

Seismic Analysis of Structures
Prof. T.K. Datta
Department of Civil Engineering
Indian Institute of Technology, Delhi

Lecture – 18
Frequency Domain Spectral Analysis (Contd.)

In the previous lecture we discussed about the response analysis of structures subjected to multi points excitation.

(Refer Slide Time: 00:34)


➤ For multipoint excitation:

$$S_{\ddot{x}_g x} = H M r S_{\ddot{x}_g} r^T M^T H^* \quad (4.77)$$

$$S_{\ddot{x}_g x} = -H M r S_{\ddot{x}_g} \quad (4.78)$$

$S_{\ddot{x}_g}$ is of size $s \times s$; r is of size $n \times s$

➤ If all displacements are not required, then a reduced is $\bar{H}(\omega)$ used of size $m \times n$ & S_{xx} is given by

$$S_x = \bar{H}_{m \times n} M r S_{\ddot{x}_g} r^T M^T \bar{H}_{n \times m}^* \quad (4.79)$$


And how we obtain the power spectral density function matrix for the responses given the power spectral density function of the ground excitation and these power spectral density function of ground excitation was converted to a power spectral density function matrix of the excitations at different supports with the help of the inference coefficient matrix r and the various vectors or the vectors of excitation at different supports that lead to a matrix called $S_{\ddot{x}_g}$ matrix. Which denotes the power spectral density function matrix of the excitations at different supports in which the diagonal terms are the power spectral density of the excitations at different supports and the cross power spectral density terms represent the cross power spectral density between different excitations at different supports, and that is we are representing with the help of the matrix $S_{\ddot{x}_g}$ that is a matrix and it has cross power spectral density terms.


Once we have these matrix defined and the r the influence coefficient matrix, that is known then one can get the power spectral density function matrix of the displacements or responses $S \times x$ this will not be $S \times \text{double dot } g$ this will be $S \times$. So, $x \times$ as $H \ m \ r \ S \times \text{double dot } g$ then the transpose of this quantities $H \ m \ r$ on this side that is $r \ T \ m \ T$ and H becomes now a complex conjugate and transpose of that.

So, that is a formulation that we have proved before and if we were wanting to obtain the cross power spectral density function between the excitation and the response displacement that is $S \times \text{double dot } g \times$ that is equal to minus $H \ m \ r \ x \text{double dot } g \ r \ x \text{double dot } g$ is a matrix of size S by S where S is the number of supports and r of course, is a coefficient matrix of the size n into S . Also we discussed in the previous lecture that many a time we are not interested in obtaining all the responses, but we may be interested in only a selected few. In that case the a reduced frequency response function matrix is obtained from the total frequency response matrix of the system and that we denote by H bar and using that H bar one can obtain the power spectral density function of the selected or matrix of the selected responses.

(Refer Slide Time: 04:38)

➤ Without the assumption of ergodicity, Eqns 4.75- 4.78 can be derived (1-3).

Example4.2: (a) For example 4.2, if a time lag of 5s is introduced between supports, find the PSDFs of u_1 & u_2 ; b) For example 3.9, find PSDFs of the degrees of freedom 4 & 5 for correlated and partially correlated excitations (with time lag 5s).



Then we saw as an example we solved a problem.

(Refer Slide Time: 04:40)


Solution: For example 4.2

$$r = \frac{1}{3} \begin{bmatrix} 1 & 1 & 1 \\ 1 & 1 & 1 \end{bmatrix}$$

$$\text{coh}(i, j) = \exp\left(-\frac{|r_{ij}| \omega}{2\pi v_s}\right)$$

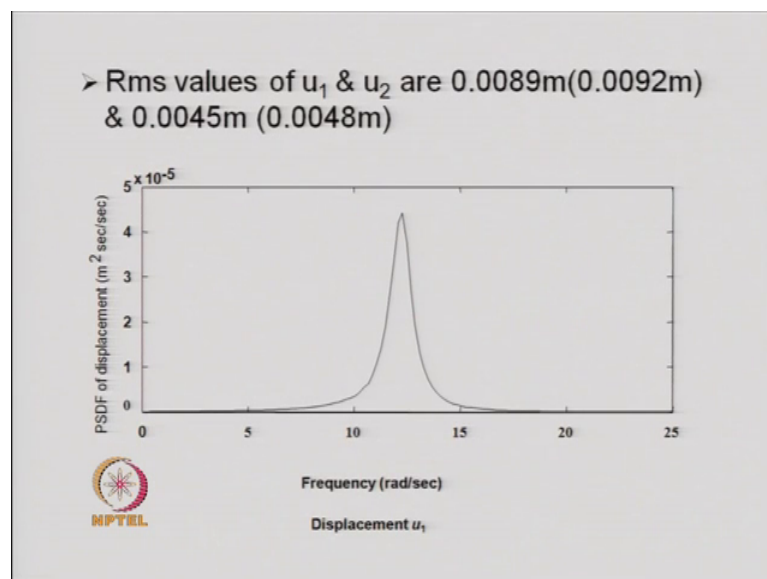
$$S_{\ddot{x}g} = \begin{bmatrix} 1 & \rho_1 & \rho_2 \\ \rho_1 & 1 & \rho_1 \\ \rho_2 & \rho_1 & 1 \end{bmatrix} S_{\ddot{x}g} \quad \rho_1 = \exp\left(-\frac{5\omega}{2\pi}\right) \quad \rho_2 = \exp\left(-\frac{10\omega}{2\pi}\right)$$

PSDFs & cross PSDFs are shown in Fig4.12



In which we had a 3 supports and two nonsupport degrees of freedom u_1 and u_2 and we discussed how we obtain the cross power spectral density function terms of the power spectral density function matrix of the excitation.

