

NPTEL ONLINE CERTIFICATION COURSES

EARTHQUAKE SEISMOLOGY

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Module 03 : Surface Waves and Dispersion Lecture 04: Dispersion, Phase and group velocities

CONCEPTS COVERED

- > Dispersion
- > Phase and group velocities



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Dispersion

- Dispersion means that wave velocities are dependent on frequency.
- As an example, Snell's Law tells us that the index of refraction of light in a medium is dependent upon the speed of light in that medium.
- However, the speed of light in a medium (like water) is often frequency-dependent.
- Hence, in a raindrop, different colors will be refracted by different amounts, giving us a rainbow.
- More technically, dispersion implies $v=v(\lambda)$.



Another way often used to represent dispersion is to give the relationship between angular velocity ω and wave number k. We know already that (phase) velocity is a function of these two parameters:

In a dispersive medium, $\omega = \omega(k)$ For instance, short period tsunami waves follow $\omega = \sqrt{gk}$. We can solve for what this means in terms of velocity.

$$w=rac{\omega}{k}=rac{\sqrt{gk}}{k}$$

Often dispersive relations are given in the format $\,f(\omega,k)=0$, so for the tsunami example, dispersion would be represented as $\,\omega-\sqrt{gk}=0$.

 $v_{phase} = rac{1}{k}$



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To explore dispersion further, we first consider the simplest example, the net effect of two harmonic waves with slightly different frequencies and wavenumbers

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$$egin{array}{lll} \omega_1 &= \omega_0 + \delta \omega \ \omega_2 &= \omega_0 - \delta \omega \end{array} extbf{and}, extbf{and}, extbf{k}_1 &= k_0 + \delta k \ k_2 &= k_0 - \delta k \end{array}$$

The total displacement of the two waves is



$$egin{aligned} &u=\cos\left(k_1x-\omega_1t
ight)+\cos\left(k_2x-\omega_2t
ight)\ &=\cos\left((k+\delta k)x-(\omega+\delta\omega)t
ight)+\cos\left((k-\delta k)x-(\omega-\delta\omega)t
ight)\ &=\cos\left[(kx-\omega t)+(\delta kx+\delta\omega t)
ight]+\cos\left[(kx-\omega t)-(\delta kx-\delta\omega t)
ight] \end{aligned}$$



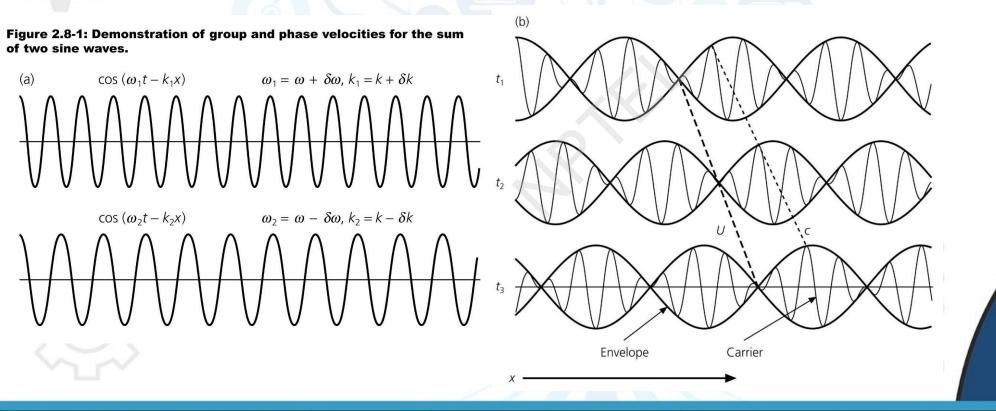
Using the trig identity $\cos{(heta+\phi)}+\cos{(heta-\phi)}=2\cos{ heta}\cos{\phi}$

the displacement becomes $u=2\cos{(kx-\omega t)}\cos{(\delta kx-\delta \omega t)}$

This is the product of two cosine waves. The first wave has wavelength, $\lambda = \frac{2\pi}{k}$ and assuming $\delta k << k$ that , the second will have a much larger wavelength defined by $\lambda = \frac{2\pi}{\delta k}$



