

Introductory Quantum Mechanics and Spectroscopy
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Lecture - 5
Part I
The Quantum Mechanics of Hydrogen Atom

Welcome back to the lectures on CY1001 or introductory chemistry. In this group of lectures consisting of several parts, I shall describe the quantum mechanics associated with the hydrogen atom. The solution of the Schrodinger equation, I will give you the results.

The solution of the Schrodinger equation has its first major achievement in arriving at the spectra of the hydrogen atom, which were known for many decades before that and it is a spectra of the hydrogen atom, which prompted Niels Bohr who came up with his first model of quantizing the energy and quantizing the angular momentum of an atom.

Schrodinger equation, of course, does this using his prescription and the wave function. And we shall see some of the details in the calculation of the energies and in the calculation of probabilities of the electrons, and so on.

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The quantum mechanics of
Hydrogen atom.

Several parts. PART I

$$\hat{H}\psi = E\psi \quad \psi(x,y,z)$$

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So, this is part one of the quantum mechanics of hydrogen atom and in the solution of this equation, $H\psi$ is equal to $E\psi$, where now ψ is three-dimensional Cartesian coordinate.

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The slide contains the following content:

Diagram: A nucleus with charge $+Ze$ and an electron with charge e^- at a distance r .

Equation: $\hat{H}\psi = E\psi \quad \psi(x, y, z)$

Reduced mass calculation:

$$\mu = \frac{m_e m_p}{m_e + m_p}$$

Approximation: $\mu \approx \frac{m_e m_p}{m_p} \approx m_e$ (since $m_p \gg m_e$)

Final result: $\mu \approx m_e$

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We do use a classical starting point of the nucleus with an electron somewhere and the nucleus having a positive charge, plus $z e$, z is 1 and the electron with a minus charge and a distance of r . I shall not describe this as a two-body problem even though that is the right way of doing it. The two-body problem and then, remove the centre mass from the two-body problem and study only the relative motion of the two particle systems, which in this case, the relative mass for the or the reduced mass for the two particle system is the mass of the electron times the mass of the proton, divided by the mass of the electron plus that of the proton, which is approximately the mass of the electron divided by the times mass of the proton divided by the mass of the proton.

Since m_p is much much greater than m_e and therefore, μ turns out to be approximately m_e when you cancel the m_p .

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The image shows a digital whiteboard with handwritten text and equations. At the top, it says $\mu \approx m_e$. Below that, m_e is written. The main text reads "Hamiltonian for electron." followed by the equation
$$\hat{H} = \frac{1}{2m_e} \hat{p}^2 - \frac{Ze^2}{4\pi\epsilon_0 r}$$
 An arrow points from the \hat{p}^2 term to the expression $p_x^2 + p_y^2 + p_z^2$ written below it. The whiteboard has a standard toolbar at the top and an NPTEL logo in the bottom left corner.

So, we shall worry about making that approximation and write the mass as nothing but the mass of the electron. Therefore, we needed the Hamiltonian for the electron.

So, let us assume, that the nucleus is stationary, does not contribute to the overall kinetic energy of the atom that is already there in the centre of mass, which is not considered here. Therefore, if you write to the kinetic energy and the potential energy operator for the hydrogen atom, it will be in terms of the operators. It will be p square 1 by $2m_e$ minus $z e$ square by $4 \pi \epsilon_0 r$, which is the classical coulombic energy of interaction between the positive and the negative charge. This p square, which is an operator, is given by p_x square plus p_y square plus p_z square, the three components of the momentum in a coordinate system, which is probably fixed in the nucleus itself for arguments.

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The image shows a handwritten derivation of the Hamiltonian operator \hat{H} for a hydrogen atom. The equation is written as:

$$\hat{H} = -\frac{\hbar^2}{2m_e} \left(\frac{\partial^2}{\partial x_e^2} + \frac{\partial^2}{\partial y_e^2} + \frac{\partial^2}{\partial z_e^2} \right) - \frac{ze^2}{4\pi\epsilon_0 r}$$

Annotations above the equation include $2m_e$ with an arrow pointing to the denominator of the kinetic energy term, and $4\pi\epsilon_0 r$ with an arrow pointing to the denominator of the potential energy term. Above the kinetic energy term, the momentum components are written as $p_x^2 + p_y^2 + p_z^2$. Below the equation, the boundary conditions for the coordinates are given:

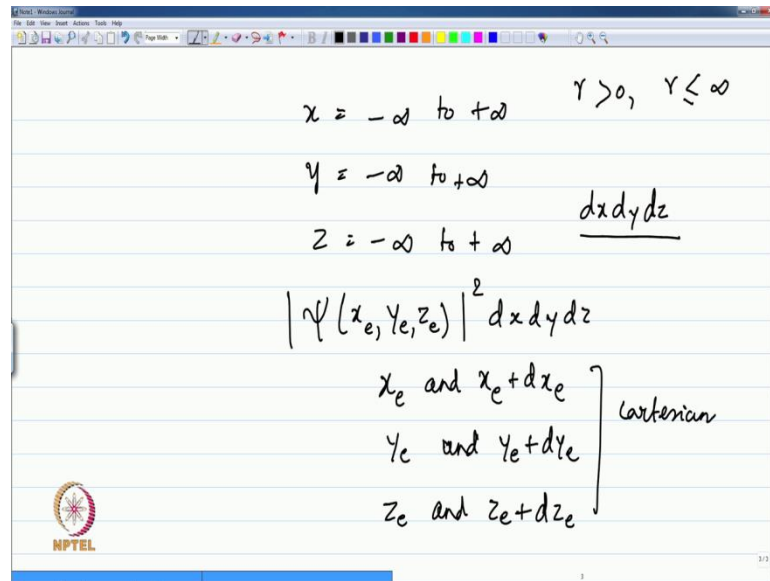
$$x = -\infty \text{ to } +\infty \quad y > 0, \quad y \leq \infty$$
$$y = -\infty \text{ to } +\infty$$
$$z = -\infty \text{ to } +\infty$$

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And then, of course, p_x is replaced by the derivative operators, so that the Hamiltonian becomes minus \hbar squared by $2m_e$ squared by ∂x_e squared plus ∂y_e squared plus ∂z_e squared, where these are the coordinates of the electron with respect to that origin.

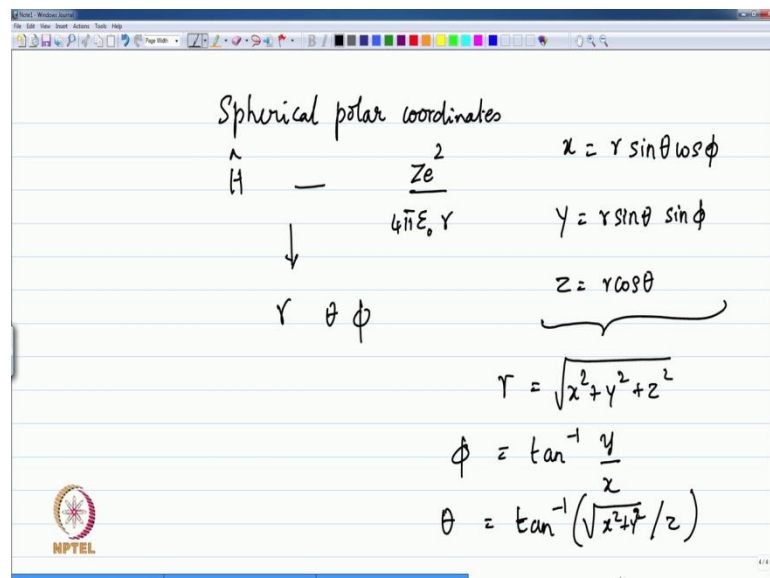
And then, you have the potential energy minus ze^2 by $4\pi\epsilon_0 r$. And r is in principle greater than 0 and less than or equal to infinity. At infinity, of course, the coulombic interaction is 0, therefore here the boundary includes the entire three-dimensional world, the whole universe. So, the boundaries are explicitly, x is from minus infinity to plus infinity, y is from minus infinity to plus infinity and also, z is from minus infinity to plus infinity.

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So, this is the three dimensional region and the volume element, that we talk about for the particle for the electron probability is psi, is the volume element is dx, dy, d z. And then the probability is psi x e, y e, z e absolute square dx dy dz as the probability of finding the electron in the region or in the cube between x e and x e plus dx e, y e and y e plus dy e, and z e and z e plus dz e. This is the three-dimensional Cartesian coordinate representation for the hydrogen electron problem, the nucleolus electron problem.