(Refer Slide Time: 05:05)



With the help of that we obtain the values of the power spectral density function or power spectral density function of the response u_1 and u_2 and from that the area under the curve provided us the root mean square value of the responses, and these root mean square values of the response were compared with the time domain analysis, note that the

excitations where the elcentro earthquake record and the elcentro earthquake power spectral density function, they are the inputs respectively for the two kinds of analysis.


So, you can see that comparison was very good.

(Refer Slide Time: 05:53)

➤ For the problem of example 3.9, required matrices & values of other parameters are given below:

$$M = \begin{bmatrix} 2.5 & 1.67 \\ 1.67 & 2.5 \end{bmatrix} m \quad \omega_1 = 5.58 \text{ rad/s} \quad \omega_2 = 18.91 \text{ rad/s}$$

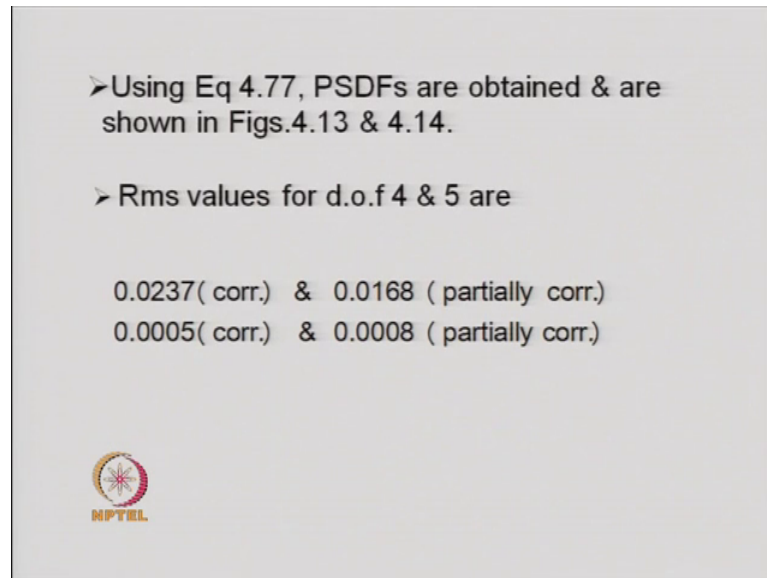
$$K = \frac{EI}{L^3} \begin{bmatrix} 19.56 & 10.5 \\ 10.5 & 129 \end{bmatrix} \quad \sqrt{\frac{EI}{mL^3}} = 2 \quad \alpha = 0.431$$

$$r = \begin{bmatrix} 0.479 & 0.331 \\ -0.131 & 0.146 \end{bmatrix} \quad C = \alpha M + \beta K \quad \beta = 0.004$$


We now solve another problem in this class, the problem of the pitch portal frame in which we had two degrees of freedom if we recall, one at the as the sway degree of freedom on the left hand side and the other one is a vertical degree of freedom at the crown. So, this is the problem that we had solved before for the ground excitations which had a time lag of 5 second, and for that we derived the r matrix and this was the r matrix and also we derived the mass matrix if we recall the mass matrix for such systems is not a diagonal mass matrix, but they are the couple mass matrix although we are using the lump masses at the different degrees of freedom.

So, the derivation of this also was worked out before this is the K matrix. So, with the help of the K matrix and the m matrix the frequencies two frequencies of the system where obtained they are 5.58 and 8.91 radiant per second and with the help of these two frequencies we obtain the values of alpha and beta to construct the C matrix and the C matrix thus was obtained then.


(Refer Slide Time: 07:41)



➤ Using Eq 4.77, PSDFs are obtained & are shown in Figs.4.13 & 4.14.

➤ Rms values for d.o.f 4 & 5 are

0.0237 (corr.) & 0.0168 (partially corr.)
0.0005 (corr.) & 0.0008 (partially corr.)



For obtaining the value of the C matrix, after we have obtained the value of the C matrix then we use the usual formulation that is the equation that I had shown before. So, we plugged into this equation and in constructing H matrix, we have $k - m \omega^2 + I C \omega^{-1}$ that becomes the h matrix.

So, once we know this h matrix then we obtain the values of $S \times x$ and the $S \times \ddot{g}$ matrix was obtained in this particular fashion and the cross terms where one cross term was exponential of $\text{minus } 5 \omega \text{ by } 2 \pi$ because between the 2 supports, the first support and second support there was a time like a 5 second and the first support and the third support the time lag was 10 second. So, therefore, they appear like this one row 1 and row 2 in to this matrix and the results where compared and for the two cases that is in one case we had a perfectly correlated excitation that is at the two supports the excitations where same and in another case we obtained the excitation, having a time lag of 5 second.

(Refer Slide Time: 09:41)

PSDFs of absolute displacements


➤ Absolute displacement is obtained by adding ground displacement to the relative displacement

$$x_a = Ix + rx_g \quad (4.80)$$

➤ PSDF of x_a is obtained using Eqn. 4.37

$$S_{x_a} = S_{xx} + rS_{x_g}r^T + IS_{xx_g}r^T + rS_{x_g}I^T \quad (4.81)$$

$$x(\omega) = -HMr\ddot{x}_g(\omega) = HM\omega^2x_g(\omega) \quad (4.82)$$

$$S_{x_g} = HM\omega^2S_{\ddot{x}_g} \quad \text{and} \quad S_{xx_g} = S_{x_g}^T \quad (4.83)$$


Next we discussed about the power spectral density function of absolute displacements and in order to obtain these power spectral density function of the absolute displacement what we have to do is that, we have to add to the relative displacement x to the displacement or quasi static displacement that is introduced at the degrees of freedom because of the ground displacements at different supports. So, in order to get that we multiply the x_g vector that is the excitation vector at the different supports with the r matrix r matrix, and that gives us the contribution of the different displacements of the excitations to the nonsupport degrees of freedom and the relative displacement was x that is the with respect to the base that displacements of the different nonsupport degrees of freedom and that is multiplied by or pre multiplied by a identity matrix I .

