^2}{2m} \frac{d^2 Y(y)}{dy^2} = E_2 Y(y)$$


Now, let us see the solutions that quantity, which I have written on the board is namely, this is the x equation and the corresponding y equation is that. Obviously, each one of them is like a one-dimensional, part, particle in a box. Therefore, the solutions for each one of them will have a running quantum number for that particular equation. The x component of the wave function will be given by the solution. It is similar to the psi of x that we wrote except, that now we call it X of x and now, this will have a quantum number going from 1, 2, 3 to some value, which we call as n 1. In an exactly, in an identical manner, the y equation will also have a free quantum number n 2, which will run from 1, 2, 3 to whatever that we take.

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$\psi(x, y)$

$E_1 = \frac{h^2}{8mL^2} n_1^2$

$E_2 = \frac{h^2}{8mL^2} n_2^2$

$E = \frac{h^2}{8mL^2} (n_1^2 + n_2^2)$

$E_1 + E_2 = E$

Separation of variables

What are the new results for the particle in the two dimensional (coordinate) square box whose side is of length L. Then we have all the solutions same one dimensional box except that they are repeated for one coordinate e dimensional problem. An important new result in addition is the requi...

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But please remember, these two quantum numbers are not independent, in the sense, they are connected to the total energy, the requirement that  $E_1 + E_2 = E$ . Now, remember the expression for  $E_1$  from particularly one-dimensional box it is  $\frac{h^2}{8mL^2} n_1^2$ , a free quantum number. In the sense, it takes 1, 2, 3 integer values. And  $E_2$  is also given by  $\frac{h^2}{8mL^2} n_2^2$  such that this equation is satisfied therefore, you have  $\frac{h^2}{8mL^2} n_1^2 + \frac{h^2}{8mL^2} n_2^2 = E$ .

So, this is the only constraint that comes out in the separation of the two-dimensional Schrodinger equation that the total energy is the sum of the two one-dimensional energies and that is possible because we do not have a potential, which couples the two dimensions. We have put  $V = 0$ , therefore the method of separation of variables. We have separated the  $x$  and  $y$  from the  $\psi$  of  $(x, y)$ . If you recall, the  $\psi$  of  $(x, y)$ , we have separated that into the  $x$  equation and the  $y$  equation, so that process is called the separation of variables. Now, how do these functions look like?

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dimensional problem. An important new result in addition is the require-  
energies have to add up to a total energy, which will lead to the idea of  
states, to be discussed below.

$$X_{n_1}(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n_1 \pi x}{L}\right), E_{n_1} = \frac{h^2 n_1^2}{8mL^2}$$

$$Y_{n_2}(y) = \sqrt{\frac{2}{L}} \sin\left(\frac{n_2 \pi y}{L}\right), E_{n_2} = \frac{h^2 n_2^2}{8mL^2}$$

$$E_{n_1} + E_{n_2} = \frac{h^2}{8mL^2} (n_1^2 + n_2^2) = E_{n_1 n_2}$$

Obviously, you have the solutions for the quantum number n 1. In terms of the one-dimensional solution, that you have seen in the previous lecture, root 2 by L sine n 1 pi x by L and the energy is given by n 1 square. And likewise, for the y with the n 2 square, and with the constraint, that the total energy E n 1 plus E n 2 is E n 1 n 2, you have seen that.

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$$Y_{n_2}(y) = \sqrt{\frac{2}{L}} \sin\left(\frac{n_2 \pi y}{L}\right), E_{n_2} = \frac{h^2 n_2^2}{8mL^2}$$

$$E_{n_1} + E_{n_2} = \frac{h^2}{8mL^2} (n_1^2 + n_2^2) = E_{n_1 n_2}$$

$$\Psi_{n_1 n_2}(x, y) = X_{n_1}(x) Y_{n_2}(y)$$

Examples of quantum numbers, wave functions and energies:

$n_1 = 1, n_2 = 1$	$\psi_{11}(x, y) = \frac{2}{L} \sin\left(\frac{\pi x}{L}\right) \sin\left(\frac{\pi y}{L}\right)$	$E_{11} = \frac{h^2}{8mL^2} (1^2 + 1^2)$
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What about the wave function? The wave function, now if you see this, the wave function psi of n 1 n 2, because it is obviously specified by the two quantum numbers n 1

and  $n_2$ , has the independent function  $x$  with the quantum number  $n_1$  and  $y$  with the quantum number  $n_2$ . Each one is in an orthogonal direction. Therefore, you see this interesting thing next line.

