

NPTEL ONLINE CERTIFICATION COURSES

EARTHQUAKE SEISMOLOGY

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Module 08 : Composition of the mantle and core Lecture 01: Density within the Earth, Adams–Williamson equation and Pressure within the earth

CONCEPTS COVERED

- Density within the earth
 - > Adams–Williamson equation
 - Setbacks
 - Solution
- > Pressure Variation within the earth
- > Summary



- → Density of Earth can be combined with velocities to derive elastic constants.
- → Average Earth's density can be found by the Earth's mass M, which can be found from the law of gravitation,
 where,

$$g = GM/a^2$$

 $g = 9.8 \text{ m/s}^2, a = 6371 \text{ km},$
 $G = 6.67 \times 10^{-11} \text{ Nm}^2 \text{kg}^{-2},$
 $M = 5.97 \times 10^{24} \text{ kg}.$

Note: The mass is the volume integral of the density because Earth's density varies only with depth \int^a

$$M=4\pi\int_0^{+}
ho(r)r^2dr$$

But the average density, ρ_o , is found by dividing the mass by the volume,

$$ho_o=M/ig[(4/3)\pi a^3ig]$$

The resulting average density of the earth is about 5.5 g/cm³



Spin axis

dv

 $\rho(r, \theta, \phi)$

 \rightarrow The fact that this value is significantly higher than the density of the surface rocks (about 3 g/cm³) is evidence for a core of much denser, and different material.

The second indication of the dense core, comes from the moment of Figure 3.8-1: Moment of inertia of a volume element. inertia about the rotation axis.

> This is defined by integrating over volumes dV, each at a distance $I = r \sin \theta$ from the spin axis,

$$egin{aligned} C &= \iiint l^2
ho(r, heta,\phi) dV = \int_0^{2\pi} \int_0^\pi \int_0^a
ho(r) ig(r^2 \sin^2 heta ig) r^2 \sin heta dr d heta d\phi \ &= rac{8}{3} \pi \int_0^a
ho(r) r^4 dr \end{aligned}$$



The ratio of the moment of inertia to the mass gives a scalar that depends on the density distribution.

Case I- If the earth were homogeneous, the density everywhere would equal the average density, $\rho(r) = \rho_0$, and

 $C = (8/15)\pi a^5 \rho_0$, $M = (4/3)\pi a^3 \rho_0$, $C/Ma^2 = 0.4$

Case II - if all the mass were in a shell at the surface.

$$ho(r)\,=\,\delta(r-a)
ho_s$$

 $C = (8/3)\pi a^4 \rho_s$, $M = 4\pi a^2 \rho_s$, $C/Ma^2 = 0.67$

i.e. a distribution with material concentrated towards the outside gives a larger ratio.



Case III -A more realistic case is a two-shell planet, with a mantle of density ρ_m and a core of density ρ_c and radius r_c . The integrals are evaluated in pieces as

$$egin{aligned} C &= rac{8}{3} \pi igg[\int_{0}^{r}
ho_{c} r^{4} dr + \int_{r_{c}}^{a}
ho_{m} r^{4} dr igg] = rac{8}{15} \pi igg[
ho_{m} a^{5} + (
ho_{c} -
ho_{m}) r_{c}^{5} igg] \ M &= rac{4}{3} \pi igg[
ho_{m} a^{3} + (
ho_{c} -
ho_{m}) r_{c}^{3} igg] \end{aligned}$$

If we consider the values similar to those for the earth i.e.,

 $ho_{c} = 12 \text{ g/cm}^{3}$, ho_{m} = 5 \text{ g/cm}^{3} ho_{c} = 3480 km,



- → It yields a moment of inertia ratio of C/Ma² = 0.35.
- → This value is less than the 0.4 which a uniform planet would have, because the material is concentrated toward the center.
- → It is similar to the value of C/Ma² for the earth field, about 0.33, determined from the earth's shape and gravity. Thus indicates the presence of a dense core.





