

Soil Dynamics
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Lecture 11
Single Degree of Freedom System (SDOF) - Part 9

Hello friends. Welcome to the class Soil Dynamics. So far, we have discussed Single Degree of Freedom System under undamped condition or damped by viscous damping. Now, our today's topic is Single Degree of Freedom System with Coulomb Damping.

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The slide is titled "Single Degree of Freedom System with Coulomb Damping". It contains the following text and diagrams:

- ▶ Coulomb damping force: $F_d = \mu N$ (oppose the relative motion) $F_d \rightarrow F_D = \mu N^T$
- ▶ μ is the coefficient of friction and N is the normal force acting on mass " m ".

Fig. 11.1 Spring-mass system with Coulomb damping: A mass m is attached to a spring with stiffness k fixed to a wall. An arrow labeled $+x$ (L-R) indicates the positive displacement direction.

Fig. 11.2 Free body diagram of mass m : Two diagrams, (a) and (b), show the forces on the mass. Diagram (a) shows the mass moving to the right with velocity \dot{x} (L-R). Forces acting on it are: weight W (down), normal force N (up), spring force kx (left), and friction force μN (left). Diagram (b) shows the mass moving to the left with velocity \dot{x} (R-L). Forces acting on it are: weight W (down), normal force N (up), spring force kx (right), and friction force μN (right).

So, first of all how coulomb damping works? We can see these spring mass system with the coulomb damping. So, here we have a mass m which is attached to the spring having stiffness k . Now, when the mass m will be allowed to move in positive x direction, in this case positive x direction means left to right, I can write it here left to right, then what will be happened?

There will be a friction between the mass and the surface on which it is resting. And this friction opposes the motion of the mass m . Then how the free body diagram looks? Here, we can see the free body diagram. In this case, you can see the mass is moving, the velocity of the mass is shown which is left to right. So, figure a shows left to right, figure b shows, you can see the direction of the velocity.

So, this is right to left. Now, when the mass is moving from left to right that means we can see in this figure it is moving from left to right. Then what is happened? The surface on which this mass resting offers a resistance in the opposite direction of motion. In this case, the direction of motion is left to right. So, resistance from right to left, you can see here.

Now, for the case, when motion is occurring from right to left then the surface on which the mass was resting, offers a resistance offers a resistance opposing to the motion. That means in this case, resistance will act from left to right. You can see in this figure, it is occurring from left to right. And what is the magnitude of the damping force because of the coulomb damping? The magnitude of the damping force, you can see F_d is equal to μN .

You can write instead of F_d here, you can make it capital D for damping. So, in place of F small d , I am writing it F capital D . So, then what will be happen? The damping force because of the coulomb damping is capital F_D and the magnitude of that force is μ times capital N . What is μ ? In this case, μ is the coefficient of friction. You can see here. And what is N ? N is the normal force acting on the mass m . Now, let us consider the motion is occurring from right to left first and then left to right.

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The slide displays a free-body diagram of a mass m on a horizontal surface. The forces acting on the mass are: weight W acting downwards, normal force N acting upwards, spring force kx acting to the left, and friction force μN acting to the right. A velocity vector \dot{x} points to the left, indicating motion from right to left. Handwritten notes include: "for x L to R (+ve)", "Motion of mass m from right to left", and "x-direction".

Derivation: $= m\ddot{x}$; Unbalanced force $= -kx + \mu N$
 Inertia force = Unbalanced force $\Rightarrow m\ddot{x} = -kx + \mu N$
 $\Rightarrow m\ddot{x} + kx = \mu N$ (Motion of mass m from 'R' to 'L')

So, when motion is occurring from right to left, what will be happened? Damping force acts from left to right, in this case you can see it. Then if we write the force equilibrium equation then what

will be that? Inertia force in this case, I am writing here, inertia force in this case is how much? Obviously, mass times acceleration. Now, you can find out the unbalanced force. We are talking in x direction. So, better I write it here also, in x direction.