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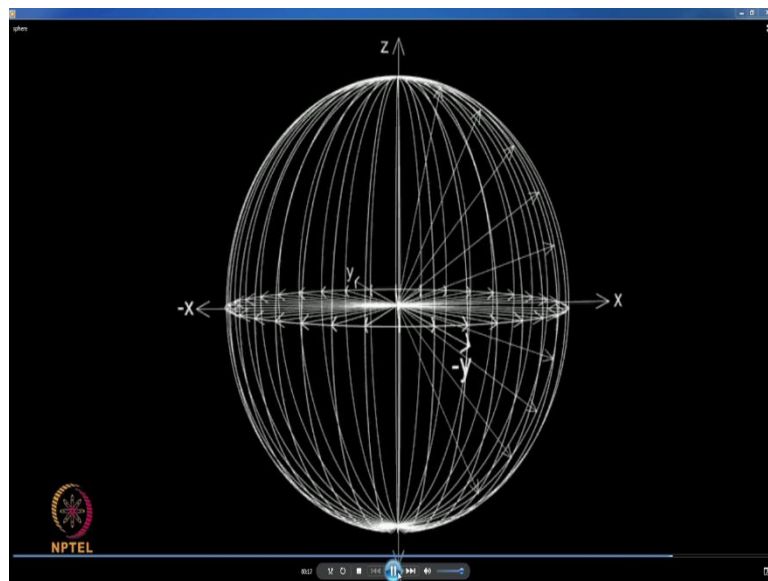


Instead of Cartesian coordinates, in the case of hydrogen atom one uses spherical polar coordinates. The reason for that is, that if you look at the hydrogen atom, the potential energy is spherically symmetric. Therefore, the important contribution to the stability of the hydrogen atom, which is the binding energy between that coulombic charges being spherically symmetrical, the system is better described using the spherical polar coordinates, which if you recall have three variables, the radius of the sphere and then the polar angles theta and phi on the sphere. The standard relations for these are, x is equal to r sine theta cos t, and y is r sine theta sine phi and z is r cos theta.

These are the equations for the transformation between polar and the Cartesian coordinates and the inverse transformation is of course, r is square root of x square plus y square plus z square. And if you take the ratio of x by y r sin theta cancels off, you have phi or phi is tan inverse y by x and the last relation is theta, which is given in terms of tan inverse square root of x square plus y square divided by z. So, the coordinate transformation allows you to either use the spherical or the Cartesian coordinate by using the relationship between them.

And this one animation gives you the relation or the visualization of the spherical polar coordinates system and the values. Let me play the animation.

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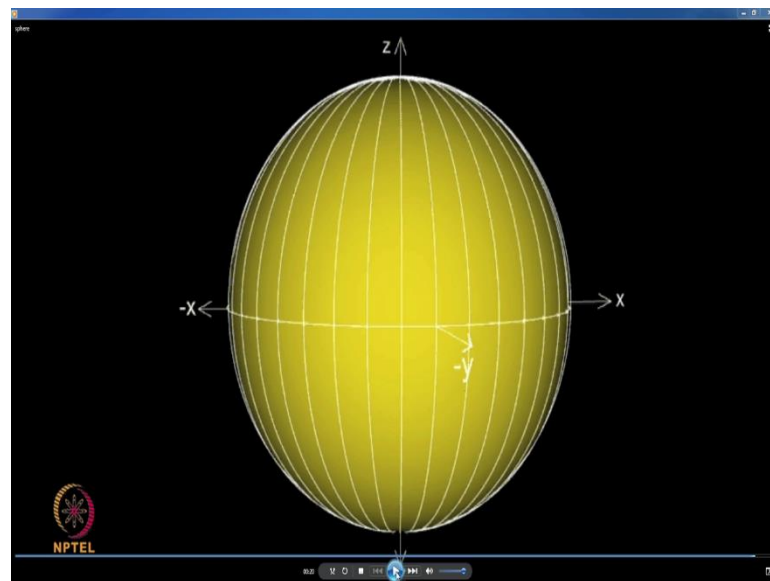


Here, the relation between the Cartesian and the spherical system is given for a, given one value of r the radius of a sphere and you can see, that if you fix a polar axis called

the z axis, then the polar angle theta is the angle theta varying from 0 to pi, as shown by these different radii. So, that is a variation of theta and theta varies from 0 to pi only.

And the other angle is, of course, the Azimuthal angle phi, which is perpendicular in a plane perpendicular to this and if you rotate this arc, semi arc by 2 pi, you generate the surface of the sphere. So, that is it, the Azimuthal angle phi with respect to a chosen x axis.

