

Nonlinear System Analysis
Dr. Ramakrishna Pasumarthy
Department of Electrical Engineering
Indian Institute of Technology, Madras

Lecture – 31
Center Manifold Theorem

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① Jacobian linearization
 $\dot{x} = f(x), \quad x \in \mathbb{R}^n$
 $x: Ax \quad (\text{linear})$
 $A = \frac{\partial f}{\partial x} \Big|_{x: x_{eq}}$

1 Center Manifold Theorem - What is it?

We have seen that when linearization about an equilibrium is illustrative of local nature of an equilibrium point if it is conclusive of either stability or instability of an equilibrium point. But what if the linearization throws out some eigen values with zero real part which does not conclude anything about stability in that particular eigen direction? This can be addressed via the notion of a center manifold.

As an illustrative example, consider the 2D system

$$\begin{aligned}\dot{x} &= ax^3 + xy - xy^2 \\ \dot{y} &= -y + bx^2 + x^2y\end{aligned}$$

Here, linearizing about the origin, we have the linearized system as

$$\begin{aligned}\dot{x} &= 0 \\ \dot{y} &= -y\end{aligned}$$

Hi everybody welcome to this lecture series on Non-linear Systems Analysis. We are in week 9 and we have done a lot of things so far now say starting from properties of non-linear differential equations still bifurcation still stability and so on. So, what I will talk of is one of the concluding lectures of stability analysis before we proceed to other topics.

What we have seen so far when we talk of stability is first we started looking and stability in terms of local linearization which is also called the Jacobi linearization or the Lyapunov's first method or the Lyapunov's direct method. There we just look at eigenvalues of the linearized

system around a particular equilibrium point and remember that stability was always defined with respect to a particular equilibrium point.

So, I do not have just either systems stable or not was defined with respect to a particular equilibrium right for example, the downward position of a pendulum was a stable equilibrium whereas, the upright position was an unstable equilibrium right.

So, we talked about linearization, we talked about defining explicit Lyapunov functions right, so and then we had some methods to derive Lyapunov functions especially the basic ones could be like an energy function or a quadratic Lyapunov function then we had the LaSalle's invariance principle.

A few converse theorems about a Lyapunov stability for non-linear systems we also saw extensively of in the previous lectures of this week on explicit Lyapunov base analysis of linear systems. And different notions of stability and how exponential stability and asymptotic stability are in a way the same concepts when it is a linear time invariant system and what exactly a Lyapunov construction of a Lyapunov function means.

So, in all this analysis or even while during the bifurcation analysis we were interested or we kind of postponed some analysis to the later part, namely when I was looking at what happens when the linearization has a 0 eigenvalue right.

So, as a parameter varied we were looking at the change of or a bifurcation was defined when you are you know when you cross the imaginary axis as your parameter varies when the linearized version or the system a matrix A crosses the has a it causes a 0 eigenvalue right.

So, what to do with those 0 eigenvalues so we will I will briefly give you an introduction to the analysis which is called the center manifold theory, which will give you a method to analyse stability of a non-linear system when the linearized version has a 0 eigenvalue. And of course, a lot of these will be used in for design purposes too, we will not explicitly look at design this course is mostly based on analysis and then the design principles will follow right.

But a lot of systems before you even do the design a lot of time is spent in analysing the systems right. So, we will briefly do the center manifold theorem ok. So, first we did the Jacobi linearization right, so we start with a non-linear system let me take it autonomous \dot{x} equal to $f(x)$ usually a state vector in \mathbb{R}^n .

And around an equilibrium point x^* I just use the same A x^* was A x^* is the linearized version. And then A , was computed as the df by dx evaluated as x equal to some equilibrium points. And then the nature of A told me a lot of things about local stability of the non-linear of the original non-linear system ok. So, most of it was the local nature of an equilibrium point.

