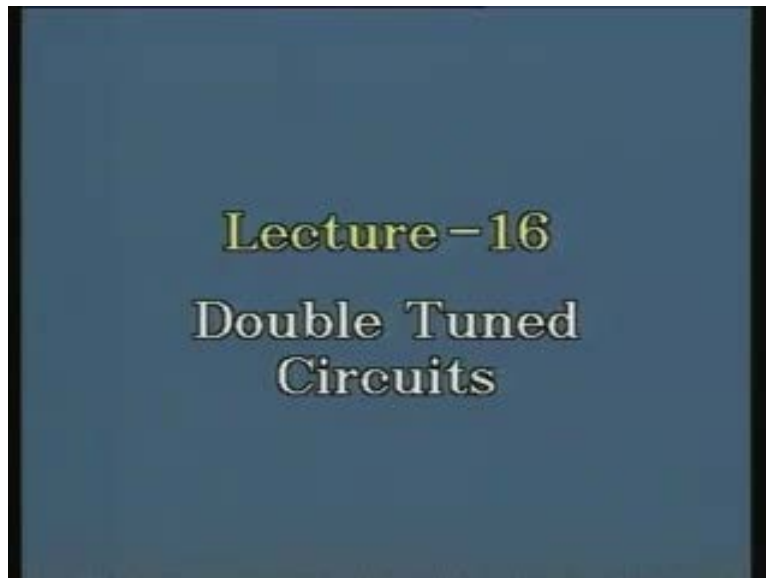


Circuit Theory
Prof. S.C. Dutta Roy
Department of Electrical Engineering
Indian Institute of Technology, Delhi

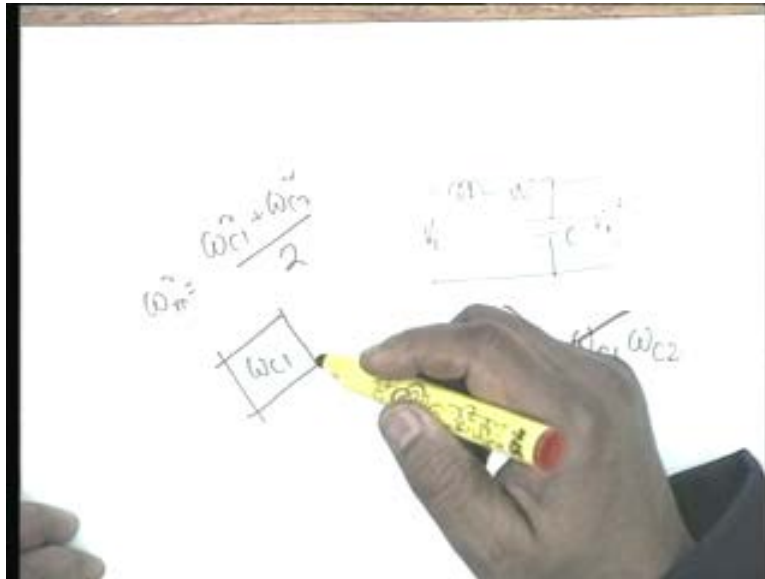
Lecture - 16
Double Tuned Circuits

(Refer Slide Time: 00:52)



This is the sixteenth lecture and our topic for discussion today is double tuned circuits. Before I taken double tuned circuits, I like to mention that in response to the challenges that I have said last time and I had said, not all statements that I make are true. I am not under oath to make all true statements. I do make intentional mistakes and one of the things that mean is pointed out, is this relationship.

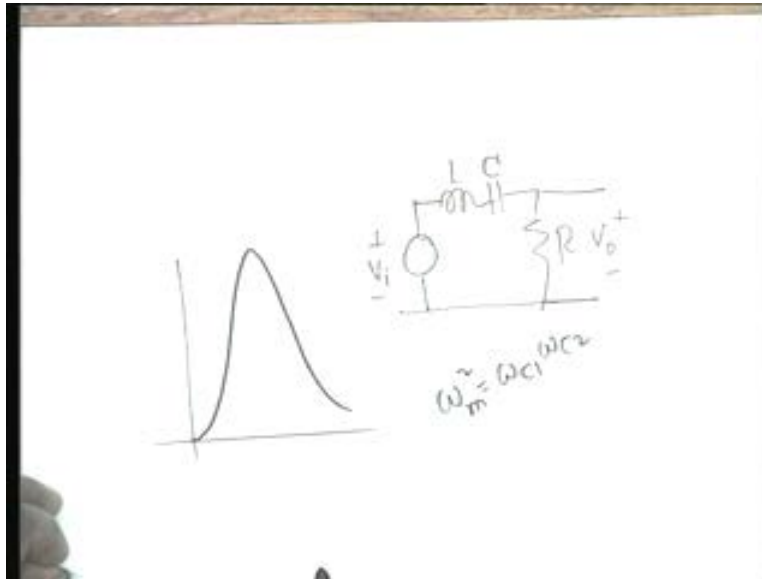
(Refer Slide Time: 1:32)



In a single tuned circuit, if we take the output across a capacitor c and this is $V I$, this is $V 0$, then the central frequency or the frequency of maximum, I had claimed that this will be the geometric mean of the 2 cut off frequencies, $\omega_c 1$ and $\omega_c 2$. This happens to be a wrong statement. In fact, what is true is that ω_m squared is equal to $\omega_c 1$ squared plus $\omega_c 2$ squared divided by 2. That is, it is the arithmetic mean of, not of the half power frequencies, but of the square of the frequency of maximum response is the arithmetic mean of these squares of the half power frequencies.

