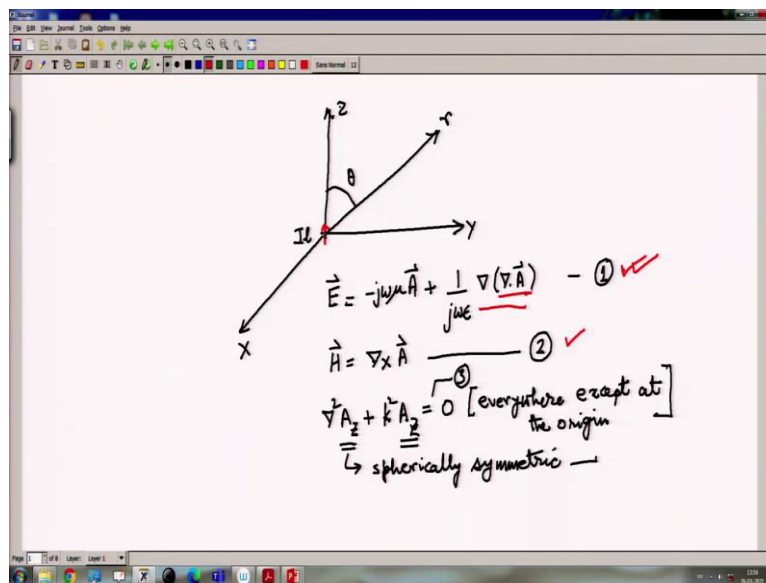


Advanced Microwave Guided-Structures and Analysis
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Lecture No. 23
Radiation from an Electric Current Source (contd.)

Welcome to this session, we will be continuing with the radiation from an electric current source, the spherical wave components for the radiated fields. Last time we investigated the magnetic field components and we saw that both the r , the radial component and the elevation component which is the θ component of the magnetic field H_θ , the magnetic field has only an azimuthal component or ϕ file.

Next, we will investigate the electric field components radiated by the current source. Let us go to the lecture. Now, in order to find the electric field components, we need to look again at question number 1.

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And for calculating the radiated electric field, we need to calculate divergence of \vec{A} and gradient of divergence of \vec{A} for the spherical wave. So, let us do that.

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$$\begin{aligned}
 \nabla \cdot \vec{A} &= \frac{1}{r^2 \sin \theta} \left[\frac{\partial}{\partial r} (r^2 \sin \theta A_r) + \frac{\partial}{\partial \theta} (r \sin \theta A_\theta) + \frac{\partial}{\partial \phi} (r A_\phi) \right] \\
 &= \frac{1}{r^2 \sin \theta} \left[\frac{\partial}{\partial r} \left(r^2 \sin \theta \frac{I_0}{4\pi r} e^{-jkr} \cos \theta \right) \right. \\
 &\quad \left. + \frac{\partial}{\partial \theta} \left(r \sin \theta (-1) \frac{I_0}{4\pi r} e^{-jkr} \sin \theta \right) \right] \\
 &= \frac{1}{r^2 \sin \theta} \left[\frac{I_0}{4\pi} \sin \theta \cos \theta \frac{\partial}{\partial r} (r e^{-jkr}) - \frac{I_0}{4\pi} e^{-jkr} \frac{\partial}{\partial \theta} (\sin^2 \theta) \right] \\
 &= \frac{1}{r^2 \sin \theta} \left[\frac{I_0}{4\pi} \sin \theta \cos \theta \left\{ r(-jk)e^{-jkr} + e^{-jkr} \right\} - \frac{I_0}{4\pi} e^{-jkr} 2 \sin \theta \cos \theta \right] \\
 &= \frac{I_0}{4\pi} e^{-jkr} \left[-\frac{jk}{r} + \frac{1}{r^2} - \frac{2}{r^2} \right] \cos \theta
 \end{aligned}$$

So, divergence of A is $\frac{1}{r^2 \sin \theta} \left[\frac{\partial}{\partial r} (r^2 \sin \theta A_r) + \frac{\partial}{\partial \theta} (r \sin \theta A_\theta) + \frac{\partial}{\partial \phi} (r A_\phi) \right]$ and that is equal to $\frac{1}{r^2 \sin \theta} \left[\frac{\partial}{\partial r} (r^2 \sin \theta \frac{I_0}{4\pi r} e^{-jkr} \cos \theta) + \frac{\partial}{\partial \theta} (r \sin \theta (-1) \frac{I_0}{4\pi r} e^{-jkr} \sin \theta) \right]$. So, we substitute the value of A_r and A_θ in A_ϕ . So, $\frac{\partial}{\partial r} (r^2 \sin \theta \frac{I_0}{4\pi r} e^{-jkr} \cos \theta) + \frac{\partial}{\partial \theta} (r \sin \theta (-1) \frac{I_0}{4\pi r} e^{-jkr} \sin \theta)$ plus 0 because A_ϕ is 0.