- → Let's look at the variation of density with depth relating the seismic velocities. Consider a region of uniform material and see how the density increases with depth as the material is self-compressed by its own weight.
- \rightarrow At a radius r, the gradient of the hydrostatic pressure P(r) is:

 $\frac{dP}{dr} = -g\rho$ where $\rho(r)$ and g(r) are the density and the acceleration of gravity at that depth.

.....Equ.1

Since,

hence,

We know that,

 $g=GM/r^2$

 $rac{dP}{dr} = rac{ho Gm}{r^2}
onumber
ho = m/V, \ \ d heta = dV/V$

$$d
ho = -ig(m/V^2ig) dV = -
ho d heta$$



So,

The bulk modulus K can be expressed as:

$$K = -\frac{dP}{d\theta} = -\frac{dP}{d\rho}\frac{d\rho}{d\theta} = \rho\frac{dP}{d\rho}$$

Combining this with the pressure derivative equation 1

$$\frac{d\rho}{dr} = -\frac{d\rho}{dP}\frac{dP}{dr} = -\frac{\rho^2 Gm}{Kr^2}$$

To include the observations of seismic velocities, we define the seismic parameter, Φ , and bulk sound speed, $\Phi^{1/2}$, such that

$$\Phi=lpha^2-(4/3)eta^2=K/
ho$$

Thus we can write the Adams–Williamson equation relating the velocity structure to the derivative of density with radius,

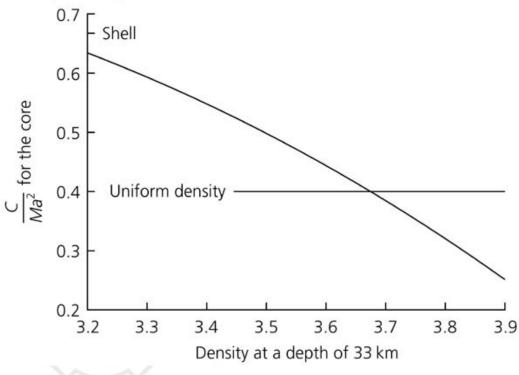
$$rac{d
ho(r)}{dr}=rac{-
ho(r)Gm(r)}{\Phi(r)r^2}=rac{-
ho(r)g(r)}{\Phi(r)}$$



Density within the earth : Setbacks

→ Adams- Williamson equation is insufficient because density increases with depth as a result of mineral phase changes as well as of self-compression.

Figure 3.8-2: Moment of inertia ratio of the earth's core as a function of density at the top of the mantle.



- → Figure shows the C/Ma² value calculated for the core as a function of the assumed density at the top of the mantle, which is the initial density for the Adams– Williamson calculation.
- → If we consider near-surface density ≈ 3.3 g/cm³, this will give C/Ma² for core greater than 0.4, implying that density decreases with depth in the core.
- → It is not true, because the solid inner core should be denser than the liquid outer core.



Density within the earth : Solution!!

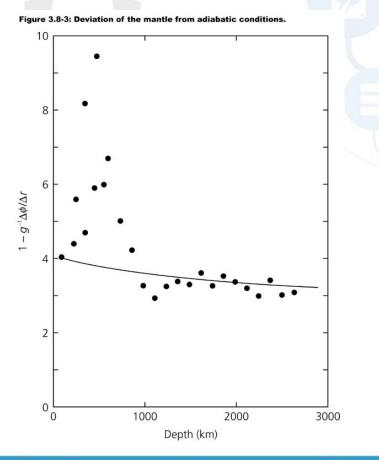
- → F. Birch (1950s) showed that at least one of two assumptions underlying the method was inappropriate.
- → One implicit assumption is that the temperature increases with depth along an adiabatic gradient, or "adiabat," which is not always true.
- → Infact, the temperature gradient in the mantle is thought to exceed the adiabatic gradient, because a superadiabatic gradient is required for the thermal convection expected in the mantle.
- → The superadiabatic gradient gap be incorporate in the Adams–Williamson equation to: $\frac{dp}{dr} = -\frac{pg}{\phi} + g\tau\alpha$

where, α is the coefficient of thermal expansion and τ is the portion of the temperature gradient exceeding the adiabetic gradient.