That is number , and number 2 I had set another challenge which has been very successfully met is how to get $\omega_c 1$. $\omega_c 2$, we had found out, both the reports, Mayank and Manish, both the reports, I got only 2, have shown how to construct this ω_c , geometrically, and I am very happy about it. May I also mention, in passing, that the statements that I have made is needs true. Once again, with a pinch of salt, is indeed true if you take the output across the resistor.

(Refer Slide Time: 3:09)



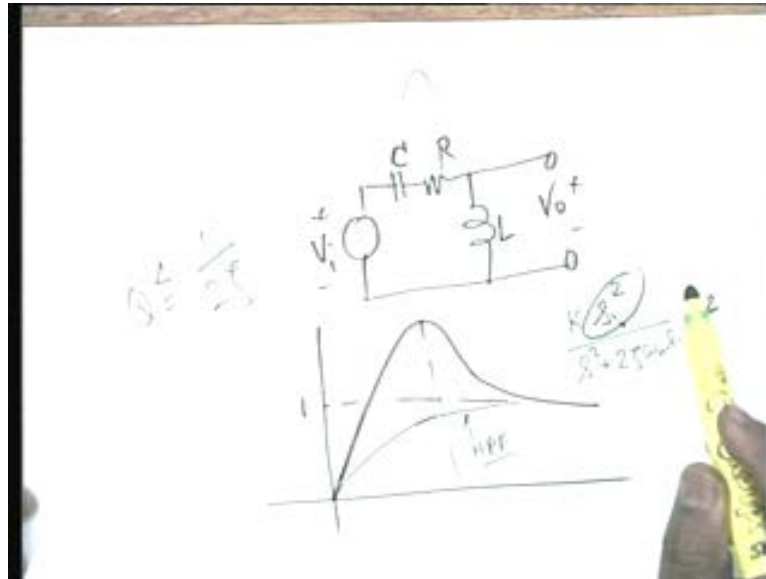
That is, if the output voltage is proportional to the current in this circuit. Suppose, I have a circuit like this, then you can see from physical reasoning that at DC, at a DC, the output shall be 0, because the capacitor will charge to the DC voltage. At infinite frequency, the output shall again be 0. Why, because the inductor offers an open circuit. $j\omega L$, ω goes to infinity therefore, the impedance goes to infinity and therefore, the response of this circuit shall truly be a band pass response. It would be like this. It would start from 0, it would end up in 0. Truly a band pass response and one can easily so that here ω_m^2 is indeed equal to $\omega c_1 \omega c_2$. It is the geometric mean.

Now you notice that same circuit, same tuned circuit, can give you a low pass response and also a band pass response. A low pass response flows if you take the output across the capacitor and ζ is less than, ζ is a greater than or less than

Student: Greater than.

Sir: Greater than 1 by, greater than the critical balance. Similarly, the same tune circuit, if you take the output across the inductor,

(Refer Slide Time: 4:48)



if you take the output across the inductor, then you see that, again from physical arguments, at DC the output should be 0, because the inductor axis is short circuit. At infinite frequency, the output should be the same as the input, that is, the transfer function magnitude will be equal to 1 and therefore, in between, in between there are maybe there may be an maximum, there may not be a maximum. For example, if zeta is less than a certain value, I will get a response like this. Were the finite, the final, the infinite frequency value is unity, starts from 0 at DC, then it may so a peaking, depending on the value of zeta and then it goes (...) to unity at infinite frequency.

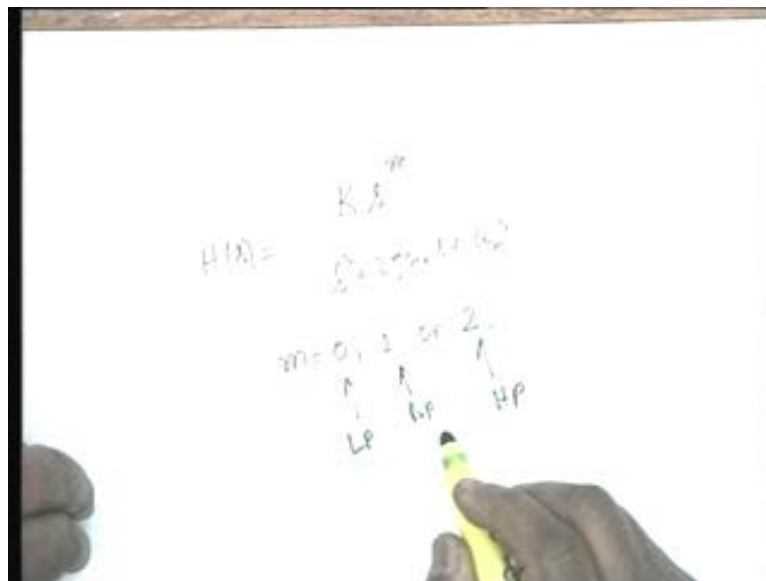
On the other hand, if zeta is greater than a certain value, that if zeta exits the critical damping, you might get a response like this, which is truly a high pass response, a high pass filter. So the single tuned circuit it an extremely interesting circuit from which you can get high pass, low pass, as well as band pass responses. You see, even in the high pass mode, it will act as a band pass if this goes high enough, if this is selective enough. Even in the low pass mode it can act as a band pass filter if the Q, once again I go back to definition of Q, Q is 1 by 2 zeta. If Q is sufficiently high or zeta is sufficiently low.