So, this becomes the addition of these two becomes the absolute value of the total value of the displacement at different degrees of freedom and once this relationship is obtained then using this relationship one can write down the value of the power spectral density function matrix like this, that is $S_{x_a}x_a$ is for this term because $I S_{xx} I^T$ will be S_{xx} itself, then $r S_{x_g} r^T$ where S_{x_g} is the power spectral density function matrix of the ground displacement and then $I S_{xx_g} r^T$ and $r S_{x_g} I^T$.

So, they are the cross power spectral density function matrix between the displacement and the ground excitation or ground displacement and this is the other part of the cross power spectral density function matrix which generally is a complex conjugate of this

and transpose that you have seen before. So, once we have this expression in this expression this is known this we have already obtained this is also given, because if the power spectral density function on matrix of the ground acceleration is given to us then from that one can construct the power spectral density function matrix of the ground displacement. And in order to get the value of S_{xg} and xg_x we use this equations that is first we obtain the value of the cross power spectral density function between the excitation and the displacement using this standard expression, in which it is $S_{\ddot{x}g}$ and $S_{\ddot{x}g}$ is replaced with the help of the S_{xg} by multiplying it with the help of omega square.

So, this and the complex conjugate of this S_{xg} is a complex conjugate of S_{xg} . If this matrix is a complex matrix then there will be a star over here star and transpose and if it is a real quantity then the star does not come into picture it is simply a transpose of that.

(Refer Slide Time: 14:09)

Example 4.4: For example 4.3, find the PSDFs of absolute displacement of d.o.f 4 & 5.


Solution: Let X & \ddot{X}_g represent

$$X^T = [x_4 \quad x_5] \quad \ddot{X}_g^T = [\ddot{x}_{g1} \quad \ddot{x}_{g2}]$$

➤ Using Eqn. 4.83

$$S_{x_g x} = H M r \omega^2 S_{\ddot{x}_g} = H M r \omega^{-2} S_{\ddot{x}_g} \quad (4.84)$$

$$S_{\ddot{x}_g} = \begin{bmatrix} c_{11} & c_{21} \\ c_{12} & c_{22} \end{bmatrix} S_{\ddot{x}_g} \quad c_{ij} = \text{coh}(i, j) \quad (4.85)$$

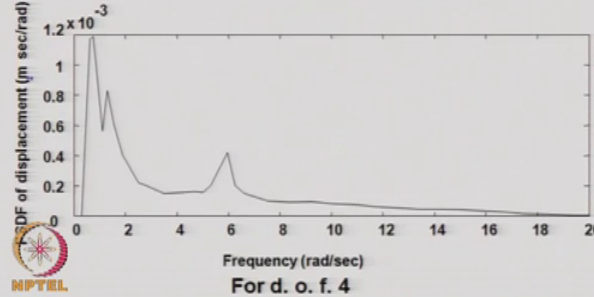
$$S_{x_g x} = \begin{bmatrix} S_{x_{g1}x_4} & S_{x_{g2}x_4} \\ S_{x_{g1}x_5} & S_{x_{g2}x_5} \end{bmatrix} \quad (4.86)$$


Now with the help of this we wanted to obtain the response of the pitch reportal frame and for that we are interested in finding out the response of the system not only for the not for the relative displacement, but for the total displacement.

(Refer Slide Time: 14:30)

➤ Using Eqn. 4.81, S_{x_a} is obtained.

➤ The PSDFs of $S_{x_{4a}}$ & $S_{x_{5a}}$ are shown in Figs 4.15 (a-b). rms values are given as 0.052m & 0.015m respectively.



So, the for obtaining the total displacement we have to know the first $S \times g \times$, $S \times$ matrix. So, that is given in this particular form and the since $S \times g$ and $S \times \ddot{g}$ they are related by a relationship that $S \times \ddot{g}$ is equal to ω to the power 4 into $S \times \ddot{g}$ that we have seen before. Then if you substitute that relationship over here then this become $H M r \omega$ minus to the power two into $S \times \ddot{g}$ not $S \times g$ and this $x \times \ddot{g}$ is known to us and $S \times \ddot{g}$ can be constructed from the time lag that is given to us that is a 5 second time lag.

So, these $S \times \ddot{g}$ matrix would be equal to this C_{11} multiplied by $S \times \ddot{g}$ mind you this $S \times \ddot{g}$ is a single quantity that is the power spectral density function of the earthquake that is the elcentro earthquake that we have consider and these when it gets multiplied by C_{11} then it gives the power spectral density function matrix between at point 1. And at for point 2 the excitation becomes C_{t2} and into $x \times \ddot{g}$ and; obviously, this C_{11} and C_{22} these values are equal to unity and C_{12} and C_{21} they are the cross power spectral density function on terms and we had seen before that for the same earthquake producing different kinds of excitations at different supports.

They are the cross power spectral density function between the excitations is equal to the power spectral density function of the earthquake that is a single quantity multiplied by the coherence function C_{ij} and square n and its function is given or defined in this particular fashion. And we use some empirical relationship for obtaining the C_{12} or C_{21} etcetera. $S \times g \times$ matrix will look like this that is the cross power spectral density function matrix between the ground displacement and the on nonsupport degrees of

freedom displacement x will be like this, that is S_{xg} that is for the first support and this is x_4 is the displacement square displacement of the pitch reportal frame.

