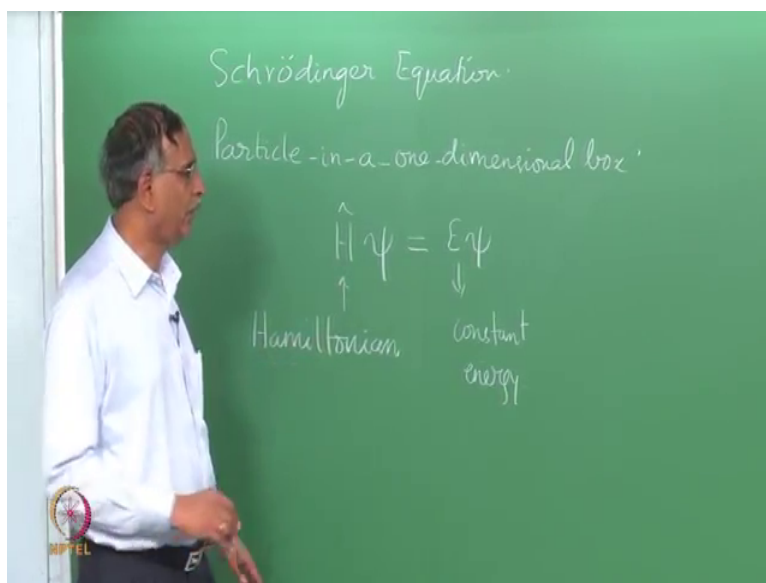


Chemistry Atomic Structure and Chemical Bonding
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Lecture - 03
Schrodinger Equation: Particle in a one Dimensional Box

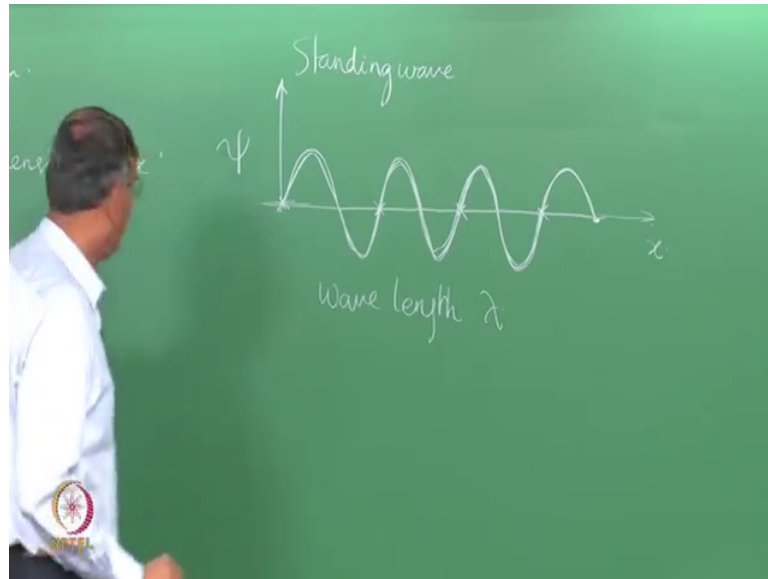
Welcome back to the lecture for the Introductory Chemistry using Schrodinger and Quantum Mechanical Methods for the Atomic Structure. So, what we would do in this and in the next segment is introduce the Schrodinger equation and also do a model problem using the particle in a one dimensional box model this is one of these simplest models that we have. Let us take a quick look at the Schrodinger equation in the lecture earlier I mentioned that.

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I would be talking about the time independent Schrodinger equation, in which this quantity was referred to as the Hamiltonian. And this as a constant, but with dimensions of energy and the function psi is the function that we wanted to find out by solving an equation of the sort, but we do not know what this is right now we have to introduce that to understand how this equation comes about or what is its origin.

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We can do a very simple example of a standing wave, and you know that a standing wave is something that happens between fixed points and the wave motion of a particle fixed to the end something of that kind and let me put it precisely. So, that the wave when it reflects it still follows. And therefore, the standing wave remains as a wave and the amplitudes do not cancel each other. So, if you if you want to look at the axis this is the coordinate or the x axis that you might want to talk about and this is the axis for the amplitude of the wave at any position x between some fixed points ok.

Obviously for this wave the length of the repeating unit is; obviously, called the wavelength λ . And here we have 1 2 yes 2 this is 1 and this is 2 and then you have 3 and 3 and half, it has to be either exactly half wavelength or a full wavelength for this to be a standing wave ok. Now the equation for the standing wave for the amplitude a or let us call that amplitude as psi in relation to what we have here. We will see later that this psi is not necessarily the same as the psi that we talked about, but for that psi if we have the maximum amplitude as a this quantity as a.