So in a any of the 3 modes, the circuit can act as an band pass filter, provided, the resistance in the circuit is not excessively large. It is this resistance which gives the dumping affect and therefore, the resistance, if the resistance is not excessively large, it may act as a band pass filter. But it may be of interest to find out, in the high pass case, in the high pass mode, what is the frequency of maximum, what are the conditions for a peaking, what is the condition for

critical damping, that is, what is the condition for a transition between high pass and band pass, that is, peaking or no peaking, and how to find out the 3 db cut off?

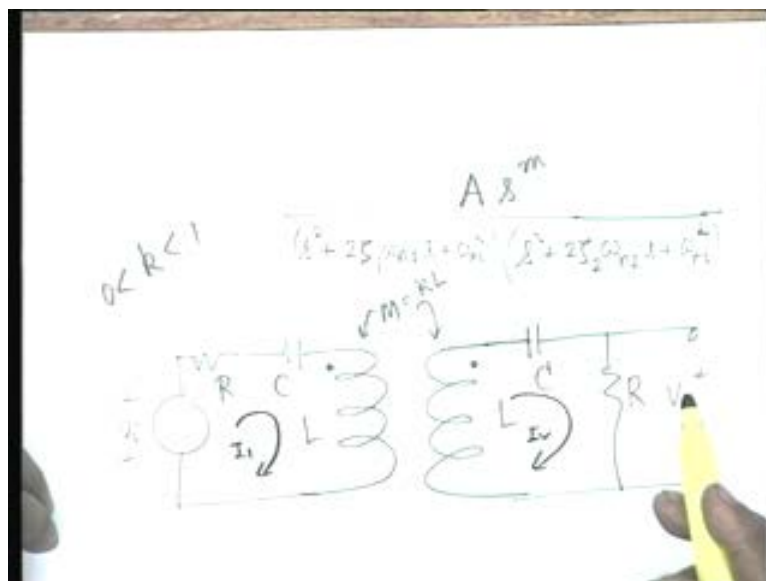
It is worth investigating. May I also mention that in the high pass mode, it is very easy to see that the transfer function would be of the form s^m divided by $s^2 + 2\zeta\omega_n s + \omega_n^2$, some constant K divided by $s^2 + 2\zeta\omega_n s + \omega_n^2$. Do you see this? That it would be s^m , because this is sL divided by $R + sL + 1/sC$ so it will be s^2 . In the low pass mode, the numerator was a constant, s^2 term was not there. In the band pass mode, the numerator has a power of s . So the single tuned circuit is indeed of versatile circuit, in which, you can get $K s^m$ divided by $s^2 + 2\zeta\omega_n s + \omega_n^2$.

(Refer Slide Time: 08:26)



A transfer function of this type, where m can either 0 1 or 2, agreed? m is equal to 0 is the low pass mode, m equal to 2 is the high pass mode and m equal to 1 is the band pass mode. It is worth investigating all these 3. Now we next consider the case of a pair, 2 pairs of complex conjugate roots, that is, we consider transfer functions in which there are 2 pairs of complex conjugate roots. For example, $\frac{2\zeta_1\omega_{n1}s + \omega_{n1}^2}{s^2 + 2\zeta_2\omega_{n2}s + \omega_{n2}^2}$. We consider 2 pairs of complex conjugate roots.

(Refer Slide Time: 9:15)



The numerator may be a constant multiplied by some power of s . Obviously, the power n , can it be negative? No, it cannot be negative. So it is either 0, 1, 2, 3 or it can also be 4, depending on how you take the output. Let us see and a practical circuit which gives this kind of a response is the, so called, double tuned circuit or a circuit which is magnetically coupled each other, the primary, as well as the secondary, both are tuned circuits. That is, 2 tuned circuits coupled to each other by means of their inductances by means of a mutual inductance.

We consider this circuit, this is my $V I$, Laplace transform voltage and then I have, let us say, an inductor and the resistor and this is my $V 0$. This is a typical mode in which it is operated, and to keep life simple, let us suppose that there are identical tuned circuits. That is, the resistance capacitance inductance are all identical and that there is a mutual inductance between them, let us say these are the dots. There is a mutual inductance which is equal to, as you know, square root of $L 1 L 2$, k times square root of $L 1 L 2$, so this would be k times L and k is the coefficient of coupling, k lies between 0 and 1.

We shall investigate this circuit, it is 1 example of producing, 1 example of a circuit which produces 2 pairs of complex conjugate poles and we shall see how this simple circuit can behave in a wide variety of manners. To be able to analyze this, let us suppose, let us

consider to loop currents I_1 and I_2 . You notice that m is the dots are in the favorable direction, but I have reverse the direction of current and therefore, the mutual inductance term shall come with the negative sign. Is this clear?