When we have  $n_1$  is 1, and  $n_2$  is 1, when we have that case, which is the starting point, what is called the lowest energy for the particle in a two-dimensional box. You can see, that the wave function is given by  $\psi_{11}(x, y)$  and it is given by the product of the two functions, that you saw the  $x$  of  $x$  and  $y$  of  $y$ , which gives you  $\sin \frac{\pi x}{L}$  and  $\sin \frac{\pi y}{L}$  by  $L$ .

Let me repeat this. When the quantum number is 1 1, the wave function is given by  $\psi_{11}$  and it is given by the product of  $\sin \frac{\pi x}{L}$  and  $\sin \frac{\pi y}{L}$ . And the energy is, of course, the sum of 1 square plus 1 square times the whole thing.

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The screenshot shows a presentation slide with the following content:

- Equation:  $\left. \frac{n_2 \pi y}{L} \right\}, E_{n_2} = \frac{h^2 n_2^2}{8mL^2}$
- Equation:  $n_1^2 + n_2^2 = E_{n_1 n_2}$
- Equation:  $\left. \right\} Y_{n_2}(y)$
- Text: "um numbers, wave functions and energies:"
- Equation:  $\psi_{n_1 n_2}(x, y) = \frac{2}{L} \sin\left(\frac{\pi x}{L}\right) \sin\left(\frac{\pi y}{L}\right)$
- Equation:  $E_{11} = \frac{h^2}{8mL^2} (1^2 + 1^2) = \frac{h^2}{4mL^2}$
- Handwritten notes:
  - $\psi_{n_1 n_2} = X_{n_1} Y_{n_2}$  with  $n_1 \neq n_2$
  - $= X_{n_2} Y_{n_1}$  with  $E \rightarrow (n_1^2 + n_2^2) \frac{h^2}{8mL^2}$
  - "Degenerate state (2)"

Therefore, the energy for this process  $E_{11}$  is  $\frac{h^2}{8mL^2}$  times 2. What is interesting is the next choice. You have  $\psi_{n_1 n_2}$  as  $X$  of  $n_1$   $Y$  of  $n_2$ . It is possible, if  $n_1$  is not equal to  $n_2$ . It is possible to have the wave function given by  $X$  of  $n_2$  and  $Y$  of  $n_1$ , because the energy is simply proportional to  $n_1$  square plus  $n_2$  square times  $\frac{h^2}{8mL^2}$ , which is the proportionality constant.

Therefore, you see, that you have the same energy, but you have two physically different states  $X$  of  $n_1$   $Y$  of  $n_2$ , and  $X$  of  $n_2$   $Y$  of  $n_1$ , both states have the same energy. This is

what is called a degenerate state. Degeneracy is 2 because there are two states, which have the same energy, but have different quantum states.

This is the introduction for the particle in a two-d box, that the degeneracy is the additional factor. Now, how do these things look like, let us simplify this picture.

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$$E_{n_1, n_2} = E_{n_1} + E_{n_2} = \frac{h^2}{8mL^2}(n_1^2 + n_2^2) = E_{n_2, n_1}$$

$$\Psi_{n_1, n_2}(x, y) = X_{n_1}(x)Y_{n_2}(y)$$

Examples of quantum numbers, wave functions and energies:

$n_1 = 1, n_2 = 1$	$\Psi_{1,1}(x, y) = \frac{2}{L} \sin\left(\frac{\pi x}{L}\right) \sin\left(\frac{\pi y}{L}\right)$	$E_{11} = \frac{h^2}{8mL^2}(1^2 + 1^2) = \frac{h^2}{4mL^2}$
$n_1 = 2, n_2 = 1$	$\Psi_{2,1}(x, y) = \frac{2}{L} \sin\left(\frac{2\pi x}{L}\right) \sin\left(\frac{\pi y}{L}\right)$	$E_{21} = \frac{h^2}{8mL^2}(2^2 + 1^2) = \frac{5h^2}{8mL^2}$
$n_1 = 1, n_2 = 2$	$\Psi_{1,2}(x, y) = \frac{2}{L} \sin\left(\frac{\pi x}{L}\right) \sin\left(\frac{2\pi y}{L}\right)$	$E_{12} = \frac{h^2}{8mL^2}(1^2 + 2^2) = \frac{5h^2}{8mL^2}$
$n_1 = 2, n_2 = 2$	$\Psi_{2,2}(x, y) = \frac{2}{L} \sin\left(\frac{2\pi x}{L}\right) \sin\left(\frac{2\pi y}{L}\right)$	$E_{22} = \frac{h^2}{8mL^2}(2^2 + 2^2) = \frac{h^2}{mL^2}$
$n_1 = 3, n_2 = 2$	$\Psi_{3,2}(x, y) = \frac{2}{L} \sin\left(\frac{3\pi x}{L}\right) \sin\left(\frac{2\pi y}{L}\right)$	$E_{32} = \frac{h^2}{8mL^2}(3^2 + 2^2) = \frac{13h^2}{8mL^2}$