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$$\psi(x) = A \sin\left(\frac{2\pi}{\lambda} x\right)$$
$$\frac{d\psi}{dx} = \frac{2\pi}{\lambda} A \cos\left(\frac{2\pi}{\lambda} x\right)$$
$$\frac{d^2\psi}{dx^2} = -\frac{4\pi^2}{\lambda^2} \psi$$

Then the wave function ψ of x is written as a $\sin 2\pi$ by λ of x . This is something that you are familiar with for a standing wave. Now this quantity ψ when you differentiate twice it satisfies the derivative equation. Let us do that for the first derivative of $d\psi$ by dx as 2π by λ times $A \cos 2\pi$ by λ of x .

And the second derivative $d^2\psi$ by dx^2 is equal to minus four π^2 by λ^2 of ψ , because this will become $\sin 2\pi$ by λ of x and that is the same thing as ψ of x . Therefore, you see that the standing wave satisfies the differential equation $d^2\psi$ by dx^2 where ψ is the amplitude of the wave with λ the wave length associated with that.

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$$\text{De Broglie: } \lambda = \frac{h}{p}$$
$$\frac{d^2 \psi}{dx^2} = - \left(\frac{4\pi^2}{h^2} \right) p^2 \psi$$
$$-\hbar^2 \frac{d^2 \psi}{dx^2} = p^2 \psi$$

Now, the Broglie if you remember in the lecture earlier gave an expression for the matter waves lambda in terms of the momentum of the particle in terms of momentum of the particle you have here. And therefore, if I write the wave equation its d square psi by dx square which is equal to minus 4 pi square by h square multiplied by p square psi or minus h bar square. We know that h by 2 pi is h bar therefore, if we bring that in is minus h bar square d square psi by dx square is equal to p squared psi ok.

This is the equation for the standing wave using the de Broglie idea and the quantization idea namely that the energy quantum for material particles light etcetera given in terms of the Plank's constant. So, the Plank's constant enters naturally here in describing what happens to the momentum square on the wave function is the same thing as the second derivative on the wave function multiplied by minus h bar square.

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$$\frac{p^2}{2m} \psi = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi$$

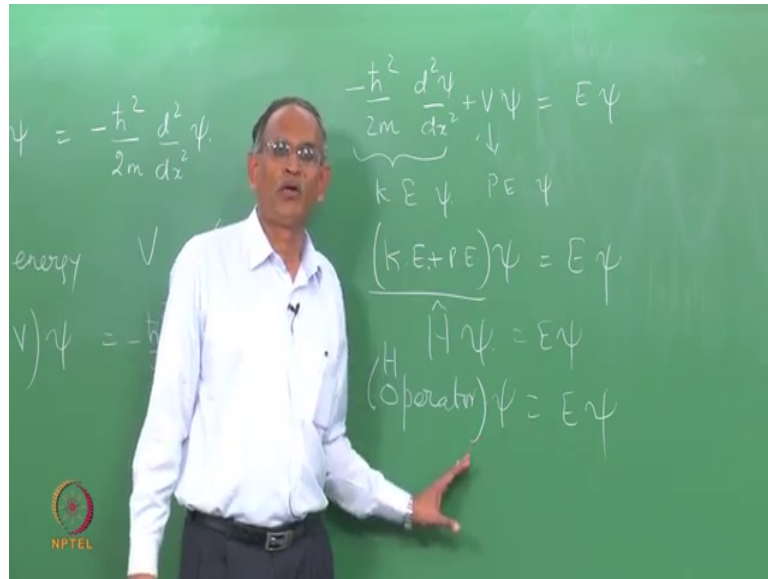
↑
Kinetic energy

$$(E-V) \psi = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi$$

Therefore, if we write the kinetic energy p^2 by $2m$ ψ that turns out to be minus \hbar^2 by $2m$ d^2 by dx^2 ψ . This being the kinetic energy this is the difference between if there is a potential energy v , then it is a difference between the total energy E and the potential energy v , which may be a function of x for whatever I mean if there is a potential we have to consider that.

Therefore what happens this p^2 by $2m$ is nothing but E minus v on ψ giving you minus \hbar^2 by $2m$ d^2 by dx^2 ψ ok. Now one last step and then you see the equation $\hbar^2 \psi$ is equal to $E \psi$ making sense to us because, now if you bring the way here.