(Refer Slide Time: 12:44)

The image shows handwritten equations on a whiteboard. The first equation is $V_i = \overbrace{\left(R + \frac{1}{sC} + sL\right)}^Z I_1 - sM I_2$. The second equation is $0 = -sM I_1 + \underbrace{\left(R + sL + \frac{1}{sC}\right)}_Z I_2$. The third equation is $H(s) = \frac{V_o}{V_i} = \frac{I_2 R}{V_i}$. The fourth equation is $I_2 = \frac{sM}{Z} I_1$.

So my equation shall be, V_i would be equal to R plus 1 over sC plus sL multiplied by I_1 , this is the self impedance term. My primary question becomes R plus 1 over sC plus sL times I_1 and then the mutual inductance terms, as I said, it should come with a negative sign $sM I_2$ and for the second loop, that is, for the secondary, 0 would be equal to minus $sM I_1$ plus R plus sL plus 1 over sC times I_2 . These are the 2 loop equations and my transfer function H of s is equal to V_o by V_i V_o is $I_2 R$.

The current I_2 flowing through the resistance R , divided by V_i and from the second equation, from equation number 2, you can see, I would like to put some abbreviations let us call this Z and this one as Z . Then obviously, you would notice from the second equation that I_2 equal to sM by Z time I_1 . From the second equation I_2 equal to sM by Z times I_1 and if I substitute this in the expression for transfer function, I get H of s equals to sM by Z times I_1 times R $I_2 R$ divided by V_i which is $Z I_1$ minus sM times I_2 .

(Refer Slide Time: 14:23)

$$\begin{aligned}
 H(s) &= \frac{\frac{sM}{Z} \cdot R}{Z \cdot 1 - \frac{s^2 M^2}{Z} \cdot 1} \\
 &= \frac{sMR}{Z^2 - s^2 M^2} \\
 Z^2 &= \left(R + sL + \frac{1}{sC} \right)^2 = \frac{(s^2 LC + sCR + 1)^2}{s^2 C^2} \\
 &= \frac{L^2}{s^2} \left(s^2 + s \frac{R}{L} + \frac{1}{LC} \right)^2
 \end{aligned}$$

So it would be s squared M squared by Z times I 1 and you can now cancel I 1 and get this as s M R divided by Z squared multiplied by Z also, minus s squared M squared. Now let us simplify Z squared. Z squared is R plus s L plus 1 over s C, whole squared, which I can write as s squared C squared and then s squared L C plus s C R plus 1, whole squared, and this I can write as if I take L C out. Then I shall get L squared by s squared, is that okay? This s squared, I shall cancel it, if I take L C out it will come L squared C squared, C squared C squared cancel, multiplied by s squared plus s R by L plus 1 over L C whole squared.

Now I introduce the notations, let me write these equations again, H of s equal to s M R divided by Z squared minus s squared M squared and I have found out this Z squared as equal to L square by s squared, s squared plus. Now R by L, if you recall, I represent this as twice zeta omega n. In terms of zeta, the dumping in factor and omega n the natural, undamped natural frequency, so I get twice zeta omega n s and I also make 1 by L C as omega n squared. Therefore, I get omega n squared. Therefore, my H of s the transfer function becomes s M R divided by L squared.

Student: It would be a whole squared sir.

Sir: It would be a whole squared, thank you. L squared, the s square that occurs here, I can take it in the numerator as s cubed. Then, in the subtraction, minus s squared M squared, I

must have s to the fourth M squared. This s squared, I am taking, I am multiplying both numerator denominator by s squared. So this becomes s to the fourth M squared and here I shall have s squared plus twice zeta omega n s plus omega n squared. A bit of algebra, not too difficult though.

I can simplify this further by putting M equal to K L. Have I missed a term? No. what I will do is, if I write this again, H of s is equal to s cubed M R, I divided by L squared, both the numerator and denominator. Then in the denominator, this becomes s squared plus twice zeta omega n s plus omega n squared. Whole squared minus, I divided by L squared so I get the s to the fourth time, K squared

(Refer Slide Time: 18:45)

The image shows a whiteboard with handwritten mathematical equations. The top equation is:

$$H(s) = \frac{s^3 MR / L^2}{(s^2 + 2\zeta\omega_n s + \omega_n^2)^2 - s^2 k^2}$$

The middle equation shows the substitution of M = KL:

$$\frac{MR}{L^2} = \frac{KL^2}{L^2} = 2\zeta\omega_n K$$

The bottom equation shows the simplified transfer function:

$$H(s) = \frac{2\zeta\omega_n K s^3}{(s^2 + 2\zeta\omega_n s + \omega_n^2)^2 - (s^2 k)^2}$$

Simply because, M is equal to K times L, so what is M R divided by L squared? The M is K L and therefore, it is K R divided by L and what is R by L?

Student: (...)

Sir: Omega n and therefore, I get twice zeta omega n times K. This quantity becomes twice zeta omega n K therefore, my write each H of s is equal to twice zeta omega n K multiplied by s cubed divided by s squared plus twice zeta omega n s plus omega n squared whole squared minus s squared K, whole squared. No simplification, no approximation so far. It is

exact. Now the denominator, you can see, it is of the form $s^2 - B^2$ and therefore, you should be able to take, the factorize it to $A + B$ and $A - B$ and the result is the following.