So, unbalanced force is how much? That is minus k x plus mu N. Why it is plus? Mu N and minus k x, minus k x is the spring force whereas damping force is plus mu N. Because we have considered left to right motion, when the when the direction of is x is left to right that is positive. So, for x left to right, we will consider positive direction. In this case, you can see spring force is acting from right to left on the body or on the mass m that is the reason it is with negative sign.

Then in equilibrium, what we can write? This inertia force is equal to the unbalanced force. Then we can write $m \ddot{x}$ is equal to minus k x plus mu N. So, once again I can rewrite it as $m \ddot{x}$ dot plus k x is equal to mu N. So, in this case, motion is occurring from right to left. Better I should write it in full sentence here also. Motion of mass in from right to left. This is the equation of motion in that time.

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Single Degree of Freedom System with Coulomb Damping

Derivation:

Inertia force = $m\ddot{x}$

Unbalanced force = $-kx - \mu N$

$\Rightarrow m\ddot{x} = -kx - \mu N \Rightarrow m\ddot{x} + kx = -\mu N$ (L-R)

$\Rightarrow m\ddot{x} + kx = -\mu N$

Motion of mass m from left to right

Force-I (0 to x_2)

$z = z_{hom} + z_p = A_3 \sin \omega t + A_4 \cos \omega t - \frac{\mu N}{k}$

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Single Degree of Freedom System with Coulomb Damping

x-direction

Derivation: $= m\ddot{x}$; Unbalanced force $= -kx + \mu N$

Inertia force = Unbalanced force $\Rightarrow m\ddot{x} = -kx + \mu N$

$\Rightarrow m\ddot{x} + kx = \mu N$ (Motion of mass *m* from 'R' to 'L') — (1)

$z = z_{hom} + z_p$

$z = A_1 \sin \omega t + A_2 \cos \omega t + \frac{\mu N}{k}$

EP. (Static Equilibrium)

for *x* L to R (+ve)

Motion of mass *m* from right to left

$z_{hom} = A_1 \sin \omega t + A_2 \cos \omega t$

$z_p = \frac{\mu N}{k}$

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Now, next case you can see, when the mass is moving from left to right, in this case mass is moving from left to right then what will be the equation of motion in this case. So, in this case, inertia force is $m \times \ddot{x}$. What about the unbalanced force? In this case, unbalanced force is because of the spring force and because of the damping force offered by the coulomb damping. Then it is $-kx - \mu N$.

Why minus? Because we have considered a when direction is left to right, we will consider it as positive direction. So, then in equilibrium condition, what we can write? Inertia force is equal to the unbalanced force which is in this case $-kx - \mu N$. Then we can write the equation of motion, the final form is $m \times \ddot{x} + kx = -\mu N$. So, this time motion of mass *m* is occurring from left to right.

So, what we have seen? Depending upon the direction of motion of the mass *m*, the equation of motion will change. So, there are two equations of motion in this case. We can write both together as $m \times \ddot{x} + kx = \pm \mu N$. When it is plus? When motion is occurring from right to left. And when it is minus? When motion is occurring from left to right.

But there is a question. So, question is that in the previous case, when motion was occurring right to left, that time you may think, okay, we can understand the direction of μN , which is the direction of the damping force because of coulomb damping, but what about, why kx is in this direction? That means spring is trying to pull the mass *m*.

In this case you can see, spring is exerting pulling force to the mass m . Why so? Actually, when we are saying right to left, basically, what is happen that I would like to explain now. This is the mass in equilibrium, static equilibrium position. So, this is your equilibrium position of mass m , static equilibrium I am writing. Now, what is happened when it is moving from right, it can move in both direction, positive x and negative x also?