Alternatively, in the time domain, the first wave will have a higher frequency $f=rac{\omega}{2\pi}$, compared to $f=rac{\delta\omega}{2\pi}.$





- Hence, the first cosine term defines a wave (sometimes called the "carrier") that is modulated by a slower varying "envelope".
- Since, the velocity with which a wave defined by $u=\cos{(ax-bt)}$ propagates is v = b/a
- The velocity of the carrier wave is v = ω/k and that of the envelope is slower $v = \frac{\delta w}{\delta k}$
- We refer to the former as "phase velocity", and the latter as the "group velocity".
- Typically, "c" is used to denote phase velocity, and U to denote group velocity.
- The relationship between $c = v_{phase} = \frac{\omega}{k}$ and $U = v_{group} = \frac{\delta\omega}{\delta k}$ will be defined by the dispersion relationship $\omega = \omega(k)$. As $\delta\omega, \delta k \to 0 \qquad U \to \frac{\mathrm{d}\omega}{\mathrm{d}k}$



What this means is that the most immediately recognizable wave (or a surface of constant phase) will propagate with a different velocity than the envelope. It turns out the actual energy propagation velocity is defined by the group velocity. In a medium with anisotropic velocity structure, the group and phase velocities actual may have slightly different directions.

Since,

$$c=rac{\omega(k)}{k}$$

The relationship between c and U at a given angular frequency ω can be derived using the chain rule.

$$U=rac{\mathrm{d}\omega}{\mathrm{d}k}=rac{\mathrm{d}}{\mathrm{d}k}(ck\,)=krac{\mathrm{d}c}{\mathrm{d}k}+crac{\mathrm{d}k}{\mathrm{d}k}=krac{\mathrm{d}c}{\mathrm{d}k}+c$$

In terms of wavelength, $\lambda=rac{2\pi}{k}\,$, can be used to derive the group velocity equation.



Using,

$$egin{aligned} rac{\mathrm{d}c}{\mathrm{d}k} &= rac{\mathrm{d}c}{\mathrm{d}\lambda}rac{\mathrm{d}\lambda}{\mathrm{d}k} \ &= rac{\mathrm{d}c}{\mathrm{d}\lambda}igg(rac{-2\pi}{k^2}igg) \end{aligned}$$

We can convert this to wavelength

$$\begin{aligned} k\frac{\mathrm{d}c}{\mathrm{d}k} + c &= k\frac{\mathrm{d}c}{\mathrm{d}\lambda}\frac{\mathrm{d}\lambda}{\mathrm{d}k} + c \\ &= k\frac{\mathrm{d}c}{\mathrm{d}\lambda}\left(\frac{-2\pi}{k^2}\right) + c \\ &= \frac{\mathrm{d}c}{\mathrm{d}\lambda}\left(\frac{-2\pi}{k}\right) + c \\ &= c - (2\pi)\lambda\frac{\mathrm{d}c}{\mathrm{d}\lambda} \end{aligned}$$





In practice, the 2π is typically dropped to get

Since at greater depths in the earth tend to have higher velocities, c increases with wavelength, so

$$rac{{
m d}c}{{
m d}\lambda} >$$

 $U=c-\lambdarac{\mathrm{d}c}{\mathrm{d}\lambda}$

So from above relation, on average in the earth, c >U.



If the medium is non-dispersive

and the phase and group velocity will be the same.

$$rac{\mathrm{d}c}{\mathrm{d}\lambda} = rac{\mathrm{d}c}{\mathrm{d}k} = 0$$

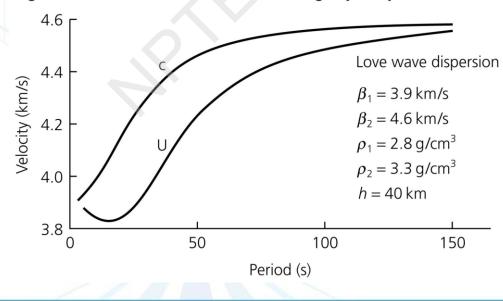


Figure 2.8-2: Fundamental mode Love wave group and phase velocities.



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