Handwritten note:  $\Psi_{n_2, n_1} \rightarrow E \rightarrow \frac{h^2}{8mL^2}(n_1^2 + n_2^2)$   
Degenerate state (2)

Now, I have a whole series of function here with which you can fill up any number of pages if you wish. You see, that  $n_1 = 2, n_2 = 1$  corresponds to the wave function  $\Psi_{2,1}$  with  $\sin(2\pi x/L) \sin(\pi y/L)$ . And  $n_1 = 1, n_2 = 2$  gives you the other function, namely,  $\sin(\pi x/L) \sin(2\pi y/L)$  and the energies are the same.

So, if the quantum numbers are identical, there is no degeneracy. But if the quantum numbers are different for a square box, because we have chosen the length  $L$  to be the same, the square box gives you the solution, that you have a minimum degeneracy of 2 if  $n_1$  is not the same as  $n_2$ .

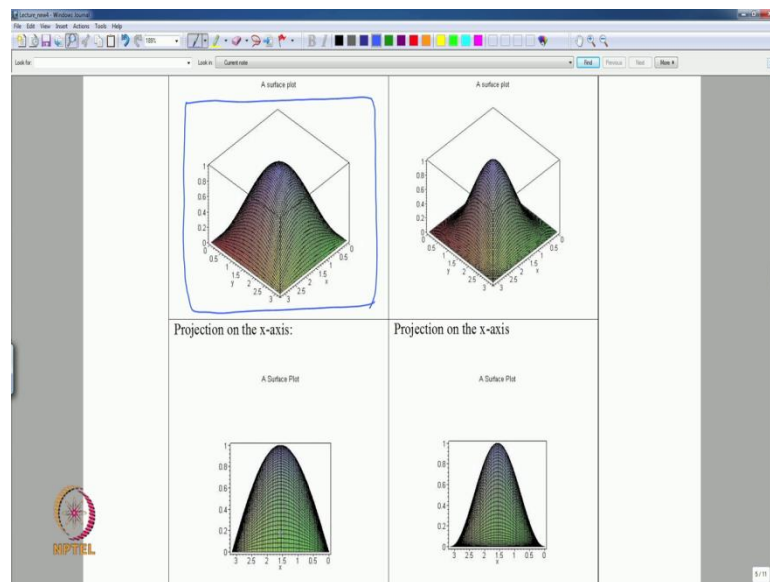


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Thus,

And you can see, that for 3 and 2, that you have the wave function sine 3 pi x by L and sine 2 pi y by L. And then, 2 and 3, which is sine 2 pi x by L sine 3 pi y by L. So, the axis choice, the quantum number choice for a given axis determines the function's state.

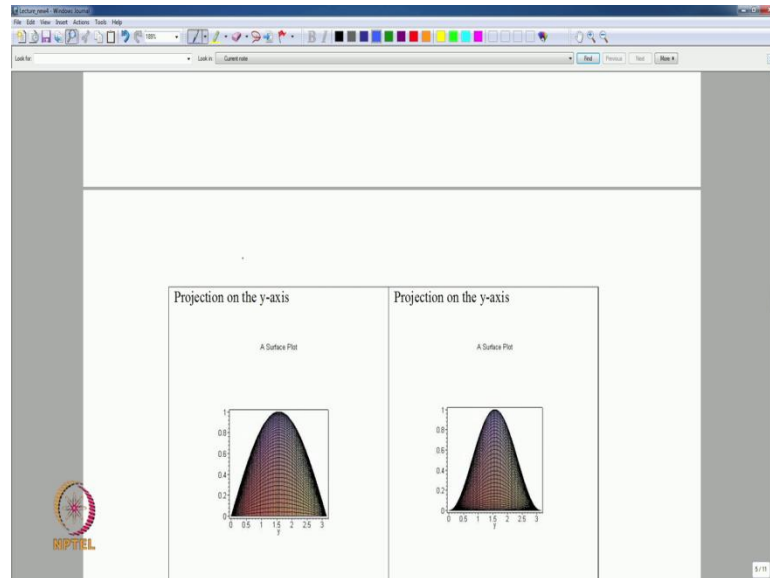
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How do these things look like if we plot them? I mean, this plot looks fancy; actually, it does not have much interpretation or meaning, but it is worth seeing the product wave function in two dimensions.