(Refer Slide Time: 20:13)

$$H(s) = \frac{2\zeta\omega_n K s^3}{[s^2(1+K) + 2\zeta\omega_n s + \omega_n^2][s^2(1-K) + \omega_n^2]}$$

$$= \frac{2\zeta\omega_n K}{1-K^2} \cdot \frac{s^3}{\left[\left(s^2 + \frac{2\zeta\omega_n}{1-K} s + \frac{\omega_n^2}{1-K} \right) \left(s^2 + \frac{2\zeta\omega_n}{1-K} s + \frac{\omega_n^2}{1-K} \right) \right]}$$

I get H of S is equal to twice zeta omega n K s^3 divided by s^4 , if I add the 2 terms $A + B$, then I get $1 + K$ plus twice zeta omega n s plus omega n squared and the second factor would be same, except for a change of sign of K . That is, I shall get $s^2 - 1 + K$ plus twice zeta omega n s plus omega n squared. Let us do a little more simplification, let us take $1 + K$ out from here and $1 - K$ out from here. Then I get twice zeta omega n K divided by $1 - K^2$, multiplied by s^3 ; do you understand what I am doing? $1 + K$ I have taken out, and $1 - K$, I have taken out.

So the product is $1 - K^2$, that is what happens, that is what appears here and the factors that will be left would be $s^2 + 2\zeta\omega_n s + \omega_n^2$ divided by $1 + K s^2 + 2\zeta\omega_n s + \omega_n^2$. This is the 1 of these factors, multiplied by is

Students: (...)

$1 + K$, yes, I must not miss that, multiplied by $s^2 + 2\zeta\omega_n s + \omega_n^2$ divided by $1 - K s^2 + 2\zeta\omega_n s + \omega_n^2$ divided $1 - K$. Now we have the expression in the

form of a constant A. A constant A multiplied by s cubed therefore; there are 3 zeros at the origin. 3 zeros at the origin because of s cube. In the denominator, I have 2 quadratic factors and each quadratic factors shall have 2 roots and therefore, I have 4 poles and, as you know, if the poles are on the real axis, the case is of very little interest. The real axis poles do not offer selectivity.

If the case would be of interest and are our trouble of introducing 4 energy storage elements; 2 inductors, 2 capacitors. In fact, 5 energy, the mutual inductance also stores energy, 5 storage elements would be successful provided the poles are complex conjugate. This is a real polynomial, real polynomial, so poles shall be complex conjugate. Now assuming that the poles are complex conjugate, I can write my H of s as equal to A s cubed, divided by s minus s 1 multiplied by s minus s 1 star. s 1 and s 1 star are the roots of the first polynomial, then s minus s 2 s minus s 2 star.

(Refer Slide Time: 23:14)

$$H(s) = \frac{A s^3}{(s - s_1)(s - s_1^*)(s - s_2)(s - s_2^*)}$$

s_1, s_1^* are roots of

$$s^2 + \frac{2\zeta\omega_n}{1+k} s + \frac{\omega_n^2}{1+k} = 0$$

$$s_1, s_1^* = \frac{-\frac{2\zeta\omega_n}{1+k} \pm j \sqrt{\frac{\omega_n^2}{1+k} - \frac{4\zeta^2\omega_n^2}{(1+k)^2}}}{2}$$

$$= -\frac{\zeta\omega_n}{1+k} \pm j \omega_n \sqrt{\frac{1}{1+k} - \frac{\zeta^2}{(1+k)^2}}$$

To understand the approximations that we shall be using a little later, let us look at s 1 and s 1 star. Now s 1 and s 1 star are the roots of the equation s squared plus twice zeta omega n divide 1 plus k s plus omega n squared divided 1 plus k equal to 0, agreed. Therefore, s 1, s 1 star would be equal to minus 2 zeta omega n divided 1 plus k plus minus j, square root of s which term shall I have first?

Student: (...)

Sir: ω_n^2 divided by $1 + k$ minus, now 4 times this. $4\zeta^2 \omega_n^2$ square divided by $1 + k$ whole square, divided by 2, you must not forget that. Now this 2 can be cancelled out, 2 with this 2 this 4 and this 4. So I get, equal to minus $\zeta \omega_n$ divided by $1 + k$ plus minus j . I can take ω_n out, ω_n from ω_n^2 , I can take this out. Square root of 1 by $1 + k$ minus ζ^2 divided by $1 + k$ whole squared, agreed.

i have not made any approximation so far. I have only made simplification and I made the assumption that the poles are complex. If they are not complex, then you see ζ must be greater than the critical value. There shall be no peaking. Even if there is peaking, it shall be horrible, the q shall be large. q is 1 by 2ζ , I am sorry, q shall be small and therefore, the circuit shall offer no selectivity. It is only for the purpose of selectivity that we take the trouble of combining energy storage elements of different kinds. Similarly, if you, let me right this again s_1, s_1^* is equal to minus $\zeta \omega_n$ divided by $1 + k$ plus minus $j \omega_n$ square root of 1 by $1 + k$ minus ζ^2 divide by $1 + k$, whole squared.

