



NPTTEL ONLINE CERTIFICATION COURSES

EARTHQUAKE SEISMOLOGY

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Module 05 : Refraction and Reflection seismology

Lecture 04: Earth as a constant velocity Distribution and Intercept-slowness formulation for travel times

CONCEPTS COVERED

- **Recap**
- **Earth as a constant velocity Distribution**
- **Intercept-slowness formulation for travel times**
- **Summary**

Recap

- Study which use reflected arrivals to determine velocities within the crust, and other geophysical parameters is known as Reflection seismology.

- Travel time curve for reflected arrival is hyperbola , for one layer $T(x)^2 = \frac{x^2}{v_0^2} + t_0^2$

- Travel time curve for n layer geometry $T(x)_{n+1}^2 = \frac{x^2}{\bar{V}_n^2} + t_n^2$,

- Relation between ray path and travel time curve $p = \frac{\sin i}{v} = \frac{dT}{dx}$ (Ray parameter)

- Root mean square and interval velocity $\bar{V}_n^2 = \left(\sum_{j=0}^n v_j^2 \Delta t_j \right) / \sum_{j=0}^n \Delta t_j$

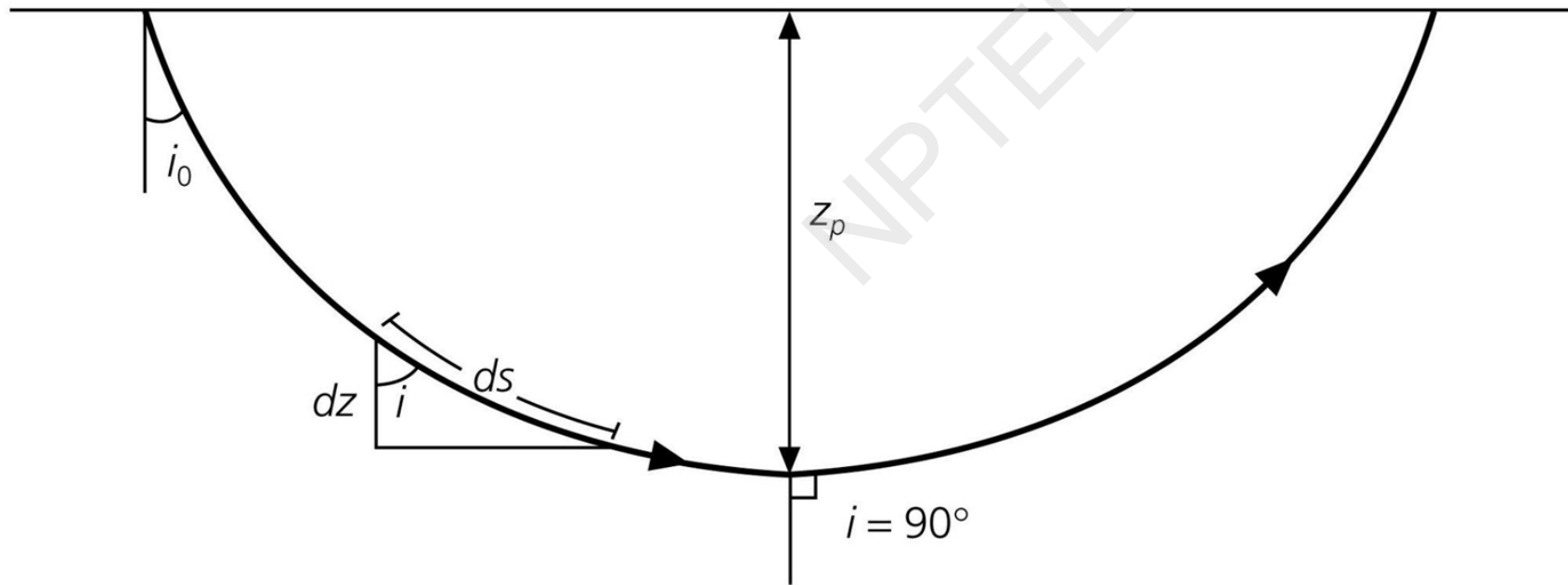
$$v_n^2 = \frac{\bar{V}_n^2 t_n - \bar{V}_{n-1}^2 t_{n-1}}{t_n - t_{n-1}}$$

Earth as continuous velocity distribution

Earth is spherical and it has continuous distribution of velocity with depth, $v(z)$, rather than a stack of discrete layers.

The ray path shown in Figure, and given by Snell's law, $p = \sin i/v(z)$, is constant.