And that is equal to $\frac{1}{r^2 \sin \theta} \left[\frac{I_0}{4\pi} \sin \theta \cos \theta \frac{\partial}{\partial r} (r e^{-jkr}) - \frac{I_0}{4\pi} e^{-jkr} \frac{\partial}{\partial \theta} (\sin^2 \theta) \right]$ and that becomes equal to $\frac{1}{r^2 \sin \theta} \left[\frac{I_0}{4\pi} \sin \theta \cos \theta \left\{ r(-jk)e^{-jkr} + e^{-jkr} \right\} - \frac{I_0}{4\pi} e^{-jkr} 2 \sin \theta \cos \theta \right]$.

And that is equal to $\frac{I_0}{4\pi} e^{-jkr} \left[-\frac{jk}{r} + \frac{1}{r^2} - \frac{2}{r^2} \right] \cos \theta$.

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$$\vec{\nabla} \cdot \vec{A} = \frac{I_0}{4\pi} e^{-jkr} \left[-\frac{jk}{r} - \frac{1}{r^2} \right] \cos\theta \quad \text{--- (9)}$$

$$\nabla(\vec{\nabla} \cdot \vec{A}) = \hat{a}_r \frac{\partial}{\partial r} \left[\frac{I_0}{4\pi} e^{-jkr} \left\{ -\frac{jk}{r} - \frac{1}{r^2} \right\} \cos\theta \right] +$$

$$\frac{1}{r} \hat{a}_\theta \frac{\partial}{\partial \theta} \left[\frac{I_0}{4\pi} e^{-jkr} \left\{ -\frac{jk}{r} - \frac{1}{r^2} \right\} \cos\theta \right]$$

$$+ \frac{1}{r \sin\theta} \hat{a}_\phi \frac{\partial}{\partial \phi} \left[\frac{I_0}{4\pi} e^{-jkr} \left\{ -\frac{jk}{r} - \frac{1}{r^2} \right\} \cos\theta \right] \quad \text{--- (10)}$$

Which can be further simplified as, divergence of A is I_0 by 4π e^{-jkr} minus jk by r minus 1 by r square, minus 1 by r square times $\cos\theta$. We call this equation 9 and then we proceed to calculate grad of divergence A and then we proceed to calculate grad of divergence A and that is equal to $\hat{a}_r \frac{\partial}{\partial r} r \frac{I_0}{4\pi} e^{-jkr} \left\{ -\frac{jk}{r} - \frac{1}{r^2} \right\} \cos\theta$ minus jk by r minus 1 by r square $\cos\theta$.

Plus $\frac{1}{r} \hat{a}_\theta \frac{\partial}{\partial \theta} \left[\frac{I_0}{4\pi} e^{-jkr} \left\{ -\frac{jk}{r} - \frac{1}{r^2} \right\} \cos\theta \right]$ plus $\frac{1}{r \sin\theta} \hat{a}_\phi \frac{\partial}{\partial \phi} \left[\frac{I_0}{4\pi} e^{-jkr} \left\{ -\frac{jk}{r} - \frac{1}{r^2} \right\} \cos\theta \right]$. So, let us call this step this is by the definition of the gradient operation in the spherical coordinate system.