(Refer Slide Time: 26:16)

$$s_{1,2} = -\frac{\zeta \omega_n}{1+k} \pm j \omega_n \sqrt{\frac{1}{1+k} - \frac{\zeta^2}{(1+k)^2}}$$

$$s_{1,2} = -\frac{\zeta \omega_n}{1-k} \pm j \omega_n \sqrt{\frac{1}{1-k} - \frac{\zeta^2}{(1-k)^2}}$$

$$\zeta^2 \ll 1, \quad k \ll 1$$

$$\frac{1}{\sqrt{1+k}} = (1+k)^{-1/2} \approx 1 - \frac{k}{2}$$

Similarly, I can write s_2, s_2^* , the only change will be k shall be replaced by minus k , that is all. So minus $\zeta \omega_n$ divided by $1 - k$ plus minus $j \omega_n$ square root of 1 by $1 - k$ minus ζ^2 divided by $1 - k$, whole squared. I now make some approximations. I now argue that a circuit of this complexity, where there are 2 coils which

are been brought close to each other, magnetically coupled, there are capacitances in both the circuits which shall be useful only if it is a high q situation. That is, if individually the tuned circuit, there are 2 tuned circuits identical tuned circuits, we assumed for simplicity, the tuned circuit themselves must be high q . That is, ζ must be small.

We do assume the ζ^2 is much less compared to unity. We assume the ζ^2 is much less compared to unity and we also assume that the couple that the coils are loosely coupled, that is, we also assume that k is much less than unity. They are loosely coupled. If that is the case, then I can ignore this k . Let me indicate the simplifications. I will ignore this k and this k in the denominator, I will take them out. I will ignore this term, as compared to 1 by 1 plus k . 1 by 1 plus k would be slightly less than unity but because of ζ^2 here, this should be very small compared to the first term, so I will ignore these 2 term also.

In addition, I will make the simplification that 1 by square root of 1 plus k is equal to 1 plus k to the power minus half, approximately equal to 1 minus k by 2. Mind you, the simplifications are based on the assumption of high q circuits, that is, low resistance. Capital R must be low, number 1. Number 2 is that the coils are not too critically coupled, too tightly coupled, they are loosely coupled. If that is so, then you notice that $s^2 + 1$ star becomes a very simple expression.

(Refer Slide Time: 29:07)

The image shows a whiteboard with handwritten mathematical equations. The equations are:

$$s_1, s_1^* = -\zeta\omega_n \pm j\omega_n \left(1 - \frac{k}{2}\right)$$

$$s_2, s_2^* = -\zeta\omega_n \pm j\omega_n \left(1 + \frac{k}{2}\right)$$

$$s_1 = -\zeta\omega_n + j\omega_n \left(1 - \frac{k}{2}\right)$$

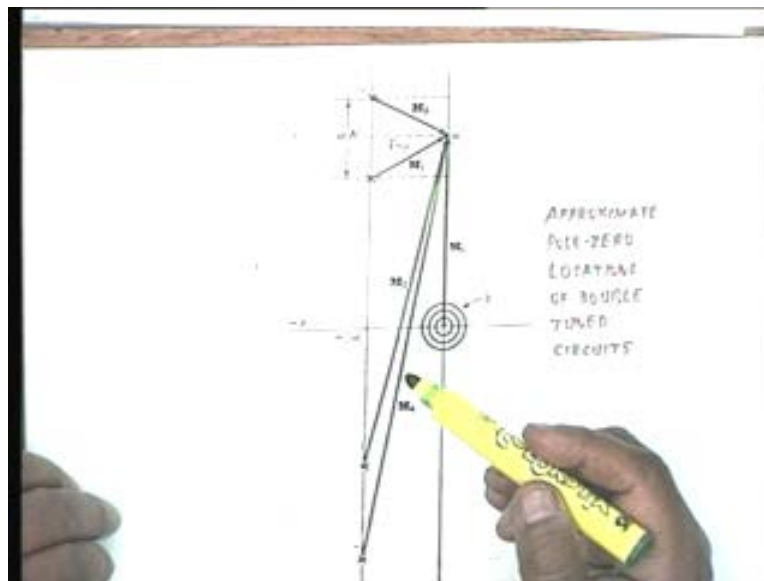
$$s_2 = -\zeta\omega_n + j\omega_n \left(1 + \frac{k}{2}\right)$$

$$s_2 - s_1 = j\omega_n k$$

It simply becomes minus zeta omega n plus minus j omega n 1 minus k by 2 and similarly, s 2, s 2 star that becomes minus zeta omega n plus minus j omega n, the only thing that happens is k changes its sign, so 1 plus k by 2 and the situation with this approximation, the location of the roots, you see that they occur, there is quad of roots, quad, 4 of them. And the real parts of all the roots are the same. That is, all the poles lie on a laying parallel to the j omega axis. The real part is same minus zeta omega n in the left half plane.

However, on this line, which is parallel to the j omega axis, there are 2 in the upper quadrant and 2 in the lower quadrant. 2 in the second quadrant and 2 in the third quadrant and the 2 in the second quadrant are symmetrical, with respect to omega n. Is not that right? One is j omega n 1 minus k by 2, the other is j omega n 1 plus k by 2 and the other 1 is negative, that is, in the third quadrant. So the situation is like this.

(Refer Slide Time: 30:48)



Let me explain, these are the 3 zeros at the origin and we indicate this by 3 concentric circles, 3 zeros. There are 4 poles with this approximation. This is 1, this 1, this 1, this is 1. 2 in the third quadrant, 2 in the second quadrant and 2 in the third quadrant, the real part is minus zeta omega n, for all of them. So this root is minus zeta omega n plus j omega n 1 minus k by 2. So this is omega n, you see, at a distance of zeta omega n. So it is half way below omega n, so this distance is omega n k. The distance between 2 poles, is that clear?

