

**Advanced Microwave Guided-Structures and Analysis**  
**Professor Bratin Ghosh**  
**Department of Electronics and Electrical Communication**  
**Indian Institute of Technology, Kharagpur**  
**Lecture 55**  
**Analysis of Guided Structures (cont.)**

So, welcome to the remaining part of the lecture on the non-radiating dielectric guide. So, today we are going to finish the explicit field expressions for the remaining modes of the NRD, as well as the discuss the characteristic equation of the NRD, from which the cut off frequencies of all the modes can be obtained.

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For the LSM odd mode:

In the dielectric region ( $|x| \leq \frac{b}{2}$ ):

$$E_x = \frac{A_0 (k_x^2 \epsilon_r - \gamma_{m0}^2)}{j\omega \epsilon_r} \sin(\gamma_{m0} x) \sin\left(\frac{n\pi y}{a}\right) e^{-jk_0 z}$$

$$E_y = \frac{A_0 \gamma_{m0} n\pi}{j\omega a \epsilon_r} \cos(\gamma_{m0} x) \cos\left(\frac{n\pi y}{a}\right) e^{-jk_0 z}$$

So, the last time we had discuss the explicit field expressions for the LSM even mode. We are going to write down the explicit field expressions for the LSM odd mode, and thereafter the LSE even mode and the LSE odd mode, not a very difficult exercise, but just for reference sake. So, you should verify yourself as you see these expressions.

So, for the LSM odd mode, we write in the dielectric region,

For the LSM odd mode:

In the dielectric region ( $|x| \leq \frac{b}{2}$ ):

$$E_x = \frac{A_{d0} (k_0^2 \epsilon_r - \beta_{m0}^2)}{j\omega \epsilon_0 \epsilon_r} \sin(\beta_{m0} x) \sin\left(\frac{n\pi y}{a}\right) e^{-jk_{z0} z}$$

$$E_y = \frac{A_{d0} \beta_{m0} n\pi}{j\omega \epsilon_0 \epsilon_r} \cos(\beta_{m0} x) \cos\left(\frac{n\pi y}{a}\right) e^{-jk_{z0} z}$$

(Refer Slide Time: 4:15)

The screenshot shows a presentation slide with the following content:

$$H_z = -\frac{A_{d0} n\pi}{a} \sin(\beta_{m0} x) \cos\left(\frac{n\pi y}{a}\right) e^{-jk_{z0} z}$$

In the air region ( $|x| \geq \frac{b}{2}$ ):

$$E_x = \frac{A_{a0} (\beta_{m0}^2 + k_0^2)}{j\omega \epsilon_0} \exp\left[-\beta_{m0} \left(|x| - \frac{b}{2}\right)\right] \sin\left(\frac{n\pi y}{a}\right) e^{-jk_{z0} z} \quad (\text{for } x > \frac{b}{2})$$

$$= -\frac{A_{a0} (\beta_{m0}^2 + k_0^2)}{j\omega \epsilon_0} \exp\left[-\beta_{m0} \left(|x| - \frac{b}{2}\right)\right] \sin\left(\frac{n\pi y}{a}\right) e^{jk_{z0} z} \quad (\text{for } x < -\frac{b}{2})$$

The handwritten derivation on the slide is as follows:

$$H_z = -\frac{A_{d0} n\pi}{a} \sin(\beta_{m0} x) \cos\left(\frac{n\pi y}{a}\right) e^{-jk_{z0} z}$$

In the air region ( $|x| \geq \frac{b}{2}$ ):

$$E_x = \frac{A_{a0} (\beta_{m0}^2 + k_0^2)}{j\omega \epsilon_0} \exp\left[-\beta_{m0} \left(|x| - \frac{b}{2}\right)\right] \sin\left(\frac{n\pi y}{a}\right) e^{-jk_{z0} z} \quad (\text{for } x > \frac{b}{2})$$

$$= -\frac{A_{a0} (\beta_{m0}^2 + k_0^2)}{j\omega \epsilon_0} \exp\left[-\beta_{m0} \left(|x| - \frac{b}{2}\right)\right] \sin\left(\frac{n\pi y}{a}\right) e^{jk_{z0} z} \quad (\text{for } x < -\frac{b}{2})$$