So, S_{xg} S_{x_4} represents the cross power spectral density function, between the excitation one support to the displacement x_4 similarly x_{g1} , x_{g5} that indicates the cross power spectral density function between the excitation at support one with the displacement x_5 and in that fashion one can also describe the cross power spectral density function between the support 2 and the displacements 4 and 5.

So, that gives us the $x_{g \times}$ matrix; however, this $x_{g \times}$ matrix can be computed easily with the help of the above relationship.

(Refer Slide Time: 19:18)

The image shows a handwritten derivation on a whiteboard. It starts with the correlation coefficient $c_{12} = \exp(-\frac{5\omega}{2\pi}) = \rho$ and $c_{11} = 1$. Then, the excitation cross power spectral density function is given as $S_{\ddot{x}g} = \begin{bmatrix} 1 & \rho \\ \rho & 1 \end{bmatrix} S_{\ddot{x}g}$. The displacement cross power spectral density function is derived as $S_{xg} = \frac{HM^T}{\omega^2} S_{\ddot{x}g} = \frac{HM^T}{\omega^2} \begin{bmatrix} 1 & \rho \\ \rho & 1 \end{bmatrix} S_{\ddot{x}g}$. This is further simplified to $S_{xg} = \frac{1}{\omega^2} \left[(K - M\omega^2) + iC\omega \right]^{-1} M^T \begin{bmatrix} 1 & \rho \\ \rho & 1 \end{bmatrix} S_{\ddot{x}g}$, where the matrix $\left[(K - M\omega^2) + iC\omega \right]$ is labeled as 2×2 . Finally, it is written as $S_{xg} = \frac{1}{\omega^2} A(i\omega) S_{\ddot{x}g}$, where $A(i\omega)$ is also labeled as 2×2 . An NPTEL logo is visible in the bottom left corner of the whiteboard image.

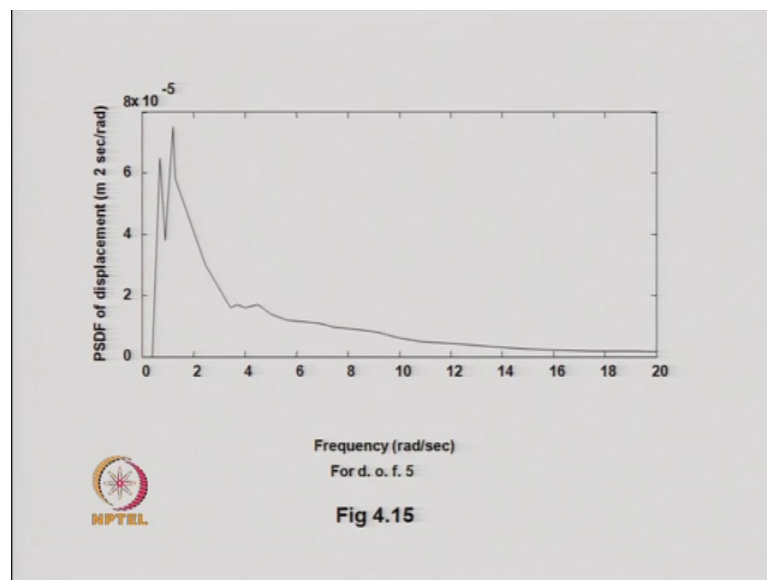
That is the relationship that we have used and once we take that relationship then what first we have to obtain is the cross power spectral density function matrix of excitation $S_{x \text{ double dot } g}$ is the elcentro power spectral density function, and the C_{11} C_{22} C_{12} terms they will be this will be 1 1; obviously, and this will be the row a row will be equal to exponential of minus 5 omega by 2 pi is equal to row.

So, one can work out these value of C_{12} for every value of omega and so, this matrix is completely defined and then we write down the cross power spectral density function matrix between the ground excitations and the displacement x and that is written as HM^T x $S_{x \text{ double dot } g}$ divided by omega square, that becomes the relationship for $S_{x \text{ double dot } g}$

therefore, the this particular term is also known and we have obtained $S_{x \times g}$ and $x_{g \times j}$ and therefore, the all the terms in this equation are known.

So, one can get the power spectral density function matrix of the absolute displacement of the pitch reportal frame for degrees of freedom 4 and 5, 4 is the sway degree of freedom and 5 is the degree of vertical degree of freedom at the crown. So, that is how these power spectral density function earns of the degree of freedom 4 and degree of freedom 5 a represents the absolute value. So, the power spectral density functions of these two quantities were obtained and plotted against frequency that is for degree of freedom 4 and this is for.

(Refer Slide Time: 24:43)



Degree of freedom 5 and one can see that the for degree of freedom 4 the value is more compare to the value of the power spectral density function of the degree of freedom 5, because this is of the order of ten to the power minus 5 the other one is of the order of 10 to the power minus 3. And the area under the curve of course, provides us the root mean square value of the absolute displacement of the degree of freedom 4 and degree of freedom 5.

So, that is how one can obtain the power spectral density functions of the absolute displacements if it is required.

(Refer Slide Time: 25:37)

PSDF of member forces

➤ Consider the frame in Fig 3.7; x is the vector of dyn. d.o.f. & θ as that of the condensed out d.o.f.

$$\theta = Ax \quad (4.87)$$

$$S_{\theta\theta} = AS_{xx}A^T \quad (4.88)$$

$$S_{x\theta} = AS_{xx} \quad \& \quad S_{\theta x} = S_{x\theta}^T \quad (4.89)$$

➤ Consider column i-j; displacement at the ends of column are

$$\delta = [x_i, \theta_i, x_j, \theta_j]^T \quad (4.90)$$

$$f = K\delta \quad (4.91)$$

Next you come to the power spectral density function of the member forces, once we obtain the power spectral density function of the displacement that is the displacement of the various translational degrees of freedom, then from that one has to obtain the power spectral density functions of the member end forces and for that we take say this particular frame and here.