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So, that is a spherical coordinate system, in which you can see the variation in theta given by these different arcs and the value of phi corresponding to each one of these arcs starting from the x-axis here at some arbitrary point and then going around the x-axis to the plus y to the minus x to the minus y and back you have 2 pi.

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$$\begin{array}{l} r \rightarrow 0 \text{ to } \infty \\ \theta \rightarrow 0 \text{ to } \pi \\ \phi \rightarrow 0 \text{ to } 2\pi \end{array} \left\| \begin{array}{l} x \\ y \\ z \end{array} \right\} \begin{array}{l} -\infty \text{ to } +\infty \\ -\infty \text{ to } +\infty \\ -\infty \text{ to } +\infty \end{array}$$
$$dx dy dz \rightarrow r^2 dr \sin \theta d\theta d\phi$$

Therefore, if we recall our lecture component, r varies from 0 to infinity, being the radius of the sphere. The sphere is from 0 radius to all over the universe. And theta varies from 0 to pi as the, as the polar angle varying from 0 to pi, as you have seen with respect to the z axis. And the phi, which goes around the circle in 2π ; phi is 0 to 2π . And these are relations in parallel to the x, y, z , all going from minus infinity to plus infinity in the Cartesian axis in taking care of the whole universal space.

Therefore, this, these are the limits and dx, dy, dz , which is a volume element in Cartesian coordinate space, will have to be expressed in terms of the volume elements in spherical polar coordinate system and that is given by $r^2 dr \sin \theta d\theta d\phi$.

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$\phi \rightarrow 0 \text{ to } 2\pi$

$dx dy dz \rightarrow r^2 dr \sin \theta d\theta d\phi$

$dx dy dz = \begin{vmatrix} \frac{\partial x}{\partial r} & \frac{\partial x}{\partial \theta} & \frac{\partial x}{\partial \phi} \\ \frac{\partial y}{\partial r} & \frac{\partial y}{\partial \theta} & \frac{\partial y}{\partial \phi} \\ \frac{\partial z}{\partial r} & \frac{\partial z}{\partial \theta} & \frac{\partial z}{\partial \phi} \end{vmatrix} dr d\theta d\phi$

← Jacobian

Those of you who are not familiar with this transformation must go back and look at the coordinate transformation and simple differentials expressed from one coordinate to the other and the relations are given by, what is known as, the Jacobian, the magnitude of the Jacobian.

The Jacobian being the partial derivative of x with respect to r, x with respect to theta, and x with respect to phi, and the partial derivative of y with respect to r with respect to theta and with respect to phi and likewise, the partial derivative with respect to z of z with respect to r and with respect to theta with respect to phi. The determinant of this multiplied by dr d theta d phi is this. This is called the Jacobian and this is in elementary transformation matrix, that transforms volume elements from one coordinate system, another coordinate system. And this Jacobian with the magnitude has the r square sine theta with the dr d theta d phi.

Therefore, when you calculate the volume elements and when you calculate the probabilities using spherical polar coordinate system, if you are using Cartesian coordinates, transform the relation from Cartesian to the polar. And these are the mathematical formulas already well known and derived from elementary differential calculus. Keep this in mind.

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$$\hat{H} = -\frac{\hbar^2}{2m_e} \left[\frac{\partial^2}{\partial x_e^2} + \frac{\partial^2}{\partial y_e^2} + \frac{\partial^2}{\partial z_e^2} \right] - \frac{Ze^2}{4\pi\epsilon_0 r}$$

↓
spherical polar coordinates

$$\frac{\partial \psi(x, y, z)}{\partial x} =$$

$$\psi(x, y, z) \Rightarrow \bar{\psi}(r, \theta, \phi)$$

Therefore, now we have this Hamiltonian expressed in terms of minus h bar square by 2 m e dou square by dou x e square plus dou square by dou y e square plus dou square by dou z e square minus Z e square by 4 pi epsilon naught r. This needs to be changed to polar coordinates, spherical polar coordinates.

That is not a trivial exercise, but it is not a very hard exercise. The derivatives, for example, dou by dou x of any function of x, y, z are expressed in another coordinate system like r theta phi, if you have to express psi in terms of r theta phi. The derivatives are expressed using the partial derivatives of the coordinates with respect to the new coordinates. So, for example, dou by dou x of psi if you want to write, the appropriate wave function in the polar coordinate, namely psi x, y, z is replaced by the corresponding substitution of the x and y and z using r theta phi using this function.