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The screenshot shows a digital whiteboard with the following handwritten equations:

$$E_y = \frac{-A_{a0} n \pi p_{m0}}{j \omega a \epsilon_0} \exp\left[-p_{m0} \left(|x| - \frac{b}{2}\right)\right] \cos\left(\frac{n \pi y}{a}\right) e^{-j k_{z0} z}$$

$$E_z = \frac{A_{a0} k_{z0} p_{m0}}{\omega \epsilon_0} \exp\left[-p_{m0} \left(|x| - \frac{b}{2}\right)\right] \sin\left(\frac{n \pi y}{a}\right) e^{-j k_{z0} z}$$

$$H_x = 0$$

$$H_y = -A_{a0} j k_{z0} \exp\left[-p_{m0} \left(|x| - \frac{b}{2}\right)\right] \sin\left(\frac{n \pi y}{a}\right) e^{-j k_{z0} z}$$

(for  $x \geq \frac{b}{2}$ )

$$= A_{a0} j k_{z0} \exp\left[-p_{m0} \left(|x| - \frac{b}{2}\right)\right] \sin\left(\frac{n \pi y}{a}\right) e^{-j k_{z0} z}$$

(for  $x \leq -\frac{b}{2}$ ).

Similarly,

The handwritten equations are identical to those in the screenshot:

$$E_y = \frac{-A_{a0} n \pi p_{m0}}{j \omega a \epsilon_0} \exp\left[-p_{m0} \left(|x| - \frac{b}{2}\right)\right] \cos\left(\frac{n \pi y}{a}\right) e^{-j k_{z0} z}$$

$$E_z = \frac{A_{a0} k_{z0} p_{m0}}{\omega \epsilon_0} \exp\left[-p_{m0} \left(|x| - \frac{b}{2}\right)\right] \sin\left(\frac{n \pi y}{a}\right) e^{-j k_{z0} z}$$

$$H_x = 0$$

$$H_y = -A_{a0} j k_{z0} \exp\left[-p_{m0} \left(|x| - \frac{b}{2}\right)\right] \sin\left(\frac{n \pi y}{a}\right) e^{-j k_{z0} z}$$

(for  $x \geq \frac{b}{2}$ )

$$= A_{a0} j k_{z0} \exp\left[-p_{m0} \left(|x| - \frac{b}{2}\right)\right] \sin\left(\frac{n \pi y}{a}\right) e^{-j k_{z0} z}$$

(for  $x \leq -\frac{b}{2}$ ).

(Refer Slide Time: 10:09)

$$H_z = -\frac{A_{10} \pi \Gamma}{a} \exp\left[-\rho_{m0} \left(|x| - \frac{b}{2}\right)\right] \cos\left(\frac{n\pi y}{a}\right) e^{-jk_0 z} \quad (\text{for } x \geq \frac{b}{2})$$

$$= \frac{A_{10} \pi \Gamma}{a} \exp\left[-\rho_{m0} \left(|x| - \frac{b}{2}\right)\right] \cos\left(\frac{n\pi y}{a}\right) e^{-jk_0 z} \quad (\text{for } x \leq -\frac{b}{2})$$

Equating the tangential magnetic fields at the air-dielectric interface, we obtain:

$$A_{10} \sin\left(\frac{q_{m0} b}{2}\right) = A_1$$

$$H_z = -\frac{A_{10} \pi \Gamma}{a} \exp\left[-\rho_{m0} \left(|x| - \frac{b}{2}\right)\right] \cos\left(\frac{n\pi y}{a}\right) e^{-jk_0 z} \quad (\text{for } x \geq \frac{b}{2})$$

$$= \frac{A_{10} \pi \Gamma}{a} \exp\left[-\rho_{m0} \left(|x| - \frac{b}{2}\right)\right] \cos\left(\frac{n\pi y}{a}\right) e^{-jk_0 z} \quad (\text{for } x \leq -\frac{b}{2})$$

So, now equating the tangential magnetic fields at the air dielectric interface will obtain the equation, which we obtained similar to the previous case for the LSM even mode. So, equating the tangential magnetic fields at the air dielectric interface, we obtain the relationship

Equating the tangential magnetic fields at the air-dielectric interface, we obtain:

$$A_{10} \sin\left(\frac{q_{m0} b}{2}\right) = A_1$$

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The screenshot shows a presentation slide with the following handwritten text and equations:

LSE even mode :  
In the dielectric region ( $|x| \leq \frac{a}{2}$ ):

$$E_x = 0$$

$$E_y = jk_z' B_{de} \cos(q_{me}'x) \cos\left(\frac{n\pi y}{a}\right) e^{-jk_z'z}$$

$$E_z = -\frac{B_{de} n\pi}{a} \cos(q_{me}'x) \sin\left(\frac{n\pi y}{a}\right) e^{-jk_z'z}$$

$$H_x = \frac{B_{de} (k_0^2 \epsilon_r - q_{me}'^2)}{j\omega\mu} \cos(q_{me}'x) \cos\left(\frac{n\pi y}{a}\right) e^{-jk_z'z}$$