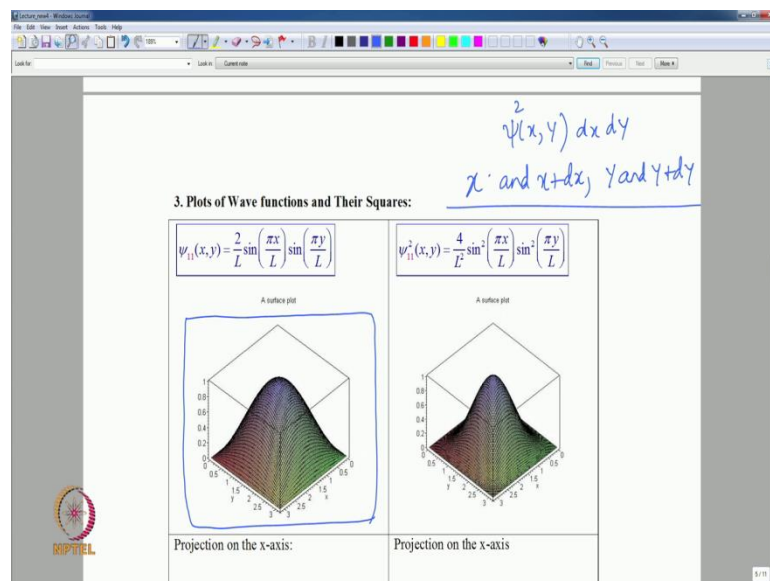
So, you see the wave function, you see the wave function  $\psi_{11}$  using this picture. It is a half wave similar to what you had in your particle-in-a-one-dimensional box in the  $x$  direction, and it is also a half wave in the  $y$  direction, as you can see through the projection in the  $x$  direction here of this graph.

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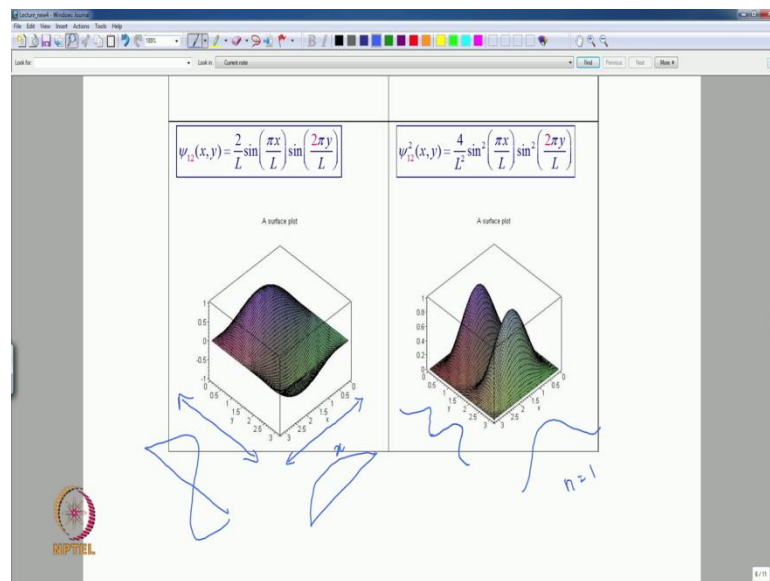
And on the  $y$  direction also you have the same thing identical. What about the  $\psi$  square?

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The psi square, which is associated with the probability, that the particle be found not in a small length region dx, but in the small area dx dy. Please remember, psi x y, if you do that psi square dx dy is the probability, that the particle will be in the small rectangular region between x and x plus d x and y and y plus d y, that is a small region. And you can see, that the psi square is given like this. Therefore, you can create, I mean, you can visualize what would be the probability exactly the same way that you have visualized the particle in a one-dimensional box except, that now we have a motion on the plane.