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The image shows a presentation slide with a white background and a blue border. At the top right, the text $\psi(x, y, z) \Rightarrow \Psi(r, \theta, \phi)$ is written. Below this, the chain rule is written as:

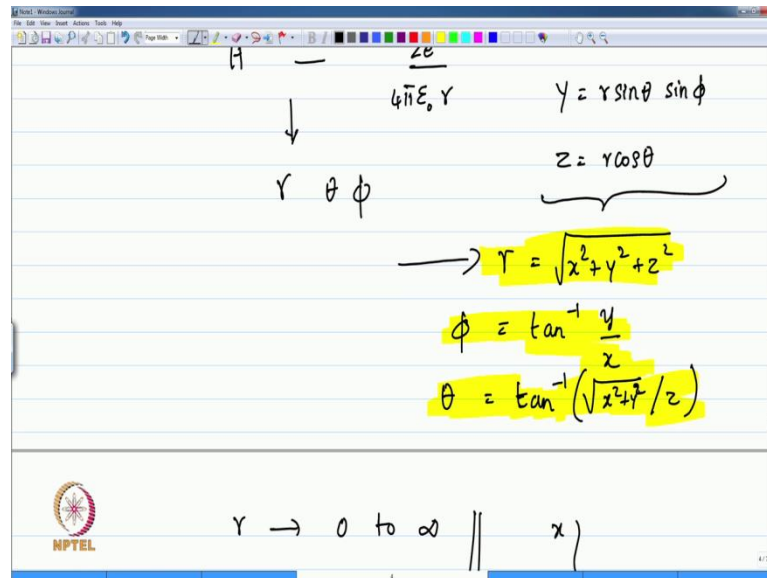
$$\frac{\partial \psi}{\partial x} = \left[\frac{\partial r}{\partial x} \frac{\partial}{\partial r} + \frac{\partial \theta}{\partial x} \frac{\partial}{\partial \theta} + \frac{\partial \phi}{\partial x} \frac{\partial}{\partial \phi} \right] \Psi(r, \theta, \phi)$$

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Then, there is a very simple partial derivative chain rule, which tells you how to calculate $\frac{\partial \psi}{\partial x}$ as nothing other than $\frac{\partial \psi}{\partial r} \frac{\partial r}{\partial x}$ plus $\frac{\partial \psi}{\partial \theta} \frac{\partial \theta}{\partial x}$ plus $\frac{\partial \psi}{\partial \phi} \frac{\partial \phi}{\partial x}$ acting on the wave function $\psi(r, \theta, \phi)$.

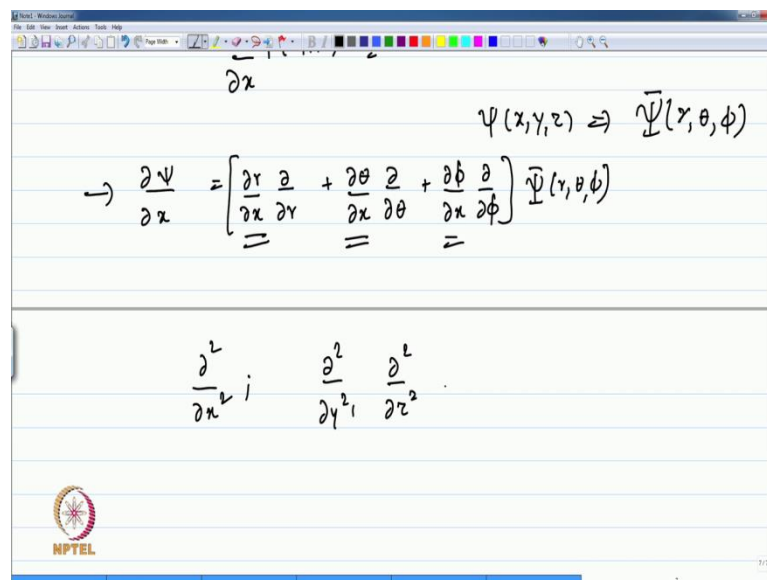
So, this is the transformation of the derivative form of the Cartesian coordinate into the corresponding polar coordinates. Of course, you can calculate $\frac{\partial r}{\partial x}$, $\frac{\partial \theta}{\partial x}$, and $\frac{\partial \phi}{\partial x}$ from the inverse relations, that you already have, already have at.