Now, in positive x direction, let us take, it reach to the extreme position. Let us say for first case, its let extreme position is x_0 . Now, what is happen? This is the extreme position in the positive x direction, when mass m is subjected to a motion. Now, what is happening? From this extreme position, now the mass will try to go back to its equilibrium position, here, and then it will cross this point and will move towards this direction.

As long as this mass m is, this equilibrium position, let us take 0 to x_0 , what is happening? If you think about the spring which is attached to the mass, how it was attached I am just showing here. Now, when mass is, this is in equilibrium position. Now, when mass is at x_0 , what is happen to this spring? This spring is stretched. And that condition will be there as long as the mass is placed in between 0 to x_0 .

So, x_0 is the extreme position, however, 0 to x_0 , at any distance from or 0 to x_0 , we can say, at any position, what will be happening? The spring will be stretched. If the spring is stretched then what will be happen to the mass m ? That will be pulled by the spring. That is the reason, in this case when mass is moving even from right to left, right to left means x_0 to this position, that time also the spring is exerting pull out force to the mass m .

And we can see the direction of kx which is the spring force is opposite to positive direction. I hope for this case, there is no confusion. Because in this case, mass is moving from left to right. That means, 0 to x_0 position, for first cycle. Then I can write it here for cycle 1. This is the position. So, obviously, spring force will act in negative x direction.

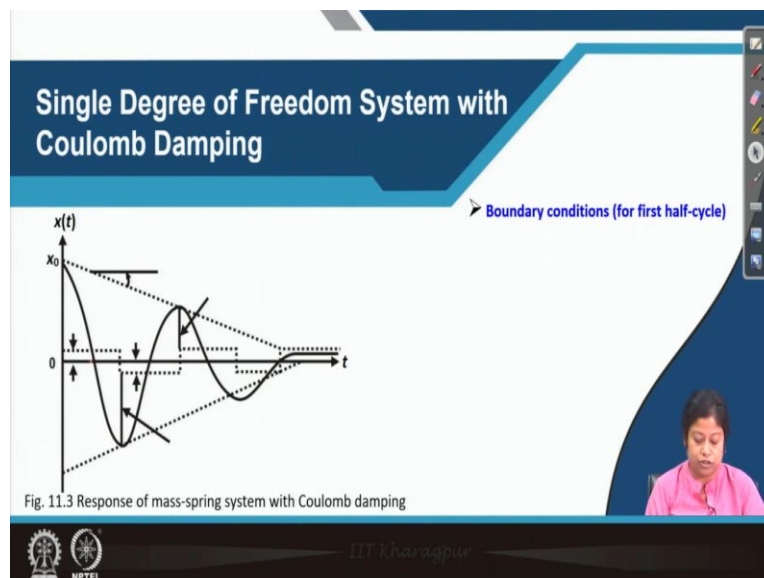
So, in this way, we can derive the equation of motion for the mass spring system with coulomb damping. Now, we need to find out the general solution for this problem. Now, for this type of problem, let me go back to write into left case. So, this is our right to left case. Now, in this case, we can see this equation is no more homogeneous equation. I can give a name to this equation.

This is equation 1. So, this equation 1 is not homogeneous equation. There is a term, non-zero term on the right hand side. So, the solution for equation 1 will be solution of homogeneous equation plus solution of particular solution, in this case. Now, what is the solution of homogeneous equation for this case? That we already have learned. So, said homogeneous in this case can be expressed as $A_1 \sin \omega_n t$ plus $A_2 \cos \omega_n t$.

And what will be the particular solution for this case? For this case, particular solution will be, you can see the equation that will be μN divided by k . Then we can write the solution Z is equal to $A_1 \sin \omega_n t$ plus $A_2 \cos \omega_n t$. Then what is left for us? In this case, it is plus, let me write it close to this, otherwise, yes, μN by k . This is for the case when motion of mass is occurring from right to left. Now, come to the case when motion of mass is occurring from left to right that means this case.