Figure 3.3-6: Ray path in a medium with smoothly increasing velocities.



Earth as continuous velocity distribution

- If velocity increases with depth, $\sin i$ and thus i increase.
- So the ray bends away from the vertical on its way down.
- Once $i = 90^\circ$, the ray turns, becomes horizontal, and then goes upward.
- At the deepest point, the turning, or bottoming, depth z_p the velocity is the reciprocal of the ray parameter, $p = 1/v(z_p)$.
- The expression for the distance traveled by the ray becomes:

$$x(p) = 2 \int_0^{z_p} \tan i \, dz = 2p \int_0^{z_p} \left(\frac{1}{v^2(z)} - p^2 \right) dz \quad \sin i = pv$$

Note: “If on some portion of the ray path the velocity decreases with depth, the ray bends toward the vertical. The ray does not turn upward until it gets below the low-velocity region.”

Earth as continuous velocity distribution

In terms of the slowness, as $u(z) = 1/v(z)$

$$x(p) = 2p \int_0^{z_p} \frac{dz}{(u^2(z) - p^2)^{1/2}}$$

Similarly, total travel time becomes

$$\begin{aligned} T(p) &= 2 \int_0^{z_p} \frac{dz}{v(z) \cos i} = 2 \int_0^{z_p} \frac{dz}{v(z)(1 - p^2 v^2(z))^{1/2}} \\ &= 2 \int_0^{z_p} \frac{u^2(z) dz}{(u^2(z) - p^2) \cos i} \end{aligned}$$

This integral is valid everywhere except at the exact bottom of the curve, where $u(z)$ equals p .

Intercept-slowness formulation for travel times

An alternative formulation in terms of intercept-slowness may provide interesting insights for data analysis.

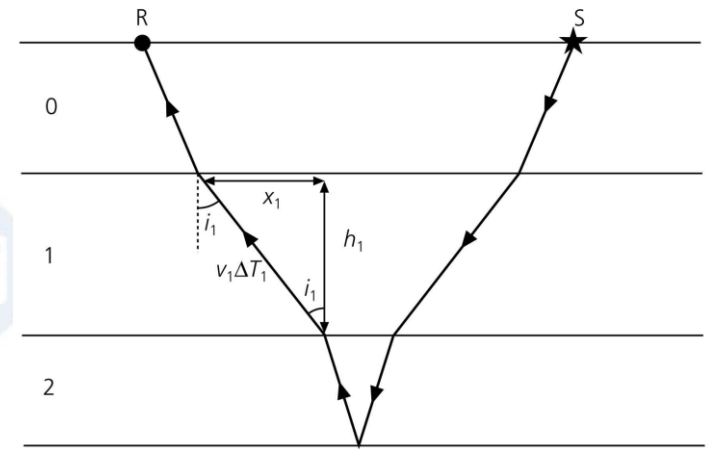
From the figure,

$$v_j \Delta T_j = (x_j^2 + h_j^2)^{1/2}$$

The incident angle i_j for this ray satisfies

$$\sin i_j = \frac{x_j}{(x_j^2 + h_j^2)^{1/2}} = \frac{x_j}{v_j \Delta T_j} \quad \text{and} \quad \cos i_j = \frac{h_j}{(x_j^2 + h_j^2)^{1/2}} = \frac{h_j}{v_j \Delta T_j}$$

Figure 3.3-3: Ray path through multilayered structure.



Hence,

$$v_j \Delta T_j = \frac{x_j^2 + h_j^2}{(x_j^2 + h_j^2)^{1/2}} = x_j \sin i_j + h_j \cos i_j$$

or

$$\Delta T_j = \frac{x_j \sin i_j}{v_j} + \frac{h_j \cos i_j}{v_j} = p_j x_j + \eta_j h_j$$

The travel time a ray spends in a layer is the sum of the horizontal slowness times the horizontal distance traveled and the vertical slowness times the vertical thickness.

$$p_j = \frac{\sin i_j}{v_j} = u_j \sin i_j \quad \text{and} \quad \eta_j = \frac{\cos i_j}{v_j} = u_j \cos i_j \quad \dots\text{equ. 1}$$

These are the components of the slowness vector that has magnitude equal to the slowness, and points in the direction of wave propagation. Hence u_j , the slowness in the layer, is

$$u_j^2 = \frac{1}{v_j^2} = p_j^2 + \eta_j^2$$

Intercept-slowness formulation for travel times

The total travel time is:

$$T(x) = 2 \sum_{j=0}^n \Delta T_j = 2 \sum_{j=0}^n p_j x_j + 2 \sum_{j=0}^n \eta_j h_j$$