So, now we do the same exercise for the LSE even mode. Let us write down the explicit field equations, in the dielectric and air regions. So, for the LSE even mode, in the dielectric region,

The handwritten equations are identical to those in the screenshot above:

LSE even mode :  
In the dielectric region ( $|x| \leq \frac{a}{2}$ ):

$$E_x = 0$$

$$E_y = jk_z' B_{de} \cos(q_{me}'x) \cos\left(\frac{n\pi y}{a}\right) e^{-jk_z'z}$$

$$E_z = -\frac{B_{de} n\pi}{a} \cos(q_{me}'x) \sin\left(\frac{n\pi y}{a}\right) e^{-jk_z'z}$$

$$H_x = \frac{B_{de} (k_0^2 \epsilon_r - q_{me}'^2)}{j\omega\mu} \cos(q_{me}'x) \cos\left(\frac{n\pi y}{a}\right) e^{-jk_z'z}$$

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$$H_y = \frac{B_0 a' g_m e' n \pi}{j \omega \mu} \sin(g_m' z) \sin\left(\frac{n \pi y}{a}\right) e^{-j k_{z e}' z}$$

$$H_z = \frac{B_0 a' k_{z e}' g_m e'}{\omega \mu} \sin(g_m' z) \cos\left(\frac{n \pi y}{a}\right) e^{-j k_{z e}' z}$$

In the air region ( $|x| \geq \frac{b}{2}$ ):

$$E_x = 0$$

$$E_y = j k_{z e}' B_0 a e' \exp[-\rho_m' (|x| - \frac{b}{2})] \cos\left(\frac{n \pi y}{a}\right) e^{-j k_{z e}' z}$$

$$E_z = -\frac{\rho_m' n \pi}{a} \exp[-\rho_m' (|x| - \frac{b}{2})] \sin\left(\frac{n \pi y}{a}\right) e^{-j k_{z e}' z}$$

$H_y$  is given by

$$H_y = \frac{B_0 a' g_m e' n \pi}{j \omega \mu} \sin(g_m' z) \sin\left(\frac{n \pi y}{a}\right) e^{-j k_{z e}' z}$$

$$H_z = \frac{B_0 a' k_{z e}' g_m e'}{\omega \mu} \sin(g_m' z) \cos\left(\frac{n \pi y}{a}\right) e^{-j k_{z e}' z}$$

Now, in the air region,

In the air region ( $|x| \geq \frac{b}{2}$ ):

$$E_x = 0$$

$$E_y = j k_{z e}' B_0 a e' \exp[-\rho_m' (|x| - \frac{b}{2})] \cos\left(\frac{n \pi y}{a}\right) e^{-j k_{z e}' z}$$

$$E_z = -\frac{\rho_m' n \pi}{a} \exp[-\rho_m' (|x| - \frac{b}{2})] \sin\left(\frac{n \pi y}{a}\right) e^{-j k_{z e}' z}$$

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$$H_x = \frac{B_0 e^{(p_1^2 + k_0^2)} e^{-p_1' (|x| - \frac{b}{2})} \cos\left(\frac{n\pi y}{a}\right) e^{j k_0 z}}{j \omega \mu}$$

$$H_y = \frac{B_0 e^{p_1' n\pi} e^{-p_1' (|x| - \frac{b}{2})} \sin\left(\frac{n\pi y}{a}\right) e^{-j k_0 z}}{j \omega \mu} \quad (\text{for } x \geq \frac{b}{2})$$

$$= -\frac{B_0 e^{p_1' n\pi} e^{-p_1' (|x| - \frac{b}{2})} \sin\left(\frac{n\pi y}{a}\right) e^{j k_0 z}}{j \omega \mu} \quad (\text{for } x \leq -\frac{b}{2})$$

$$H_y = \frac{B_0 e^{p_1' n\pi} e^{-p_1' (|x| - \frac{b}{2})} \sin\left(\frac{n\pi y}{a}\right) e^{-j k_0 z}}{j \omega \mu} \quad (\text{for } x \geq \frac{b}{2})$$

$$= -\frac{B_0 e^{p_1' n\pi} e^{-p_1' (|x| - \frac{b}{2})} \sin\left(\frac{n\pi y}{a}\right) e^{j k_0 z}}{j \omega \mu} \quad (\text{for } x \leq -\frac{b}{2})$$

$$H_z = \frac{B_0 e^{k_0 z} e^{p_1' n\pi} e^{-p_1' (|x| - \frac{b}{2})} \cos\left(\frac{n\pi y}{a}\right) e^{-j k_0 z}}{\omega \mu} \quad (\text{for } x \geq \frac{b}{2})$$

$$= -\frac{B_0 e^{k_0 z} e^{p_1' n\pi} e^{-p_1' (|x| - \frac{b}{2})} \cos\left(\frac{n\pi y}{a}\right) e^{j k_0 z}}{\omega \mu} \quad (\text{for } x \leq -\frac{b}{2})$$