(Refer Slide Time: 26:23)

The diagram shows a frame structure with four supports labeled $x_{g1}, x_{g2}, x_{g3}, x_{g4}$. The horizontal displacement is x_1 and the vertical displacement is x_6 . The displacement vector is $\delta = [x_i, \theta_i, x_j, \theta_j]^T$. The force vector is $f = K\delta$. The stiffness matrix in local coordinates is $S_{ff} = K S_{ss} K^T$. The transformation matrix is $\bar{S} = T S$. The relationship between θ and x is given by $\theta = -\frac{S_{\theta x}}{S_{\theta\theta}} x = A x^T$.

Say this is the frame in this frame we have say 1, 2, 3, 4, 4 supports and at these 4 supports we have got 4 different excitations because of the time lag between them and

these are nonsupport degrees of freedom that is x_1 to x_6 they are nonsupport degrees of freedom and say we are interested in finding out the power spectral density function of the bending moment at i and j . So, that is of our interest.

So, then what we do in the beginning that we write down the matrix, the stiffness matrix of the entire structure with these rotations and once we know the rotations then from there one can get the relationship between the rotation and the displacement and this matrix is called A matrix A . Say once we know the relationship between the theta matrix theta vector and the x vector through matrix A then one can obtain the power spectral density function of the rotations in terms of the power spectral density function of the displacements using this equation and the cross power spectral density function between the displacement and the rotation is given by this $A S_{xx}$ and the $S_{\theta x}$ will be simply the transpose of this.

Now, if $x_{\theta x}$ is a complex matrix then of course, there will be a this will be a complex conjugate and then transpose. Now once we have obtained this $S_{\theta\theta}$ that is the power spectral density function matrix of rotation and the power spectral density function matrix of displacement and the cross power spectral density function matrix between rotation and displacement. So, once they are known then from that from those matrices one can select the specific terms of the power spectral density function to construct the power spectral density function matrix, which will be used in finding out the power spectral density function of member end forces.


Now, for example, if I take this column i, j that I had shown in the figure then the values which are required to find out the bending moment at the ends i and j are x_i theta i x_j and theta j . So, these are the particular displacements and rotations which are required. So, we take out that and we write down the relationship between the member end forces and the member end displacement with the help of the usual equation f is equal to K into δ .

(Refer Slide Time: 30:14)

➤PSDF matrix of the member end forces are

$$S_{ff} = KS_{\delta\delta}K^T \quad (4.92)$$

➤If $\bar{\delta}$ are in local coordinates, then

$$\bar{\delta} = T\delta \quad (4.93)$$
$$S_{ff} = KS_{\delta\delta}K^T = KTS_{\delta\delta}T^TK^T \quad (4.94)$$


Once we are able to write that then S_{ff} that is the power spectral density function matrix of the member end forces becomes equal to K into $S_{\delta\delta}$ K^T , this K is known that is the member stiffness matrix and $S_{\delta\delta}$ is the or $\delta\delta$ they are the member end displacement.

So, by knowing $S_{\delta\delta}$ once can easily find the power spectral density function of the member end forces. In many cases we want the member end forces in the local coordinate that is the member end coordinates. Then of course, there at a further transformation is required; that means, if the $\bar{\delta}$ is the displacements in the local coordinate and if δ are the displacement in the global coordinate there you related through this relationship that is $\bar{\delta} = T\delta$ where T is the transformation matrix and substituting for the $\bar{\delta}$ that is one can obtain the power spectral density function matrix of the member end forces in local coordinate, to be equal to K into $S_{\bar{\delta}\bar{\delta}}$ and $S_{\bar{\delta}\bar{\delta}}$ is further converted to $S_{\delta\delta}$ by including the transformation matrix T . So, it is possible to obtain the power spectral density function of the member end forces both in the global and the local coordinates that is they are power spectral density function of individual quantities may be bending moment share force etcetera.

(Refer Slide Time: 32:22)


Example 4.5 : For the example 3.10, find PSDFs of displacements 1 & 2; also find the PSDF of bending moment at the centre.

Solution:

$$K = \begin{bmatrix} 56 & -16 & 8 \\ -16 & 80 & -16 \\ 8 & -16 & 56 \end{bmatrix} \text{ m} \quad \bar{C} = \begin{bmatrix} 0.813 & -0.035 & 0.017 \\ -0.035 & 0.952 & -0.035 \\ 0.017 & -0.035 & 0.811 \end{bmatrix} \text{ m}$$

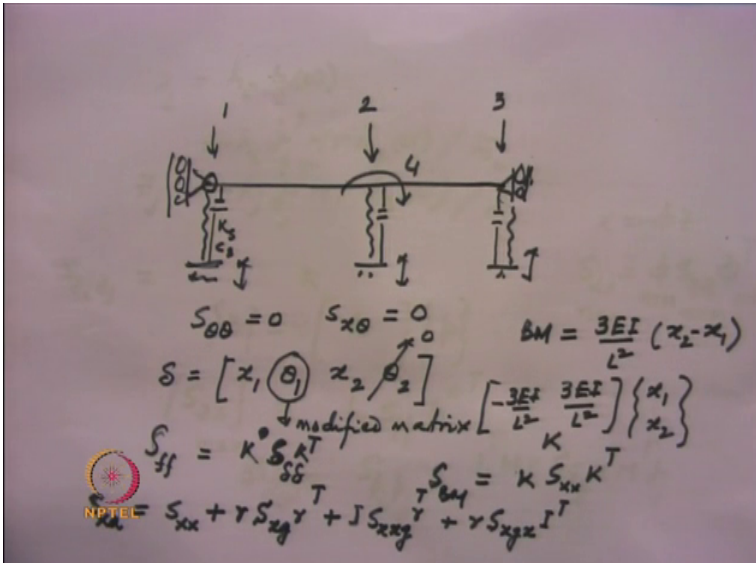
$\omega_1 = 8.0 \text{ rad/s}; \omega_2 = 9.8 \text{ rad/s}$ and $\omega_3 = 12.0 \text{ rad/s}$

➤ Following the same procedure, PSDFs of 1 & 2 are obtained and are shown in Fig 4.16(a-b).