By Snell's law, the horizontal ray parameter is constant along the ray path, so $p_j = p$, and

$$T(x) = px + 2 \sum_{j=0}^n \eta_j h_j$$

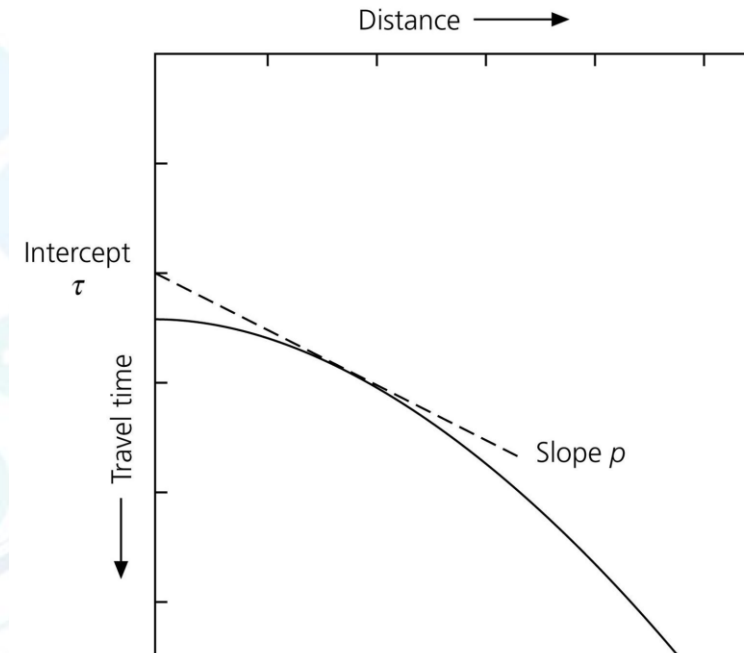
We define, $T(x) = px + \tau(p)$ equ 2

where the function,

$$\tau(p) = 2 \sum_{j=0}^n \eta_j h_j = 2 \sum_{j=0}^n \left(\frac{1}{v_j^2} - p^2 \right)^{1/2} h_j = 2 \sum_{j=0}^n (u_j^2 - p^2)^{1/2} h_j \text{equ 3}$$

- Because p is the slope of the travel time curve (dT/dx) and hence of a line tangential to it at the point (T, x) , τ is the intercept of the tangent line with the time axis.
- In general τ and p differ for different points on the travel time curve, so the travel time curve can be described by the values of either (T, x) or (τ, p) .
- “Thus the function $\tau(p)$ is called the *intercept-slowness* representation of the travel time curve.”

Figure 3.3-7: Relation between travel time curve, tau, and ray parameter.



On differentiating equ.2 w.r.t. p

$$T(x) = px + \tau(p) \quad \text{.....equ 2}$$

$$\frac{d\tau}{dp} = \frac{dT}{dp} - p \frac{dx}{dp} - x(p) = \frac{dT}{dx} \frac{dx}{dp} - p \frac{dx}{dp} - x(p) = -x(p)$$

“Thus, just as p is the slope of the travel time curve, $T(x)$, the distance, x , is minus the slope of the $\tau(p)$ curve.”

Intercept-slowness formulation for travel times

We may also show that the $\tau(p)$ formulation will lead to the travel time curve for the reflected wave in a layer over a halfspace. For this case, $x_0 = x/2$, so, using Eqn 1,

$$p = \frac{x/2}{v_o \left[(x/2)^2 + h_o^2 \right]^{1/2}}, \quad \eta_o = \frac{h_o}{v_o \left[(x/2)^2 + h_o^2 \right]^{1/2}}$$

Hence,

$$\begin{aligned} T(x) &= px + 2\eta_o h_o = \frac{(x^2/2) + 2h_o^2}{v_o \left[(x/2)^2 + h_o^2 \right]^{1/2}} \\ &= \frac{\left((x/2)^2 + h_o^2 \right)^{1/2}}{v_o} \end{aligned}$$