Then for this case, Z will be the same way Z homogeneous plus Z particular. And that means Z for homogeneous equation plus particular solution. Then we can write it here as $A_3 \sin \omega_n t$ plus $A_4 \cos \omega_n t$ minus μN divided by k . This is for the case when motion of mass from left to right. Now, we have two equations and two solutions, depending upon the direction of motion of mass from 0 to x_0 position.

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Single Degree of Freedom System with Coulomb Damping

Derivation:

Inertia force = $m\ddot{x}$

Unbalanced force = $-kx - \mu N$

$\Rightarrow m\ddot{x} = -kx - \mu N \Rightarrow m\ddot{x} + kx = -\mu N$ (L→R)

$m\ddot{x} + kx = \pm \mu N$

\Downarrow

$z = z_{hom} + z_p = A_3 \sin \omega_n t + A_4 \cos \omega_n t - \frac{\mu N}{k}$

Motion of mass m from left to right

From eq. I (0 to x_0)

So, here you can see the situation. So, for first half cycle, x 0 to 0 and then 0 to minus x 0, in the other direction. So, before coming here, now, we need to find out the solution. Let us go to the white board.

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$$m\ddot{x} + kx = \mu N \quad (R \rightarrow L) \quad x = A_1 \sin \omega_n t + A_2 \cos \omega_n t + \frac{\mu N}{k} \quad (2a)$$

$$\dot{x} = A_1 \omega_n \cos \omega_n t - A_2 \omega_n \sin \omega_n t$$

$$m\ddot{x} + kx = -\mu N \quad (L \rightarrow R) \quad x = A_3 \sin \omega_n t + A_4 \cos \omega_n t - \frac{\mu N}{k} \quad (3b)$$

For Eq. (3a) at $t=0$ $x = x_0$ & $\dot{x} = 0$

$$x_0 = A_1(0) + A_2(1) + \frac{\mu N}{k} \Rightarrow A_2 = x_0 - \frac{\mu N}{k}$$

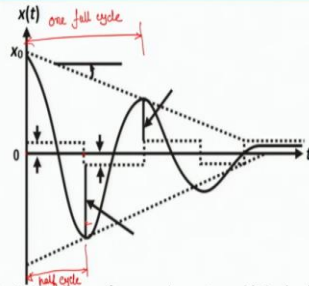
$$\dot{x}|_{t=0} = 0 = A_1 \omega_n(1) + A_2(-\omega_n)(0) \Rightarrow A_1 = 0$$

Eq. (3a) can be re-written as: $x = \left(x_0 - \frac{\mu N}{k}\right) \cos \omega_n t + \frac{\mu N}{k}$ (4a)

at $t = \frac{\pi}{\omega_n}$ $x = -\left(x_0 - \frac{\mu N}{k}\right) + \frac{\mu N}{k}$

$$= -\left(x_0 - \frac{2\mu N}{k}\right)$$

Single Degree of Freedom System with Coulomb Damping



Boundary conditions (for first half-cycle)

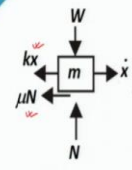
@ $t=0$ $x=x_0$ $\dot{x}=0$ (velocity)
 t varies from 0 to $\frac{\pi}{\omega_n}$

Fig. 11.3 Response of mass-spring system with Coulomb damping

Single Degree of Freedom System with Coulomb Damping

Derivation:

Inertia force = $m\ddot{x}$
 Unbalanced force = $-kx - \mu N$



Motion of mass m from left to right
 Friction \rightarrow (0 to x_0)

$$\Rightarrow m\ddot{x} = -kx - \mu N \Rightarrow m\ddot{x} + kx = -\mu N \quad (L \rightarrow R)$$

$$\boxed{m\ddot{x} + kx = \pm \mu N}$$

$$\chi = x_{hom} + x_p = A_3 \sin \omega_n t + A_4 \cos \omega_n t - \frac{\mu N}{k}$$