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from (1):

$$\vec{E} = -j\omega\mu \vec{A} + \frac{1}{j\omega\epsilon} \nabla(\vec{\nabla} \cdot \vec{A})$$

$$E_r = -j\omega\mu \frac{I_0}{4\pi r} \cos\theta + \frac{1}{j\omega\epsilon} \frac{\partial}{\partial r} \left[\frac{I_0}{4\pi} e^{-jkr} \left\{ -\frac{jk}{r} - \frac{1}{r^2} \right\} \cos\theta \right]$$

$$= -j\omega\mu \frac{I_0}{4\pi r} e^{-jkr} \cos\theta + \frac{I_0}{4\pi} \frac{1}{j\omega\epsilon} (-jk) \cos\theta \frac{\partial}{\partial r} \left(\frac{e^{-jkr}}{r} \right)$$

$$- \frac{I_0}{4\pi} \frac{1}{j\omega\epsilon} \cos\theta \frac{\partial}{\partial r} \left(\frac{e^{-jkr}}{r^2} \right)$$

$$= -j\omega\mu \frac{I_0}{4\pi r} e^{-jkr} \cos\theta + \frac{I_0}{4\pi} \frac{1}{j\omega\epsilon} (-jk) \cos\theta \left(-\frac{jk}{r} e^{-jkr} - \frac{1}{r^2} e^{-jkr} \right)$$

$$- \frac{I_0}{4\pi} \left(\frac{1}{j\omega\epsilon} \right) \cos\theta \left(-\frac{jk}{r^2} e^{-jkr} - \frac{2}{r^3} e^{-jkr} \right)$$

So, now we know that from equation 1, E is equal to minus j omega mu A plus 1 by j omega epsilon grad of divergence A, grad of divergence grad of divergence. So, Er therefore, the r component of the electric field will be equal to minus j omega mu Az cos theta plus 1 by j omega mu, the r component of grad of divergence A which we derived before which is del r II by 4 pi e to the power minus jkr minus jk by r minus 1 by r square times cos theta.

And that is equal to minus j omega mu substituting for z, II by 4 pi r e to the power minus j kr cos theta plus II by 4 pi 1 by j omega epsilon minus jk cos theta del del r will work on e to the power minus j kr by r minus II by 4 pi 1 by j omega epsilon cos theta del del r is going to work on e to the power minus j kr by r square. And that will give rise to minus j omega mu II by 4 pi r e to the power minus j kr cos theta plus II by 4 pi 1 by j omega epsilon minus jk cos theta times minus jk by r e to the power minus j kr minus 1 by r square e to the power minus j kr, for the first term.

For the second term and then for the third term minus II by 4 pi 1 by j omega epsilon, 1 by j omega epsilon cos theta minus jk by r square e to the power minus j kr minus 2 by r cube to the power minus j kr.

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$$E_r = \frac{I_0}{4\pi} e^{-jkr} \cos\theta \left[-\frac{j\omega\mu}{r} + \frac{1}{j\omega\epsilon} (-jk) \left(-\frac{jk}{r^2} - \frac{1}{r^3} \right) - \frac{1}{j\omega\epsilon} \left(\frac{-jk}{r^2} - \frac{2}{r^3} \right) \right]$$

$$= \frac{I_0}{4\pi} e^{-jkr} \cos\theta \left[-\frac{j\omega\mu}{r} - \frac{k^2}{r} \frac{1}{j\omega\epsilon} + \frac{jk}{j\omega\epsilon} \cdot \frac{1}{r^2} + \frac{1}{j\omega\epsilon} \frac{jk}{r^2} + \frac{2}{j\omega\epsilon} \cdot \frac{1}{r^3} \right]$$

$$\therefore \frac{k^2}{j\omega\epsilon} = \frac{\omega^2\mu\epsilon}{j\omega\epsilon} = -j\omega\mu$$

$$E_r = \frac{I_0}{2\pi} e^{-jkr} \cos\theta \left(\frac{k}{\omega\epsilon} \frac{1}{r^2} + \frac{1}{j\omega\epsilon r^3} \right)$$

$$\therefore \frac{k}{\omega\epsilon} = \frac{\omega\sqrt{\mu\epsilon}}{\omega\epsilon} = \eta$$

This can be further simplified as Er equal to II by 4 pi e to the power minus j kr cos theta times minus j omega mu by r plus 1 by j omega epsilon times minus jk times minus jk by r square minus 1 by r square minus 1 by j omega epsilon minus jk by r square minus 2 by r cube.

And this can be further written as $\frac{I l}{4 \pi \epsilon_0 r^2} e^{-j k r} \cos \theta$ minus $j k r$ times minus $j \omega \mu$ by r minus k^2 by r plus $j \omega \epsilon$ plus $j k$ by $j \omega \epsilon$ plus 1 by r^2 plus 1 by $j \omega \epsilon$ plus $j k$ by r^2 plus 2 by $j \omega \epsilon$ times 1 by r^3 . Since, now k^2 by $j \omega \epsilon$ is $\omega^2 \mu \epsilon$ by $j \omega \epsilon$ that is equal to minus $j \omega \mu$. E_r can be written as $\frac{I l}{2 \pi \epsilon_0 r^2} e^{-j k r} \cos \theta$ times k by $\omega \epsilon$ times 1 by r^2 plus 1 by $j \omega \epsilon$ times r^3 .