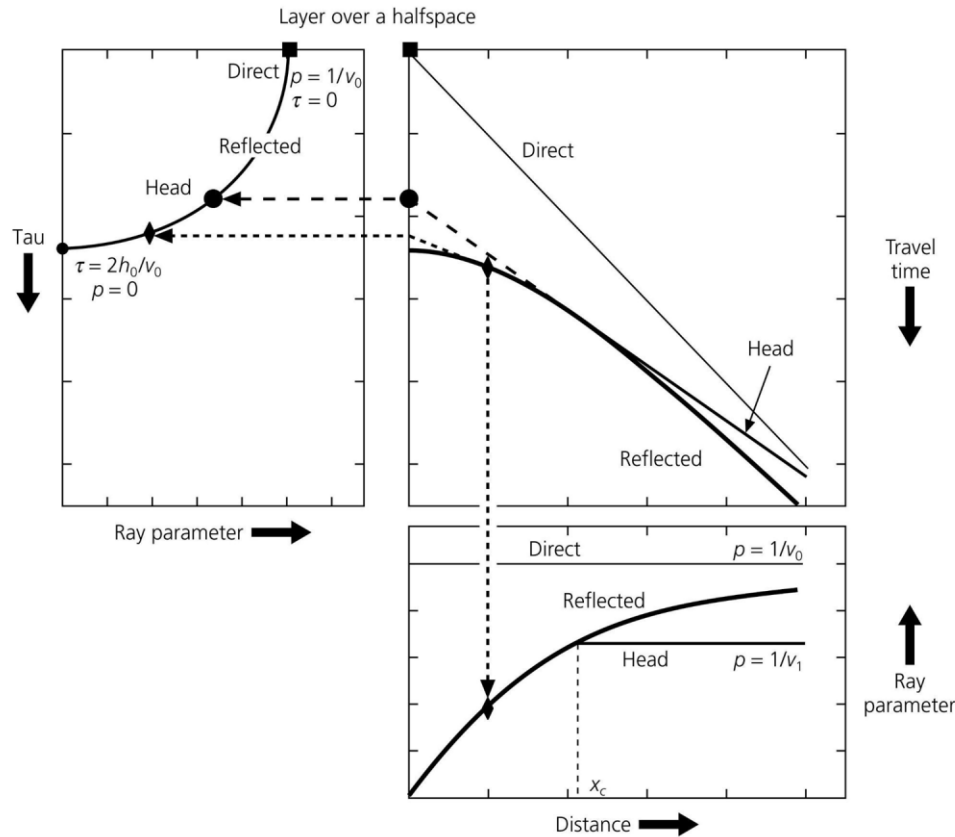
which is the familiar hyperbola. To see how this travel time curve appears when written as $\tau(p)$, we write Eqn 3 for a layer over a halfspace:

$$\tau(p) = 2 \left(\frac{1}{v_o^2} - p^2 \right)^{1/2} h_o \quad \text{.....equ 4}$$

Intercept-slowness formulation for travel times

This can also be written as: $\frac{v_0^2 \tau^2}{4h_0^2} + v_0^2 p^2 = 1$ ellipse whose axes are the τ and p axes

Figure 3.3-8: Relation between tau-p, and travel time curves.



→ It intersects the τ axis at ($\tau = t_0 = 2h_0/v_0$ at $p = 0$), which show the travel time curve has zero slope and the time axis intercept is the vertical two-way travel time, corresponds to the zero-offset point $x = 0$.

→ Ellipse intersect p axis at ($\tau = 0, p = 1/v_0$), where the travel time curve has slope $1/v_0$ and time axis intercept 0, is the $\tau(p)$ position of the linear travel time curve for the direct wave.

Intercept-slowness formulation for travel times

Hence the line for the direct wave maps to a point in the $\tau(p)$ plane that is on the ellipse describing the reflected wave.

To understand why this occurs, we use the fact that distance is minus the derivative of the $\tau(p)$ curve and differentiate Eqn 4, giving

$$x(p) = -\frac{d\tau}{dp} = 2ph_o \left(\frac{1}{v_o^2} - p^2 \right)^{-1/2} \quad \text{.....equ 5}$$

so at the point $p = 1/v_o$, $x = \infty$. This makes sense, because as $x \rightarrow \infty$, the reflected wave is asymptotic to the direct wave.

The head wave is easily mapped into the $\tau(p)$ plane, because its travel time curve is:

$$\begin{aligned} T_H(x) &= x/v_1 + 2h_o \left(1/v_o^2 - 1/v_1^2 \right)^{1/2} \\ &= x/v_1 + \tau_1 \end{aligned}$$

a line with slope equal to the reciprocal of the halfspace velocity, $p = 1/v_1$ and intercept τ_1 .

Intercept-slowness formulation for travel times

Thus the head wave maps into a point on the ellipse describing the reflected wave, corresponding to the critical distance x_c where the head and reflected waves are the same.