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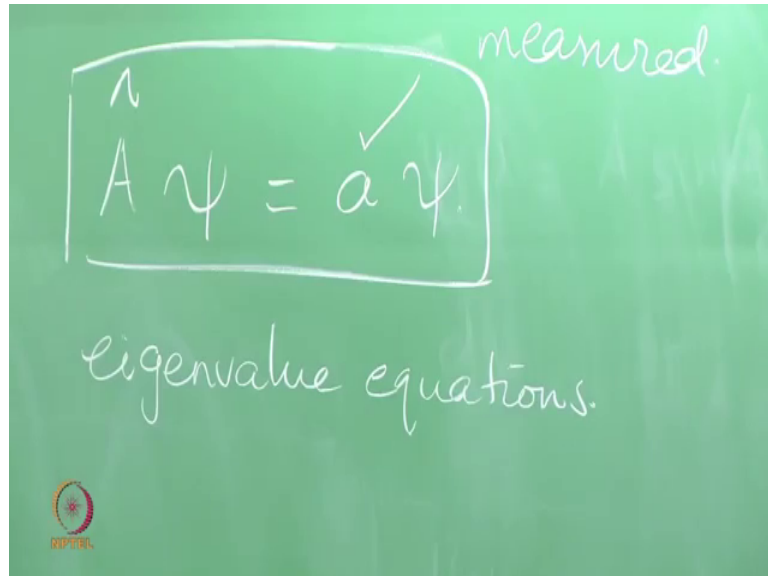
Just rewrite the equation you have minus h bar square by 2 m d square psi by dx square plus v of psi is equal to E of psi, please remember we had already written this as the kinetic energy. And this is on psi this is the potential energy on psi. And therefore, you say that this is nothing but kinetic energy plus potential energy on psi, I am sorry giving you a constant times E psi and. So, you see that this is nothing but the Hamiltonian on psi giving you E psi.

This is a very simple justification I do not think we cannot really say that we have derived it from any fundamental principles or whatever, it is a justification to see from a simple standing wave picture and using that the Broglie principle or the proposition with the Plank's constant. It looks like the particle wave function satisfies the equation Hamiltonian, but the Hamiltonian looks somewhat odd it has a derivative instead of the p square by 2 m that we have now we have put that derivative here. And therefore, the Hamiltonian is a derivative acting on the wave function. And the potential which is of course, a function of the position of whatever particle or the system that you talk about the potential generally multiplies the wave function, but the two together is actually an operator acting on psi.

The Hamiltonian operator acting on psi giving you a constant times psi Schrodinger equation is a very specific equation for the Hamiltonian operator. And such equations in

mathematics are known as eigenvalue equations for whatever quantities that appear here, suppose instead of h there is any other operator that we are going to look at.

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$$\hat{A} \psi = a \psi$$

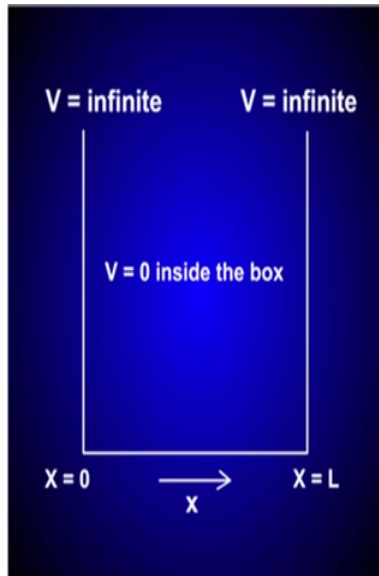
measured.

eigenvalue equations.

$A \psi$ any operator giving some constant times ψ please remember this constant has to have the same dimension as the as the operator a here. In the same way that this constant has the energy dimension for the Hamiltonian operator which is also energy.

Any such equation in which a can be measured experimentally such equations are called Eigen value equations. Eigen value equations and the Schrodinger equation the time independent Schrodinger equation is the eigenvalue equation for the Hamiltonian or the energy operator this is the picture that you have to have. So, let me give you some small problems associated with whatever we have done right after this, but then we will go to the next part namely how do we solve this for the specific case of a simple model, now what is a model.

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Let us look at the modeling now of the particle in a one dimensional box, I have a small drawing here that tells you that we have a particle in a finite region.