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$$E_r = \frac{I l}{2 \pi} e^{-j k r} \cos \theta \left(\frac{1}{r^2} + \frac{1}{j \omega \epsilon r^3} \right)$$

$$E_\theta = j \omega \mu A_z \sin \theta + \frac{1}{j \omega \epsilon} \frac{1}{r} \frac{\partial}{\partial \theta} \left[\frac{I l}{4 \pi} e^{-j k r} \left\{ -\frac{j k}{r} - \frac{1}{r^2} \right\} \cos \theta \right]$$

$$= j \omega \mu \frac{I l}{4 \pi r} e^{-j k r} \sin \theta + \frac{1}{j \omega \epsilon} \frac{1}{r} \frac{I l}{4 \pi} e^{-j k r} \left(-\frac{j k}{r} - \frac{1}{r^2} \right) (-\sin \theta)$$

$$E_\theta = \frac{I l}{4 \pi} e^{-j k r} \left(\frac{j \omega \mu}{r} + \frac{1}{r^2} + \frac{1}{j \omega \epsilon r^3} \right) \sin \theta$$

$$E_\phi = 0$$

Now, since k by $\omega \epsilon$ equal to $\omega \sqrt{\mu \epsilon}$ by $\omega \epsilon$ equal to η is root of μ by ϵ , the intrinsic impedance of free space we can write the value of E_r as $\frac{I l}{2 \pi \epsilon_0 r^2} e^{-j k r} \cos \theta$ times η by r^2 plus 1 by $j \omega \epsilon$ times r^3 . So, this completes the derivation of the radial component of the electric field radiated by the electric current point source.

Now, for the theta directed electric field we have that equal to $j \omega \mu A_z \sin \theta$ plus 1 by $j \omega \epsilon$ times 1 by r times $\frac{\partial}{\partial \theta}$ of $\frac{I l}{4 \pi} e^{-j k r}$ times $\cos \theta$ minus $j k$ by r minus 1 by r^2 times $\cos \theta$. So, this is the theta component of grad of divergence of A times 1 by $\epsilon_0 \omega \mu$. So, this is the theta component of 1 by $\epsilon_0 \omega \mu$ times grad of divergence of A .

And that now on substitution of A_z becomes equal to $j \omega \mu \frac{I l}{4 \pi r} e^{-j k r} \sin \theta$ and we need to do just this theta differentiation plus 1 by $j \omega \epsilon$ times 1 by r times $\frac{\partial}{\partial \theta}$ of $\frac{I l}{4 \pi} e^{-j k r}$ times $\cos \theta$ minus $j k$ by r minus 1 by r^2 times the differentiation with respect to θ which is minus $\sin \theta$ times the differentiation of $\cos \theta$ with respect to θ , which is minus $\sin \theta$.

And then combining the terms together we have and then combining the terms together we can finally write E_θ as $\frac{1}{4\pi\epsilon_0} \frac{j\omega\mu_0 I_0 a^2 \sin\theta}{r^3} \left(\frac{1}{r} + j\omega\epsilon_0 r \right) \sin\theta$. So, this is the expression for E_θ .

Finally, E_ϕ is 0 because E does not have any component along the ϕ direction. So, the first term of the electric field which is $-\frac{j\omega\mu_0 I_0 a^2 \sin\theta}{4\pi\epsilon_0 r^3}$ that is 0 and the second term if you look at the expression for grad of divergence of A , if you look at the second term, at the third term, if you look at the third term of the grad of divergence of A , that is $\frac{\partial}{\partial \phi} \frac{\partial \phi}{\partial \phi}$ of this term. So, this term has no ϕ component. So, therefore, the ϕ component of grad of divergence of A is 0.

And therefore, E_ϕ becomes equal to 0. So, this completes the derivation of the three electric field components radiated by the point current source. We note that E_ϕ is 0 and the electric field is along the radial direction and the θ direction. So, we conclude this lecture. Next we are going to go to the magnetic current source. Thank you.