To see this, note that for $p = 1/v_1$, Eqn 5 gives:

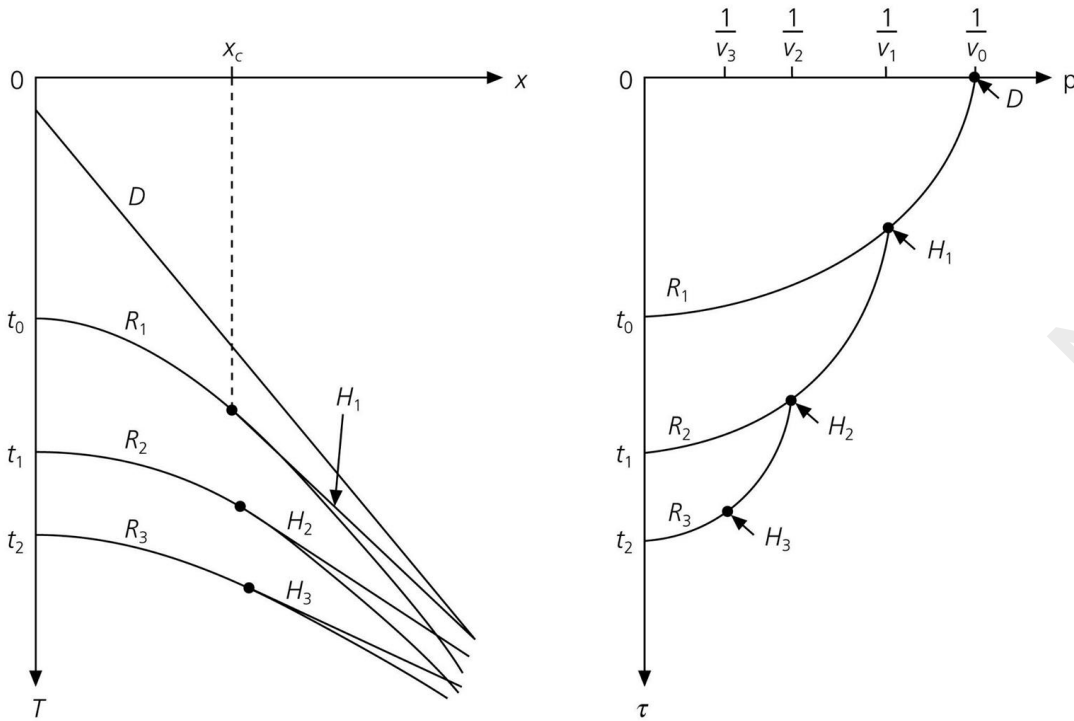
$$x(p) = -\frac{d\tau}{dp} = 2h_o v_o (v_1^2 - v_o^2)^{-1/2} = x_c$$

This point divides the ellipse describing the reflected wave into a subcritical portion, between the τ axis and the head wave, and a post critical portion, between the head wave and the p axis.

Intercept-slowness formulation for travel times

For multiple layers, the $\tau(p)$ curves corresponding to reflections off successive layers are all portions of different ellipses.

Figure 3.3-9: Tau-P and travel time curves for multiple layers.



For a continuous velocity distribution, the summation for τ (Eqn 3) becomes an integral

$$\begin{aligned} \tau(p) &= 2 \int_0^{z_p} \eta(z) dz = 2 \int_0^{z_p} \left(\frac{1}{v^2(z)} - p^2 \right)^{1/2} dz \\ &= 2 \int_0^{z_p} (u^2(z) - p^2)^{1/2} dz \end{aligned}$$

Formulating travel time curves as $\tau(p)$ is useful for some techniques that invert for velocity structure.

Summary

- For a continuous increasing velocity model, at the deepest point, the turning, or bottoming, depth z_p , the velocity is the reciprocal of the ray parameter, $p = 1/v(z_p)$.

- Distance travelled by the ray:
$$x(p) = 2 \int_0^{z_p} \tan i \, dz = 2p \int_0^{z_p} \left(\frac{1}{v^2(z)} - p^2 \right) dz$$

- Time taken to cover distance $x(p)$ by the ray:
$$\Delta T_j = \frac{x_j \sin i_j}{v_j} + \frac{h_j \sin i_j}{v_j} = p_j x_j + \eta_j h_j$$

- The function $\tau(p)$ is called the *intercept-slowness* representation of the travel time curve.
- p is the slope of the travel time curve, $T(x)$, the distance, x , is minus the slope of the $\tau(p)$ curve.
- The head wave maps into a point and hyperbola maps as ellipse on τ - p plane

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**THANK
YOU!**