Now, these in this is example these that displacement.

(Refer Slide Time: 32:56)



So, one and two that is the if you recall that is the example problem for the yes I think yeah, this is the problem if we recall we had solved for the for finding out the responses in time domain, where we specified the ground motions at this 3 supports and they are the same ground motion and this was a modal of a pipeline buried pipeline. And these supports there were 3 supports were taken, the soil damping and stiffness matrix or the damping and the stiffness are replaced by spring and dashpots and these values were

given to us and the degrees of freedoms where 1 2 3 they are the translational degrees of freedom and a rotation here.

So, one can easily condense out this rotational degree of freedom and once you condense out this rotational degree of freedom, then with the help of the translational or matrix stiffness matrix corresponding to translation degrees of freedom 1 2 3 one can obtain the responses for the degree of freedom 123.

Now, the stiffness matrix for the system was this 3 by 3 stiffness matrix, we obtained a damping matrix using C bar is equal to αm plus βk , and the 3 frequencies that that are obtained from this k matrix and m matrix so, that where 8 9.8 and 12 radiant per second.

(Refer Slide Time: 35:04)


➤ For finding PSDF of bending moment θ at the center is required

$$x^T = [x_1 \quad x_2 \quad x_3]$$

$$\theta = \frac{3}{4L} [-1 \quad 0 \quad 1]x = Ax$$

$S_{\theta\theta}$ & $S_{x\theta}$ are zero as δ is zero at the center.

➤ Displacement vector θ is

$$\delta = [x_1, \theta_1, x_2, \theta_2]^T$$


And then we obtained the displacement power spectral density function for the 3 displacement and using the same procedure that we have discussed before, but now we are interested in finding out the power spectral density function of the bending moment at the center.

Now, for that we recognize the fact that at the center if all the 3 degrees of or at all the 3 supports the ground excitations are the same then the we expect that there should be not be any rotation over here and in fact, it can be easily shown with the help of the condensational matrix. So, if there is no rotation then $S_{\theta\theta}$ will be 0 and $x S x$

theta also will be equal to 0, that is the power spectral density function of the rotation and the cross power spectral density function of the rotation and displacement they will be equal to 0.

Now, if we wish to find out say bending moment here, then the displacement that we should take this will be the displacement 1 and then if there is a rotation over here say theta 1 this is the theta 1 and then the displacement here x 2 and then the rotation here theta 2. So, if you know them then with the help of that one can obtain the power spectral density function matrix of the bending moment at this point and at this point.

Now, since at this point the bending moment is equal to 0. So, we need not include this theta 1 into this vector and we can modify our stiffness coefficient for bending moment for this member accordingly. So, that these theta 1 need not be considered and since theta 2 is 0 then we get only x 1 and x 2 these are the two displacements if we know this two displacements only one can find out the bending moment at this point, and that is that bending moment can be written by simply by this equation $\frac{3EI}{L^3} x_2 - \frac{3EI}{L^2} x_1$ in the matrix form these becomes $\begin{bmatrix} \frac{3EI}{L^2} & -\frac{3EI}{L^3} \\ -\frac{3EI}{L^3} & \frac{3EI}{L^2} \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix}$ where x 1 and x 2 are the displacements at this two points and this can be further written in this particular form that is the S bending moment, that is the power spectral density function of the bending moment is equal to $K S x x K T$.

So, if we know S x x once can find out the bending moment at this particular point. Now the important part of this analysis is that if the ground motions are not the same at different points then the relative displacements at this point and in this point and from that one cannot obtain the bending moment at this particular point, that is at the center that the way that we have obtained the bending moment.

In that case what we have to do we have to find out the absolute displacement or the absolute displacement power spectral density function of point 2 and the power spectral density function of the absolute displacement at 1 that must be known and the cross power spectral density function must be known or in other words we must have a power spectral density function matrix of the absolute displacements at 123 degrees of freedom. And for this particular problem we need not bother about the rotation here because the rotation becomes 0 if the degrees of freedom are or the excitations are same at these 3 points, but if they are not same and we are wanting to find out the bending moment at

this particular point then we wish to have the power spectral density function matrix of the absolute displacements at 1 2 3 and a power spectral density function of the rotation.

However the power spectral density function of the rotation at this point will not be effected by the ground displacements, that is taking place at this point at different points with a different displacements will not be effected by that this theta can be obtained S theta can be obtained from the relative displacement itself.