Single Degree of Freedom System with Coulomb Damping

x-direction

Derivation: $= m\ddot{x}$; Unbalanced force $= -kx + \mu N$

Inertia force = Unbalanced force $\Rightarrow m\ddot{x} = -kx + \mu N$

$\Rightarrow m\ddot{x} + kx = \mu N$ (Motion of mass m from 'R' to 'L') — (1)

$x = x_{hom} + x_p$

$x = A_1 \sin \omega t + A_2 \cos \omega t + \frac{\mu N}{k}$

EP. (Static Equilibrium)

for x L to R (+ve)

Motion of mass m from right to left

$x_{hom} = A_1 \sin \omega t + A_2 \cos \omega t$

$x_p = \frac{\mu N}{k}$

So, we have two equations of motions. I am writing here once again. First one is $m \ddot{x} + kx = \mu N$, for right to left. And solution is $A_1 \sin \omega t + A_2 \cos \omega t + \frac{\mu N}{k}$. The second equation is for motion left to right and that is $m \ddot{x} + kx = -\mu N$.

So, in this case, our general solution is Z . That is, $A_3 \sin \omega t + A_4 \cos \omega t + \frac{\mu N}{k}$. Now, let us see this slide. So, for the first half cycle, first half cycle means from where to where? Up to where? Because full cycle means from here to here. This is full cycle or I can write one full cycle. This is half cycle.

Now, for first half cycle, one boundary condition we can set here. You can see, at t equals to 0. That means at this point what is the value of x ? Yes, here it is x , not Z . Please, correct. I made a mistake here. This is x homogeneous. That means, solution for homogeneous equation plus particular solution. Same thing for the previous case also. Two three places mistakenly I have written Z which should be x , for this case.

So, here it is x , x . Same thing here also, it is x , x for homogeneous equation x particular solution. Here also I have corrected now. Now, come to the boundary condition. So, what we can see here? At t is equal to 0, what is a value of x ? x is x_0 . And if we will draw a slope here, that means \dot{x} , that will be how much? \dot{x} will be 0 here. And \dot{x} is nothing but the velocity.

So, now, we have two boundary conditions at t is equal to 0. One boundary conditions says x is equal to x_0 and the other boundary condition says velocity is 0 at t is equal to 0. Then what we can do here? Give me name for this one, let us take this is 3 a and this is 3 b. So, then we can write for equation or expression 3 a, at t is equal to 0, x is equal to x_0 and \dot{x} is 0.

So, first t is equal to 0, t is equal to 0, x is x_0 . So, I am writing equation 3 a once again, x_0 is equal to A_1 times, this is 0, because t is 0 plus A_2 times 1 plus μN by k . So, from this we can get A_2 is equal to x_0 minus μN by k . Likewise, if we will find out \dot{x} at t is equal to 0, which is 0 here, now, what is the expression for \dot{x} ?

The expression for \dot{x} is A_1 times ω_n times cosine $\omega_n t$. I can write here also. A_1 times ω_n times cosine $\omega_n t$ plus A_2 times minus ω_n times sin $\omega_n t$. This is the expression for velocity. Now, at t is equal to 0, velocity is 0. So, we can write then A_1 ω_n times 1 plus A_2 times minus ω_n times 0. Because sin $\omega_n t$ is 0 at t is equal to 0. Then from this we can get A_1 is 0.

That means, now, equation 3 a can be rewritten as x is equal to A_1 is 0, so, first term will not be there. A_2 is how much? x_0 minus μN by k . So, I am writing x_0 minus μ in by k . Let me write it better way, μN by k times cosine of $\omega_n t$ plus μN by k . So, this is our μ equation 4 a. Now, what about equation 3 b, what will be the boundary condition for equation 3 b? So, when I am talking about first half cycle that means t , from where to where?

t is varying from from 0 to π by ω_n . This is π by ω_n . This is the domain for the first half cycle. Now, for the second half cycle, at the beginning of second half cycle, that means when t is 0 for the second half cycle, what will be the value of displacement, this one? So, let us see what is the value of displacement?