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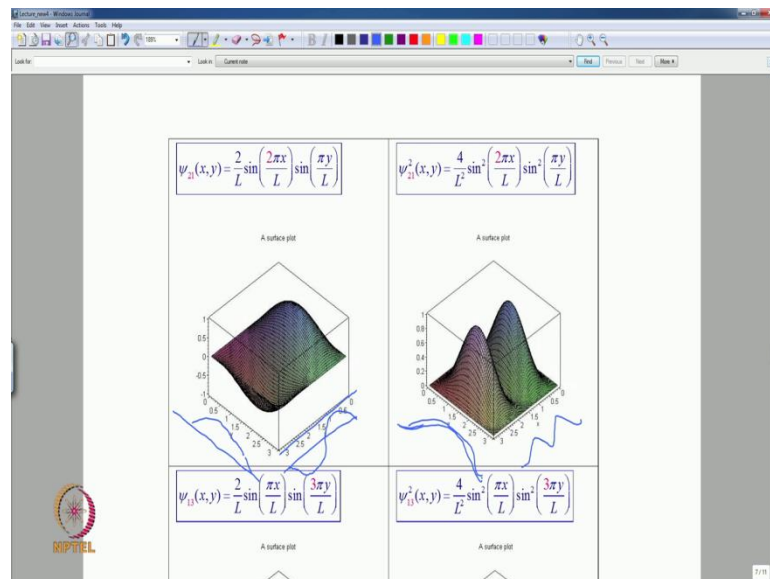
And now, what is interesting is when you go to different quantum numbers where there is degeneracy psi 1 2. If you look at this, psi 1 2 is quantum number 1 for the x direction, and quantum number 2 for the y direction. Therefore, this is the quantum number, this is the quantum number for the x direction and you can see, that it is a half wave, which is either up or down. It is either positive or negative, the reason being, the y direction wave is a full wave. So, in this direction, what you have is if I may draw this, the wave function looks like that. In this direction, the wave function looks like that.

Therefore, when you take the product of these two functions, a negative side makes this wave function negative for half the length and therefore, you see, that for half the length you have either a positive wave function or you have a negative wave function, that is, only for the wave function. We know, that the wave function is not that important, it is a square of the wave function, which is important for probability interpretation. And you

can see, that psi square, which removes this negative character of the function gives you now very beautifully the 2 n equal to 1 case for the x-axis. And the n equal to 2, if you remember the graph that you had for n equal to 2 for the y-axis and this is the x-axis.

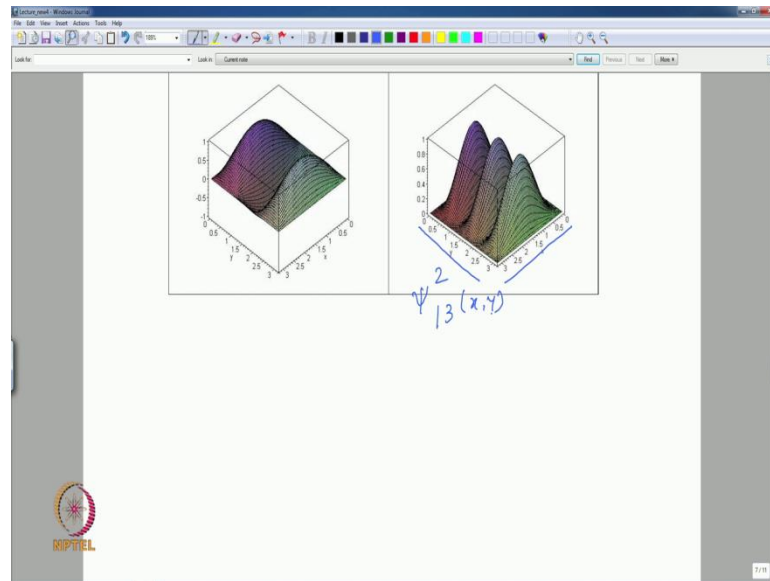
Therefore, the features are captured, the wave function features are captured when you do the surface plot and you can see, that the pictures can be created for a large number of them, but there is a limit, two dimensions. And in three dimension, we probably can use colour at the most to distinguish the function from the three axis, but that is it. You cannot visualize this for n dimensions.

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So, let us conclude this part of the particle in a two-dimensional box with some examples of the wave functions and the squares of the wave function for different quantum numbers. So, here is a 2 1 as opposed to 1 2. And you see all that happens is, that for a 2 1, the wave function along the x-axis is like this and the wave function along the y-axis, it is like that. And you can see, that actually, sorry, this is in the wrong direction, so let me erase that because your 0 starts from here, therefore have that. And this is the y-axis; that is the reason why part of it is negative and the other part is positive. The square of the wave function you can see, that there are two humps along the x-axis and along the y-axis. It is a quantum number one, so you have only one similar to the one-dimensional y-axis. And let us see one or two more examples and let me stop with that.

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This is, I mean, the exercise here, what does this picture represent? There is one here along the x-axis, and there are three peaks. Therefore, you have, this is a, y is 3 and x is 1, so it is psi 1 3 square (x, y).

So, the lecture notes give you many more such pictures, but in the next part of this lecture we will see what do all these things mean in terms of probability calculations and in terms of a new idea called the expectation values. We will stop here for this particular part of the lecture.

Thank you.