No, shall we go back to this? You see, what I am saying is, s_1 , s_1 is, let us say, $\omega_n - \zeta \omega_n + j \omega_n \sqrt{1 - \zeta^2}$ and s_2 is $\omega_n - \zeta \omega_n + j \omega_n \sqrt{1 - \zeta^2}$ by 2. So what is $s_1 - s_2$? It is simply, or $s_2 - s_1$, it is simply $j \omega_n \sqrt{1 - \zeta^2}$. So the difference between the 2 poles, it is simply, the distance is $\omega_n \sqrt{1 - \zeta^2}$ and it is on the $j \omega_n$ axis so $j \omega_n \sqrt{1 - \zeta^2}$. This distance is $\omega_n \sqrt{1 - \zeta^2}$ and you see that the simplification that it achieves is that, the 2 poles in the upper quadrant becomes symmetrical about the natural, undamped natural frequency of the tuned circuits about ω_n .

Similarly, these 2 in the lower quadrant, I am going to make further simplifications, so it is essential that you understand this simplification and the only thing that we assumed is that ζ is much less than 1 and k is much less than 1, actually ζ^2 . You see ζ , it is very easy to satisfy ζ^2 much less than 1 because if ζ is, let us say, point 5, ζ^2 is point 0.25. If ζ is point 1 and ζ^2 is point 0.1, 1 hundred. So under these conditions the poles and zeros of the double tuned circuits assume this particular form.

Now here we are going to make further simplifications, further simplifications like this. Suppose, I want to find out the magnitude responds at this frequency ω_n . Then as you know, my transfer function is $A s^3$ divided by $(s - s_1)(s - s_1^*)(s - s_2)(s - s_2^*)$.

(Refer Slide Time: 33:55)

$$H(s) = \frac{A s^3}{(s - \lambda_1)(s - \lambda_1^*)(s - \lambda_2)(s - \lambda_2^*)}$$

$$H(j\omega) = \frac{A (j\omega)^3}{(j\omega - \lambda_1)(j\omega - \lambda_1^*)(j\omega - \lambda_2)(j\omega - \lambda_2^*)}$$

$$|H(j\omega_n)| = \frac{A \omega_n^3}{\omega_n \omega_n \omega_n \omega_n}$$

Therefore, H of $j\omega$ would be equal to $A j\omega$ cubed divided by $j\omega$ minus s_1 , $j\omega$ minus s_1^* , $j\omega$ minus s_2 , multiplied by $j\omega$ minus s_2^* and as far as magnitude is concerned, all we do is, we draw a vector from the point $j\omega$ to the zeros and to the poles. We draw the vectors. These are the vectors, from the poles, from the pole to the point at which you want find out the magnitude response. From these 2 poles to the point $M_1 M_2 M_3 M_4$, these are the magnitudes of these vectors and M_0 is the vector drawn from the origin to from this zeros to ω_n and therefore, my H of $j\omega_n$, under this condition, shall be $A M_0^3$ cubed, 3 zeros divided by $M_1 M_2 M_3 M_4$, the magnitude. I can also find the phase from this diagram. Let us consternate on the magnitude at the moment, any question?

Student: Sir, what is this M_0 cubed?

Sir: M_0 cubed comes from $j\omega$ cubed M_0 is this vector, from the 0 to ω_n , it is the 0 vector and since it occur 3 times, we are taking m_0 cubed. Now I want you look at this diagram very carefully. You see, when ω , when our frequency, this is the situation to ω_n , now when our frequencies around ω_n , what does it mean, a narrow band centered around the undamped natural frequency? Then M_0 , even if ω shifts a little on the upper side or lower side, M_0 can be approximated as ω_n , agreed. When ω , let me write this again, I will come back to this diagram again and again.

(Refer Slide Time: 36:58)

$$\begin{aligned} \omega &\approx \omega_n \\ |H(j\omega)| &= \frac{A M_0^3}{M_1 M_2 M_3 M_4} \\ M_0 &\approx \omega_n \\ M_1, M_2 &\approx 2\omega_n \\ |H(j\omega)| &= \frac{A \omega_n^3}{M_1 M_2 4\omega_n^2} \\ &= \frac{A \omega_n / 4}{M_1 M_2} \end{aligned}$$

So H of $j\omega$, ω around ω_n , this is what, this is how we indicate it. ω around ω_n , H of $j\omega$ which is equal to $A M_0$ cubed M_0 is now a function of ω , divided by $M_1 M_2 M_3 M_4$. When ω is around ω_n , M_0 can be approximated by ω_n , agreed? When the frequencies close to ω_n . When the frequencies around this point, it can be approximated by ω_n . Not only that, if ζ is low, this figure is an exaggerated figure. ζ is low, the lines $j\omega$ axis and this line are very close to each other. If they are close to each other, naturally, M_3 and M_4 should also be close to each other.