So, here I can write at t is equal to π by ω_n , what will be x . So, cosine $\omega_n t$ at π t is equal to π by ω_n is nothing but cosine π . Cosine π means minus 1 plus μN by k . So, it is coming minus x_0 minus 2 μN by k . So, these boundary conditions now we can use for the second half cycle.

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Single Degree of Freedom System with Coulomb Damping

➤ Boundary conditions (for second half-cycle)

$$x = A_3 \sin \omega_n t + A_4 \cos \omega_n t - \frac{\mu N}{k} \quad (0 \leq t \leq \pi/\omega_n)$$

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Single Degree of Freedom System with Coulomb Damping

➤ Boundary conditions (for first half-cycle)

@ $t=0$ $x=x_0$ $\dot{x}=0$ (velocity)

t varies from 0 to $\frac{\pi}{\omega_n}$

Fig. 11.3 Response of mass-spring system with Coulomb damping

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That means here itself. Let me write this value here also, $x_0 - \frac{2\mu N}{k}$. Now, that means, for second half cycle, what we can write? For second half cycle, our equation was, x is equal to $A_3 \sin \omega_n t + A_4 \cos \omega_n t - \frac{\mu N}{k}$. This is the general solution where t varies from, once again, 0 to π by ω_n .

So, when I am writing t here, same thing, t varies from 0 to π by ω_n . That means, from here to here. This is second half cycle. So, here x at t is equal to 0, I need to write the expression. So, let us go back to the white board.

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For Second half cycle
 $x = -\left(x_0 - \frac{2\mu N}{K}\right)$ at $t=0$

$$-\left(x_0 - \frac{2\mu N}{K}\right) = A_3(0) + A_4(1) - \frac{\mu N}{K}$$


$$\Rightarrow A_4 = -\left(x_0 - \frac{3\mu N}{K}\right)$$

for $\dot{x} = 0$ at $t=0$ $A_3 = 0$

$$x = -\left(x_0 - \frac{3\mu N}{K}\right) \cos \omega_n t - \frac{\mu N}{K} \quad (4b)$$

at $t = \frac{\pi}{\omega_n}$ (for Second half cycle)

$$x = (-)(-1) \left(x_0 - \frac{3\mu N}{K}\right) - \frac{\mu N}{K}$$

$$= x_0 - \frac{4\mu N}{K}$$


$$m\ddot{x} + kx = \mu N \quad (R \rightarrow L) \quad \begin{cases} x = A_1 \sin \omega_n t + A_2 \cos \omega_n t + \frac{\mu N}{K} & -(3a) \\ \dot{x} = A_1 \omega_n \cos \omega_n t + A_2 (-\omega_n) \sin \omega_n t & \end{cases}$$

$$m\ddot{x} + kx = -\mu N \quad (L \rightarrow R) \quad \begin{cases} x = A_3 \sin \omega_n t + A_4 \cos \omega_n t - \frac{\mu N}{K} & -(3b) \end{cases}$$

For Eq. (3a) at $t=0$ $x = x_0$ & $\dot{x} = 0$

$$x_0 = A_1(0) + A_2(1) + \frac{\mu N}{K} \Rightarrow A_2 = x_0 - \frac{\mu N}{K}$$

$$\dot{x}|_{t=0} = 0 = A_1 \omega_n(1) + A_2(-\omega_n)(0) \Rightarrow A_1 = 0$$

Eq. (3a) can be re-written as: $x = \left(x_0 - \frac{\mu N}{K}\right) \cos \omega_n t + \frac{\mu N}{K}$ — (4a)

at $t = \frac{\pi}{\omega_n}$ $x = -\left(x_0 - \frac{\mu N}{K}\right) + \frac{\mu N}{K}$