H x is given

$$H_x = \frac{B_{ae} (p_{me}'^2 + k_0^2)}{j\omega\mu} \exp[-p_{me}' (|x| - \frac{b}{2})] \cos(\frac{n\pi y}{a}) e^{-jk_0 z}$$

$$H_y = \frac{B_{ae} p_{me}' n\pi}{j\omega\mu} \exp[-p_{me}' (|x| - \frac{b}{2})] \sin(\frac{n\pi y}{a}) e^{-jk_0 z} \quad (\text{for } x \geq \frac{b}{2})$$

$$= -\frac{B_{ae} p_{me}' n\pi}{j\omega\mu} \exp[-p_{me}' (|x| - \frac{b}{2})] \sin(\frac{n\pi y}{a}) e^{-jk_0 z} \quad (\text{for } x \leq -\frac{b}{2})$$

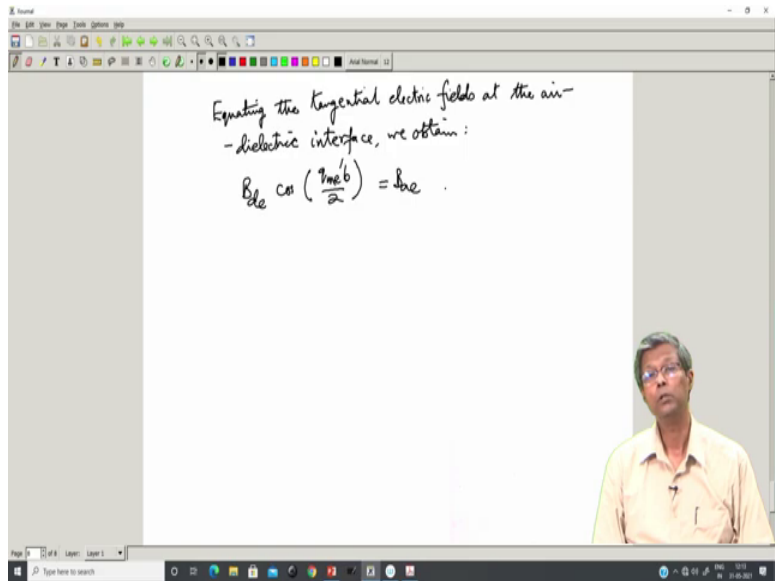
$$H_y = \frac{B_{ae} p_{me}' n\pi}{j\omega\mu} \exp[-p_{me}' (|x| - \frac{b}{2})] \sin(\frac{n\pi y}{a}) e^{-jk_0 z} \quad (\text{for } x \geq \frac{b}{2})$$

$$= -\frac{B_{ae} p_{me}' n\pi}{j\omega\mu} \exp[-p_{me}' (|x| - \frac{b}{2})] \sin(\frac{n\pi y}{a}) e^{-jk_0 z} \quad (\text{for } x \leq -\frac{b}{2})$$

$$H_z = \frac{B_{ae} k_{ze}' p_{me}'}{\omega\mu} \exp[-p_{me}' (|x| - \frac{b}{2})] \cos(\frac{n\pi y}{a}) e^{-jk_0 z} \quad (\text{for } x \geq \frac{b}{2})$$

$$= -\frac{B_{ae} k_{ze}' p_{me}'}{\omega\mu} \exp[-p_{me}' (|x| - \frac{b}{2})] \cos(\frac{n\pi y}{a}) e^{-jk_0 z} \quad (\text{for } x \leq -\frac{b}{2})$$

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So, in this case we equate the tangential electric fields at the air dielectric interface. In order to find the relationship between the coefficients in the dielectric region and the air region. So, equating the tangential electric fields at the air dielectric interface, we obtain



Equating the tangential electric fields at the air-dielectric interface, we obtain:

$$E_{de} \cos\left(\frac{\pi n_e' b}{2}\right) = E_{ae} .$$