These poles, you see, k is a small quantity, so loose coupling and therefore, this poles are also very close to each other, which means, that when frequencies around ω_n , both M_3 and M_4 can be taken to be approximately equal and approximately equal to twice M_0 , is that clear? So I do this approximation M_3 and M_4 both are approximately equal to twice ω_n . If I do that then H of $j\omega$ becomes equal to magnitude becomes equal to $A \omega_n$ cubed divided by $M_1, M_2, 4 \omega_n$ squared. That is equal to $A \omega_n$ divided by 4 divided by $M_1 M_2$. I hope you realize what I have done.

Student: Excuse me sir, in that you are assuming ω_n also to be small?

Sir: No.

Student: Sir otherwise what would happen M_3 and M_4 will not be equal to $2\omega_n$, $2\omega_n$?

Sir: Why not? You see, this line, this line is very close. Just a minute, I understand your confusion. You see, first we say because k is small ω is not necessarily small because k is small, these 2 poles are very close to each other, so M_3 and M_4 are approximately equal. Then I argue that these 2 lines are very close to each other.

Student: Sir, but this is ζ times ω_n .

Sir: Which is ζ times? So correct, ζ is a small quantity and therefore, they are very close to each other, therefore, M_3 and M_4 are approximately twice M_0 . Is not it approximation and a fairly good approximation, in practice, as far as design of this circuit is the concerned.

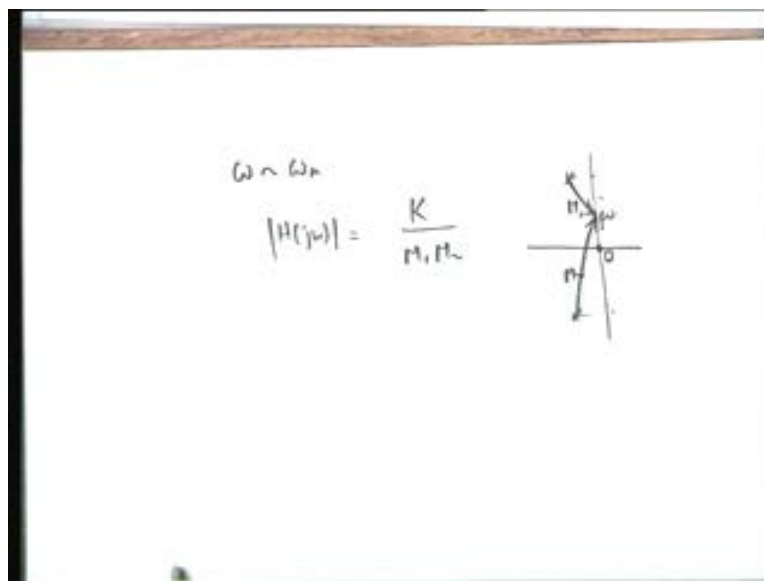
Student: Sir that means, we have approximated that s_2 and s_1 equal to $j\omega$

Sir: Quite so, quite so.

Student: Both the poles are at $j\omega$.

Sir: Quite so, that is correct and therefore, if you notice carefully what we have done is, this is constant, let us call it a K , and therefore, for the double tuned circuit, with ω approximately equal to ω_n around ω_n , the magnitude function has been approximated as K by M_1 and M_2 , which physically means, I go back to this diagram, which physically means that it is these 2 poles which are close to ω , approximately equal to ω_n , affect the magnitude response much more than the zeros and the other 2 poles.

(Refer Slide Time: 40:45)



This is all that I have expressed in terms of this mathematical approximation, which is obviously true. If ω shifts a little bit, M_1 and M_2 , the relative change in M_1 and M_2 are much greater than the relative change in M_3 , M_4 or M_0 . This is what we have done. No, I have done something else. If you recall this single tuned circuit, in the single tuned circuit, we had 2 poles like this and we had expressed the magnitude response as $K \sin \psi$ over $M_1 M_2$.

Student: (...)

Sir: Oh, it was simply $K \sin \psi$ because M_1 and M_2 are included but there also the magnitude response if this is ω . If this is $j\omega$ and this is M_1 and this is M_2 , magnitude responses is simple a constant divide by $M_1 M_2$ and therefore, virtually, the double tuned circuits has been converted to a single tuned circuits, with 1 difference. In a single tuned circuit, this was 0 and therefore what we have done is we have raised the origin or the real axis.

We have raised it by an amount ω_n and therefore, the totals, the response of the double tuned circuits at frequencies around ω_n , which are all concerned, will be basically determined, exactly in the same manner as that of a single tuned circuit. In other words, I can now bring in the concepts of the peaking circle. The peaking circle, for example, would be a circle, what would be the center?

Student: Sir, ω_n .

This should be the center and this as the radius, you draw a circle. If it cuts the $j\omega$ axis, then there shall be peaking. If it does not cut the $j\omega$ axis, then there shall be no peaking. That is not correct. Unfortunately, there shall be peaking. You see, this is not 0, this is not 0, this is ω_n .

Student: Sir, what did you say about raising the?

Sir: Raising, you see, in this single tuned circuit, this center was 0, was the origin of the complex plain. Whereas the center here is ω_n . So it is as if the real axis has been raised by a level, has been transformed to ω_n . Once that is established, this behaves exactly like single tuned circuit. However, whether the peaking circuit cuts the $j\omega$ axis or does not cut the $j\omega$ axis, it does not matter. There shall still be a peaking and this can be understood with reference to another diagram, which we shall project next time. It is 2 0 1, so we continue on Thursday.