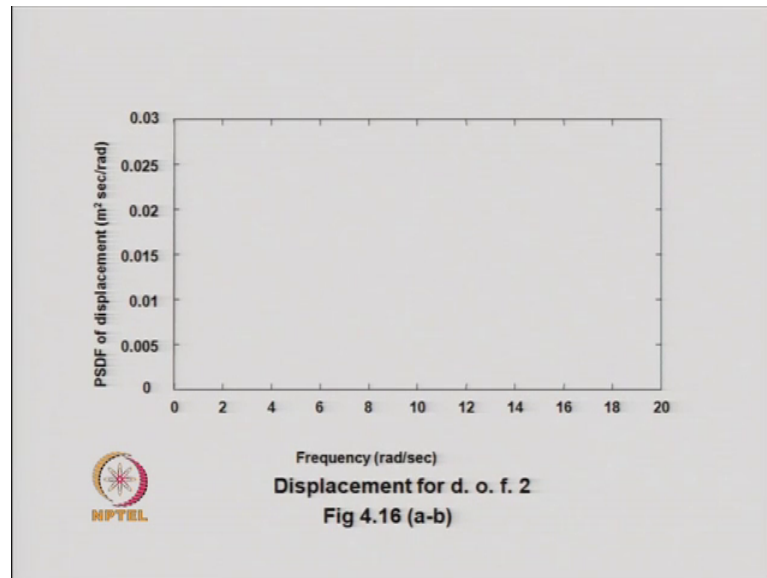
So, therefore, what we do is that in obtaining the cross power spectral density function matrix of these absolute displacements, we go back to the previous problem that we have solved in which we know we should know the power spectral density function on matrix of the relative displacement then power spectral density function matrix of the ground displacement then $S_{x \times g}$ and $S_{g \times x}$ that we had solved in the previous problem. And with the help of that we obtain the power spectral density function matrix of the absolute displacements and then we obtain the bending moment that is $S_{\text{bending moment}}$ here in place of $S_{x \times S}$ here, this $S_{x \times x}$ is the power spectral density function matrix of related displacement that would be replaced by the power spectral density function of the absolute displacements.

So, the problem in which we have got the same excitations at a different supports, the we need we do not have to find out this absolute displacement power spectral density function matrix, we can only obtain the power spectral density function matrix of the relative displacement. However, the problem becomes lightly complex, if we have a power spectral density of the ground excitations at different supports are different and for that remember that one has to find out the power spectral density function matrix of the absolute displacements or total displacements of different degrees of freedom. And also we have to find out the condensed out a power spectral density function matrix of the condense out degrees of freedom that that is like rotations theta theta theta theta and the cross power spectral density function between theta and x.

So, all those things must be known then only one can obtain the power spectral density function matrix of the member end forces. Thus we see that to find out the power spectral density of function of the member forces for excitations which are same as different supports the formulation is simpler, but if you are wanting to find out the power spectral density function of member end forces for a multi support excitation problem,

then one has to obtain the power spectral density function on matrix of the total displacement and then from there we have to take out the relevant terms to form the power spectral density function matrix $S_{\delta\delta}$ and which can be multiplied with the stiffness matrix k and k_t to obtain the power spectral density function of the member end forces.

(Refer Slide Time: 44:56)



So, the these problems the power spectral density function of the absolute displacements where and the bending moment where plotted in this figures.

(Refer Slide Time: 45:02)

- Using modified stiffness for pinned end, θ_1 is eliminated ; x_1 & x_2 are abs displacements.
 - PSDF of bending moment obtained using Eqn. 4.92 is shown in Fig 4.16c. rms values of x_1 , x_2 & B.M. are 0.0637m, 0.1192m & 1.104m .
- NPTEL

And the displacements and bending moment values of the displacement and bending moments where obtained for different quantities.


(Refer Slide Time: 45:16)

Modal spectral analysis

- Advantages of modal analysis have been described before.
- For multipoint excitation, the i^{th} modal equation can be written as

$$\bar{m}_i \ddot{z}_i + 2\xi_i \omega_i \bar{m}_i \dot{z}_i + \bar{m}_i \omega_i^2 z_i = -j_i^T M r \ddot{x}_g \quad (4.95)$$

$$\ddot{z}_i + 2\xi_i \omega_i \dot{z}_i + \omega_i^2 z_i = \frac{-j_i^T M r}{\bar{m}_i} \ddot{x}_g = p_i \quad i = 1, \dots, r \quad (4.96)$$

 $j_i^T M_j$

Next we come to what is called the modal spectral analysis, now since we use the mode super position technique you uncouple the different equations.

(Refer Slide Time: 45:34)

$$z_i = h_i p_i(\omega)$$

$$= -h_i j_i^T M r \ddot{x}_g(\omega) / \bar{m}_i$$

$$z_j = -h_j j_j^T M r \ddot{x}_g(\omega) / \bar{m}_j$$

$$S_{z_i z_j} =$$


$$\{z\} = [h] \{p\}$$

$$[S_{zz}] = h S_{pp} h^* T$$

$$S_{p_i p_j} = j_i^T M r S_{\ddot{x}_g}^T M j_j^T$$

$$x = \phi z$$

$$S_{xx} = \phi S_{zz} \phi^T$$



Of motion of a multi degree freedom system and can simplify the problem that is a multi degree freedom problem can be converted to a set of single degree freedom problems and for each single degree freedom problem one can obtain the responses in the

generalized coordinate and from there one can obtain the displacement in the structural coordinate using the mode shapes.

So, extending that concept one can develop also a modal spectral analysis, the name modal spectral analysis comes because we are using the mode shapes of the structures for developing this method and that is why it is known as the modal spectral analysis and it has the same advantages and that we observed for the case of usual time domain analysis and the frequency domain analysis.

So, for multi support excitation we have these standard expressions that we have discussed in connection with the mode superposition technique. So, the i th generalized displacement velocity and accelerations are z_i , \dot{z}_i and \ddot{z}_i and \bar{m}_i is the i th generalized mass and \mathbf{j}_i^T is the transpose of the first mode shape \mathbf{r} is the influence coefficient matrix associated with $\ddot{\mathbf{x}}_g$ considering $\ddot{\mathbf{x}}_g$ to be a multi support excitation, then one can convert this to this form by dividing by \bar{m}_i and this entire thing is called p_i that is the generalized force at the for the i th mode.


So, or in other words this is the i th modal equation and in that \bar{m}_i is defined in this particular way that is already known to you now.