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From this you can calculate the derivative of r with respect to x , y and z ; the derivative of ϕ with respect to x , y and z ; and the derivative of θ with respect to x , y and z .

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Therefore, the partial derivatives that you need to calculate for expressing the kinetic energy in spherical polar coordinate system involves three such quantities, namely $\frac{\partial^2}{\partial x^2}$, $\frac{\partial^2}{\partial y^2}$, and $\frac{\partial^2}{\partial z^2}$, which is operating this once more, but being careful, that the terms contain already r , θ , and ϕ . And therefore, the partial derivatives have to be taken carefully.

And you have to do the same thing for double square by double y square and double square by double z square.

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$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

$$\nabla^2 = \frac{\partial^2}{\partial r^2} + \frac{\partial^2}{\partial \theta^2} + \frac{\partial^2}{\partial \phi^2}$$

$$\frac{\partial \psi}{\partial y} = \left[\frac{\partial r}{\partial y} \frac{\partial}{\partial r} + \frac{\partial \theta}{\partial y} \frac{\partial}{\partial \theta} + \frac{\partial \phi}{\partial y} \frac{\partial}{\partial \phi} \right] \psi(r, \theta, \phi)$$

$$\frac{\partial \psi}{\partial z} = \left[\frac{\partial r}{\partial z} \frac{\partial}{\partial r} + \frac{\partial \theta}{\partial z} \frac{\partial}{\partial \theta} + \frac{\partial \phi}{\partial z} \frac{\partial}{\partial \phi} \right] \psi(r, \theta, \phi)$$

Therefore, let me summarize this particular part of the lecture with the corresponding expressions, namely, double psi by double y as double r by double y double by double r plus double theta by double y double by double theta plus double phi by double y double by double phi acting on the wave function psi of r theta and phi.

And similarly, double psi by double z as double r by double z double by double r plus double theta by double z double by double theta plus double phi by double z double by double phi acting on the corresponding wave function psi r theta phi.

These are the derivative equivalents and you calculate likewise the double square terms, the double square by double y square terms and the double square by double z square terms.

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Summary Hydrogen atom - spherical polar

$$\hat{H} = -\frac{\hbar^2}{2m_e} \left\{ \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right\} - \frac{Ze^2}{4\pi\epsilon_0 r}$$

$$\hat{H} \Psi(r, \theta, \phi) = E \Psi(r, \theta, \phi) \quad \hat{H} \psi(x, y, z) = E \psi(x, y, z)$$

So, the summary of doing that calculation and if you are doing it for the first time, about two to three hours is what the time that you have to give in order to add all these terms and cancel and arrive at the final form. But I will write the final magic form that everybody uses for solving the hydrogen atom Hamiltonian in polar spherical coordinates.

The Hamiltonian is minus h bar square by 2m e 1 by r square dou by dou r of r square dou by dou r plus 1 by r square sine theta dou by dou theta sine theta dou by dou theta plus 1 by r square sine square theta dou square by dou phi square, all of which is the transformation of the derivatives to the spherical polar form.

And therefore, this is nothing but the kinetic energy term in terms of the spherical polar coordinates with the potential energy minus Z e square by 4 pi epsilon naught r and the equation, that you are looking for solving is the H psi r theta and phi is equal to E psi r theta and phi. Instead of the H psi x, y, z in terms of E psi x, y, z, the wave functions are different in the different polar, different coordinate systems.

But please remember, the energy, which is independent of the coordinate representation will not be different between different coordinate system. How you represent your coordinates should not lead to any changes in the Eigen value for the hydrogen electron. And therefore, the, the traditional method is to use the spherical polar coordinates and

that allows the wave function to be separated into an r dependent wave function only, a θ dependent wave function only, and the ϕ dependent wave function only.

If you recall, the particle in the two-dimensional box where we had an x, y dependent wave function being separated into an x only wave function term and on y only wave function term and we were able to get the energies and the solutions, etcetera. Therefore, separation of variables is far more detailed here in the case of hydrogen atom.

Now, let me stop with this as the focal point for the next part of the lecture on what is called the substitution of the wave function in terms of the three radial only, polar θ angle dependent only and Azimuthal angle ϕ dependent only functions, and how we separate this into three different equations. We will not solve them, but in the second part we will look at the solution. And in the third part, we will see some physical representations of the wave functions themselves, the real and the imaginary parts. So, let me stop with part one here, we will continue exactly from this in the next part.

Until then thank you.