Density within the earth : Solution!!

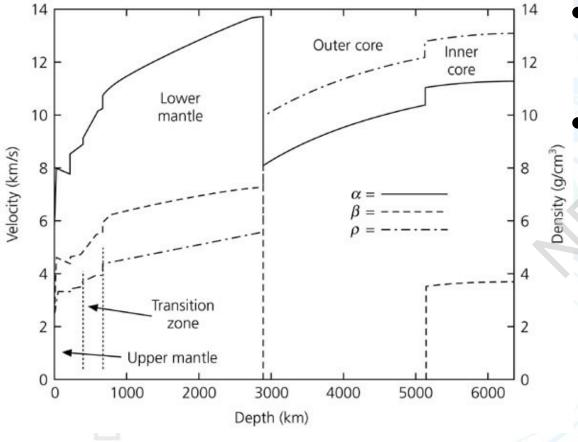
→ This correction for higher temperature lowers the calculated mantle densities and increase the C/Ma² for the core. So, it makes the issue of core density structure worse.



- → Hence the assumption of homogeneous material whose density changes only by self-compression must be incorrect. Birch showed that inhomogeneity can be identified using the function 1 – (1/g)dφ/dr.
- → Below 1000 km the mantle behaves as a homogeneous material, while at shallower depths it does not.
- → This is because the mineral phase transitions expected at the 410 and 660 km discontinuities involve denser atomic packings, and therefore transitions to higher densities, than predicted by the Adams–Williamson equation.



Figure 3.8-4: Preliminary Reference Earth Model.



- As a result, density models of the earth include rapid changes in the transition zone.
 - Within the lower mantle, outer core, and inner core, density increases smoothly with depth according to the Adams–Williamson equation.



Pressure variation within the Earth

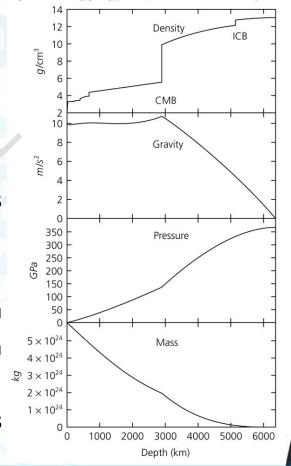
A density profile lets us compute a pressure profile, and thus use the results of experiments showing which mineral phases exist at particular pressures.

 $P(r)=-\int_{0}^{r}g(r)
ho(r)dr$

It can be computed as:

Points to remember

- Gravity averages 9.8 m/s² at earth's surface, and is zero at earth's center.
- Gravity increases slightly across the mantle, reaching a maximum of 10.7 m/s² at the CMB because of the high density of the core relative to the mantle.
- The core has only 16% of the earth's volume, but has almost one-third of the mass.





Summary

- The ratio of the moment of inertia to the mass gives a scalar that depends on the density distribution.
- Adams–Williamson equation relating the velocity structure to the derivative of density with radius,

$$rac{d
ho(r)}{dr}=rac{-
ho(r)Gm(r)}{\Phi(r)r^2}=rac{-
ho(r)g(r)}{\Phi(r)}=$$

- The drawback of Adams– Williamson equation is that, it does not take the account of changes in the mineral phase with depth.
- Modified Adams–Williamson equation which includes superadiabatic gradient is: α : coefficient of thermal expansion

- $rac{d
 ho}{dr}=-rac{
 ho g}{\phi}+g aulpha$ 7: portion of the temperature gradient exceeding the adiabatic gradient.
- Inhomogeneity in the earth can be identified using the function $1 (1/g)d\phi/dr$.
- Pressure inside the earth can be computed as:

$$\mathcal{P}(r) = -\int_0^r g(r)
ho(r) dr$$



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