$$= -\left(x_0 - \frac{2\mu N}{K}\right)$$

Single Degree of Freedom System with Coulomb Damping

► Boundary conditions (for second half-cycle)

$$x = A_3 \sin \omega_n t + A_4 \cos \omega_n t - \frac{\mu N}{k} \quad (0 \leq t \leq \pi/\omega_n)$$

$$x = -\left(x_0 - \frac{3\mu N}{k}\right) \cos \omega_n t - \frac{\mu N}{k}$$

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For second half cycle, x is equal to minus x_0 minus $2\pi N$ by k at t is equal to zero. Then what we can write? Minus x_0 minus $2\pi N$ by k time that is equal to A_3 times 0 plus A_4 times 1 which gives us A_4 is equal to, how much? $3\mu N$ by k . So, and for \dot{x} equals to 0 at t is equal to 0, we can get A_3 will be 0, just like the previous case. That means first half cycle.

So, finally, what we can see? The general solution for the second half cycle becomes A_4 which is 0 here. So, I am not writing $A_3 \sin \omega_n t$ here. I am directly writing $A_4 \cos \omega_n t$. So, the value of A_4 is a minus of x_0 minus $3\mu N$ by k times cosine $\omega_n t$ plus μN by k . It will not be plus, it will be minus. So, this is the general solution you can see here the first one.

Now, the second one that means second half cycle general solution is this one. So, we can give it the name 4 b. Then once again see here, at the end of the second half cycle that means at this point what will be the value of our x ? That means at t is equal to π by ω_n for second half cycle x will be minus, in this case cosine $\omega_n t$ at t is equal to π by ω_n is minus 1.

So, there are two minus times minus 1 times x_0 $3\mu N$ by k minus μN by k . So, finally, what we are getting is x_0 minus $4\mu N$ by k here. So, this is nothing but, let me write, this is x_0 minus $4\mu N$ divided by k . And these distance is μN by k , this one. So, in this way finally, we are getting the general solution which is, see it once again here, minus x_0 minus $3\mu N$ by k times cosine $\omega_n t$ minus μN by k .

So, that means, what we have seen today? At the end of the first full cycle, the amplitude reduced by an amount $4 \mu N$ divided by k . Because at the starting, amplitude was x_0 , and at the end of the first cycle, it is x_0 minus $4 \mu N$ divided by k which indicates that amplitude is reduced by an amount $4 \mu N$ divided by k .

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SUMMARY

In this lecture we discussed the followings:

- ✓ Spring-mass system with Coulomb damping
- ✓ In Coulomb damping: $x_n - x_{n-1} = \frac{4\mu N}{k}$
- ✓ Damped natural frequency (in case of Coulomb damping): $\omega_n = \sqrt{\frac{k}{m}}$

$\mu = \text{Coeff. of friction}$
 $N = \text{Normal force acting on 'm'}$

$m\ddot{x} + kx = \mu N$
 $\Rightarrow \ddot{x} + \frac{k}{m}x = \frac{\mu N}{m}$

So, come to then the summary of today's class. We learned spring mass system with coulomb damping. How we will form the equation of motions and from that how we will set the boundary conditions and from that what will be the solution. Then what we have seen? The difference between two consecutive amplitude in case of coulomb damping is $4 \mu N$ divided by k , where μ is coefficient of friction, N is normal force acting on mass m .

This is small m and k is the spring stiffness. Another interesting thing which we have started in this case is that in case of coulomb damping, the damped natural frequency remains same as the undamped natural frequency. That means ω_n is equal to square root of k by m , in this case.

That is the reason, actually, I have taken the general solution as $A_1 \sin \omega_n t$ plus $A_2 \cos \omega_n t$. If you see the equation of motion, equation of motion for one case, it was like this. So, that means, basically, it can be written like this. So, with this, now, I can conclude today's class.

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REFERENCES

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These are the references which I have used for this lecture. Thank you.