(Refer Slide Time: 48:24)

\mathbf{r} is the influence coefficient matrix of size $m \times n$.

- $h_i(\omega)$ for each modal equation can be easily obtained.
- $Z_i(\omega)$ can be related to $p_i(\omega)$ by

$$z_i(\omega) = h_i(\omega) p_i(\omega) \quad i = 1, \dots, r \quad (4.97a)$$

$$p_i(\omega) = \frac{-\mathbf{j}_i^T \mathbf{M} \mathbf{r} \ddot{\mathbf{x}}_g(\omega)}{\bar{m}_i} \quad (4.97b)$$


Once we have a single degree freedom equation for a mode i then one can write down z_i omega to be is equal to h_i omega into p_i omega and once we have this then we can easily find out $S_{z_i z_i}$ very easily, that is the power spectral density function of the generalized displacement z that will be is equal to h_i absolute square multiplied $S_{p_i p_i}$ omega and one has to simply find out what is the $S_{p_i p_i}$ omega ok.

Now, these $S_{p_i p_i}$ omega since it has a relationship like this, then one can obtain using the formulation that is a y vector or y is equal to a multiplied by x vector $x \cdot x$ double dot g for that using that relationship we can find out S_y that is the S_p over here, that is the power spectral density function of the force S_p that will be equal to a into S_x double dot g into a^T where a will be equal to these entire this $j_i^T m_r$ divided by m bar that will be the matrix a . So, that one can easily obtain the value of S_p and once you obtain the value of S_p then one can get the power spectral density function for z . Generally we rather than obtaining the power spectral density function matrix of a single z or that is for a particular value of mode we wish to find out the entire z matrix that is $S_{z z}$ that is the power spectral density function matrix of all the z s and say that matrix size will be equal to m by m where m is the number of modes that is considered.

So, in that case we can write down this basic relationship, that is z vector will be equal to a h matrix small h matrix which will be a diagonal matrix, because all the values of z they are uncoupled. So, therefore, with this h matrix will be a diagonal matrix, p is a modal load vector having the size of the same number of modes m and once we have that then we can write down $S_{z z}$ to be is equal to h into $S_{p p}$ into h^T and this is this is known these and these quantities are known only one has to find out what is the $S_{p p}$ matrix.

Now, $S_{p p}$ matrix any term of the $S_{p p}$ matrix can be obtained from this relationship that is $S_{p_i p_j}$ that is equal to $j_i^T m_r \cdot x$ double dot g into $r^T M^T$ and j_j that is between i and j the cross power spectral density function of force that is the modal force p and modal force p_j the cross power spectral density function between them. So, that can be given by this equation and only thing is that you note is that this is the i and this is the j .

So, one should only take note of these i and j over here that is this will be the i th mode shape and this will be j th mode shape and rest of $m_r M$ and r etcetera they are known.

So, one can obtain the value of $S_{p_i p_j}$ mind you $S_{x \ddot{g}}$ here is again a matrix of the power spectral density function of ground acceleration and which requires again the requires a determination of the cross power spectral density function terms, which comes from the coherence function.

So, once we get $S_{p_i p_j}$ where terms or every term of the S_p matrix then $S_{z z}$ is known and once $S_{z z}$ is known then we can go back to our modal equation that is x is equal to $\sum_{i=1}^r \phi_i z_i$ or more super position equation and if this holds good then $S_{x x}$ will be equal to $\sum_{i=1}^r \phi_i S_{z z} \phi_i^T$ and the say if we take only the m mode number of modes $S_{z z}$ will be m by m and there will be n by m and m by n respectively and we will get a power spectral density function of the displacements in the structural coordinate and we will get this matrix size will be n by n .

(Refer Slide Time: 54:43)

> Elements of PSDF matrix of z are given by


$$S_{z_i z_j} = \frac{h_i h_j}{\bar{m}_i \bar{m}_j} j_i^T M r S_{\ddot{g}} r^T M^T j_j \quad i=1, \dots, r, \quad j=1, \dots, r \quad (4.98)$$

> Using modal transformation rule

$$x = \phi z$$

$$S_{xx} = \phi S_{zz} \phi^T \quad (4.99)$$

ϕ is $m \times r$




So, we can see that in the case of the modal spectral analysis, what we do is that we have the same formulation that we have done for the mode super position technique for solving the problem in the frequency domain and time domain, and extend it to obtain the power spectral density function of the generalized coordinate by using small h that is the frequency response function of a single degree freedom system. And we can obtain term by term that is we can obtain the $S_{z z}$ for or each mode or one can obtain the power spectral density function matrix of $S_{z z}$ and the that would be given by this mind you this $S_{z z}$ is not a diagonal matrix it will have cross terms. So, this cross terms denote the

correlation between the power spectral density functions of z_1 and z_i and z_j . So, I stop at this.

(Refer Slide Time: 56:09)

Example 4.6: For problem example 3.11, find PSDF of d.o.f 1 (tower top) and d.o.f 2 (centre of deck). It is assumed that a uniform time lag of 5s between the supports exists.

Solution

$$K = \begin{bmatrix} 684 & 0 & -149 \\ 0 & 684 & 149 \\ -149 & 149 & 575 \end{bmatrix} m, \quad M = \begin{bmatrix} 20 & 0 & 0 \\ 0 & 20 & 0 \\ 0 & 0 & 60 \end{bmatrix} m$$
$$r = \begin{bmatrix} -0.781 & -0.003 & 0.002 & -0.218 \\ -0.218 & 0.002 & 0.003 & -0.781 \\ -0.147 & -0.009 & 0.009 & 0.147 \end{bmatrix}$$


Today and we will solve a problem that is the problem that we solved for a cable stayed bridge and that we will take an example to illustrate the use of the modal spectral analysis.