The potentials are infinite at 2 points namely points with x equal to 0 and the point x is equal to L meaning that the particle is confined to a region of a box of length L and the particle motion, or the particle coordinate is only one coordinate or one variable namely x . Let us assume for the time being that the potential inside the box is 0 ok. So, this is what we call as the particle in a one dimensional box with infinite barriers and what does this particle give you.

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$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + V\psi = E\psi$$

↑
 $V = \infty \quad \psi = 0$

$x = 0, x = L \quad \psi(x) = 0$

Inside: $V = 0$

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} = E\psi$$

Now, let us look at the equations we have minus \hbar squared by $2m$ $d^2\psi/dx^2$ plus V of ψ is equal to E of ψ , if the potential is infinite then ψ has to be 0.

In order to satisfy that therefore at the boundaries x is equal to 0 x is equal to L the wave function ψ of x is 0. Inside the box we have V is 0. Therefore, what we have is minus \hbar squared by $2m$ $d^2\psi/dx^2$ is equal to $E\psi$ ok; the total energy because there is no potential inside the box if you solve this in a very quick.

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$$\frac{d^2\psi}{dx^2} + k^2\psi = 0 \quad k^2 = \frac{2mE}{\hbar^2}$$
$$\psi(x) = A \cos kx + B \sin kx$$

A, B arbitrary Const.

$$\psi(0) = 0 \quad A = 0$$
$$\psi(L) = 0 \quad \psi(x) = B \sin(kL) = 0$$
$$\sin kL = 0 \quad \text{or} \quad kL = \underline{\underline{n\pi}}$$

Manner namely $\frac{d^2 \psi}{dx^2} + k^2 \psi = 0$, where $k^2 = \frac{2mE}{\hbar^2}$ this is the k^2 is positive; obviously, and therefore, what you have here is a simple derivative equation for second order and you know such functions can be obtained the solutions can be obtained from either trigonometric function or the exponential with imaginary argument.

Let us use the trigonometric function namely a sin, let's write that to be consistent we have $a \cos kx + b \sin kx$. Where a and b are arbitrary constants, arbitrary constants. Now if you look at that solution with the boundary condition that you have namely $\psi(0) = 0$ immediately you have a is equal to 0, because $\cos kx$ is 1 and $\sin kx$ goes to 0; therefore, a is equal to 0. If you have ψ at L which is the other extreme of the box please remember this model at x is equal to L at this point therefore, we have $\psi(L) = 0$ which implies that since a is already 0 ψ of x is $b \sin kL$ and that is equal to 0 ok.

We do not want b to be 0 because if a and b are 0 that is anyway it is a trivial solution for any such a differential equation does not give you any anything of interest I mean there is no meaning there is no interpretation. Therefore we are going to consider the case; obviously, a non-trivial solution would be not equal to 0, which means $\sin kL$ has to be 0 or kL has to be an integer time's π n is an integer kL is equal to $n\pi$.

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$$kL = n\pi \quad n = 1, 2, 3, \dots$$

$$k = \frac{n\pi}{L} \quad k^2 = \frac{2mE}{\hbar^2}$$

$$\frac{n^2 \pi^2}{L^2} = \frac{2mE}{\hbar^2} \quad \downarrow$$

$$E = \frac{\hbar^2 n^2}{8mL^2} \quad \psi(x) = B \sin kx$$

$$= B \sin\left(\frac{n\pi x}{L}\right)$$

And n has to be; obviously, we do not want n equal to 0 which is also the case of triviality and so, what we have is n equal to 1 2 3 etcetera. Integers or please remember k

is equal to $n\pi$ by L look at this k^2 if you recall is $2mE$ by \hbar^2 . Therefore, this gives you immediately that $m^2\pi^2$ by L^2 is equal to $2mE$ by \hbar^2 times the $4\pi^2$ that we have cancel things off and you immediately get the solution namely E is equal to $\hbar^2 m^2$ by $8mL^2$.

And what is the solution for the wave function ψ of x is $b \sin kx$ which is $b \sin n\pi x$ by L because k is $n\pi$ by L ok. So, this is the simplest solution, but two important results one is that the energy for the particle in the box which is subject to boundary conditions that the wave function vanishes at some boundaries subject to that the particle energy appears to be quantized is not arbitrary. You recall the dimension the quantity \hbar^2 by $n^2 \hbar^2$ by $8mL^2$ the quantity.

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Handwritten notes on a digital whiteboard:

- Energy: $\left(\frac{\hbar^2}{8mL^2}\right)$ energy m, L
- Quantized energy values: $1, 4, 9, 16, 25$
- Wave function: $\psi(x) = B \sin\left(\frac{n\pi}{L}x\right)$
- Max Born: $\psi(x)\psi(x)dx = \psi^2(x)dx$
- Probability \rightarrow x and $x+dx$

Has the dimension of the energy and it has the only two inputs, which is which are the inputs for this problem namely the mass of the particle m and the length of the box L and the other constant is of course, Planck's constant. So, now the energy seems to be quantized in terms of the 2 physical parameters that we introduced which particle a larger part a heavier particle or a lighter particle in a smaller box or in the larger box, but with all the other conditions being the same. Namely, potentials being 0 inside the potentials being infinite given that you see that the energy is discretized and the energy is in the units of \hbar^2 by $8mL^2$ this is the fundamental unit for this box and then it is 1

4 9 16 25 as the value of n becomes 1 2 3 four etcetera therefore, particle, particle energies are discretized.

The second part is the other namely the wave function is given in terms of $b \sin n \pi x$ by L. Now what is this wave function? From the beginning of this lecture you might think that this wave function is essentially a function telling you how the particle is oscillating that is not true. That picture was a starting point for us to get an idea that the Schrodinger equation is like this the wave function that we have here is not a function representing how the particle is moving it is just a function associated with that particle, what is the meaning of it Max Born gave the interpretation namely that wave function by itself does not have any meaning, but $\psi^* \psi$ in this case ψ is real. Therefore, $\psi^* \psi$ or ψ^2 .

In a small interval dx gives the probability of the particle being in the position between x and $x + dx$. The probability of locating the particle between x and $x + dx$ that is the number given by the product of the wave function with itself in this case, because its real that Max Born suggested that $\psi^2 dx$ gives the probability that the system we found in the interval x and $x + dx$ that is all there is to it. Therefore, let me conclude immediately what we should be because if $\psi^* \psi$.

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$$\int_0^L \psi^*(x) \psi(x) dx = \int_0^L |\psi(x)|^2 dx$$

$$\int_0^L \psi(x)^2 dx = \frac{1}{2} \Rightarrow B^2 \int_0^L \sin^2\left(\frac{n\pi x}{L}\right) dx = 1$$

$$B = \sqrt{\frac{2}{L}}$$

Which is the same as $\psi^2 dx$ is a probability then if you add all the probabilities from 0 to L, because the particle can have any position between the

endpoint, but not at the endpoint from anywhere as close to the endpoint as possible, but as close to the other endpoint. Therefore, if you integrate the total probability is this being a continuous function you have 0 to L psi x square dx. That probability has to add to 1 because we have made sure that the potentials are infinite in our model therefore, the particle cannot be found outside of the region before the probability that the particle stays inside the box is 1.

This gives you immediately your value for b because you have b squared sin square n pi x by L dx between 0 and L. And that is equal to 1 which gives you a value b is equal to root 2 by L. Therefore, you have got two results for the particle in the box.

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Handwritten notes on a whiteboard showing the wave function and energy levels for a particle in a 1D box. The wave function is given as $\psi_n(x) = \sqrt{\frac{2}{L}} \sin \frac{n\pi x}{L}$ and the energy levels are given as $E_n = \frac{\hbar^2 n^2}{8mL^2}$. Below the equations, a horizontal line separates the formulas from the text. The text reads "Particle 1d." followed by two arrows pointing to "Quantization/discretization of E" and "Probability." The NPTEL logo is visible in the bottom left corner.

Namely, the wave function is root 2 by L sin n pi x by L and E the particles the energy is given by h square n square by eight m L square. Now, because the energy is given by the quantum number m let me use a highlighter here, because it is given by n and n can take any number of values and for that n the corresponding wave function is sin n pi x by L.

We see that there are many solutions to the wave function and many solutions to the energy this will also turn out to be a general property when we solve the Hamiltonian equation the Schrodinger equation for the systems in all the other models that in s in one step you will get all the different types of all the possible energies and all the possible wave functions and the best way to I mean a convenient way I would not call it a

convenient way is to label the wave function with the quantum number $\psi_n(x)$ or E_n for a given quantum number n .

So, let me summarize and then stop for this lecture namely the particle in a 1 d box has two results a quantization of energy or discretization due to boundary conditions and of energy E and a probability statement for determining the position of the particle in the box at various locations. Let us continue this in the next part and complete the remaining that we needed to do in terms of what are called the measurable and then how do we interpret this probability. So, on for various values we will do that in the second part until then.

Thank you.