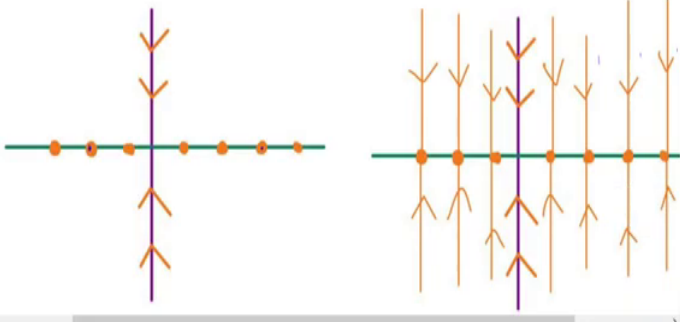
It was concluded we could say if all the eigenvalues of A r on the left half plane then the non-linear system is locally stable right around that equilibrium point. Even if 1 eigenvalue of A is on the right half plane then the non-linear system is also unstable around that equilibrium point and we saw a much of examples. But we really never we never really discussed of what if the linearization has eigenvalues with 0 real part we just not really conclude anything about stability right.

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Here, linearizing about the origin, we have the linearized system as

$$\begin{cases} \dot{x} = 0 \\ \dot{y} = -y \end{cases}$$

It is already in the Jordan form with the x-axis being the eigen direction with zero eigen value and y-axis being the eigen direction with -1 as eigen value. So, we see that in the y-direction, trajectory is stable and converges to the origin. But we are not able to conclude anything about the stability of an equilibrium in the x-direction.



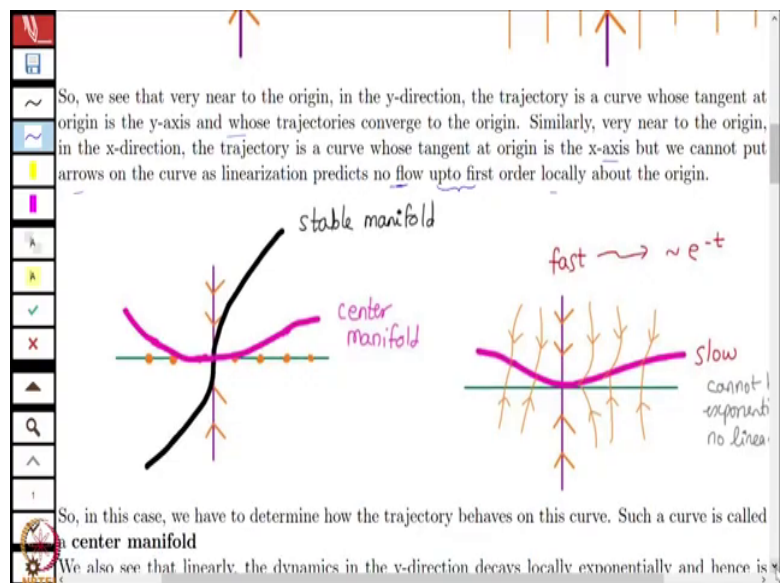
So, we will start motivating with the help of an example and then build up slowly to a theory you will understand what this means as, so now by now we are we know and a to analyse systems in terms of vector fields ok.

So, this nodes and posted on line I did not make proper slides thanks to the lockdown which we are facing here at IIT Madras well almost lockdown ok. So, I start with an example of this form x what is x cube x a example in r^2 and I just linearize about the origin to get something like this right.

So, x dot equal to 0 and y dot is the negative of y right ok, so it is kind of looks in a nice Jordan form ok. So, first looking at y what do I see right, so in the y direction the trajectories are stable and they converge to origin right.

So, $\dot{y} = -y$ will ensure that all $y(t)$ goes to 0 as t goes to infinity. But we really cannot control anything about the stability of an equilibrium in the x direction because if I am just here, I really do not know which direction to move. If I am just here, $x(0) = 0$, there is no direction which I could go.

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So, little observations here which will help us build up towards the theory right, so near the origin in the y direction the trajectory is a curve whose tangent at the origin is the y axis here right so this is how the y axis would be a right.

And similarly near to the origin the trajectory in the x direction the trajectory is a curve whose tangent at origin is the x axis. But we cannot put arrows on the curve as the linearization

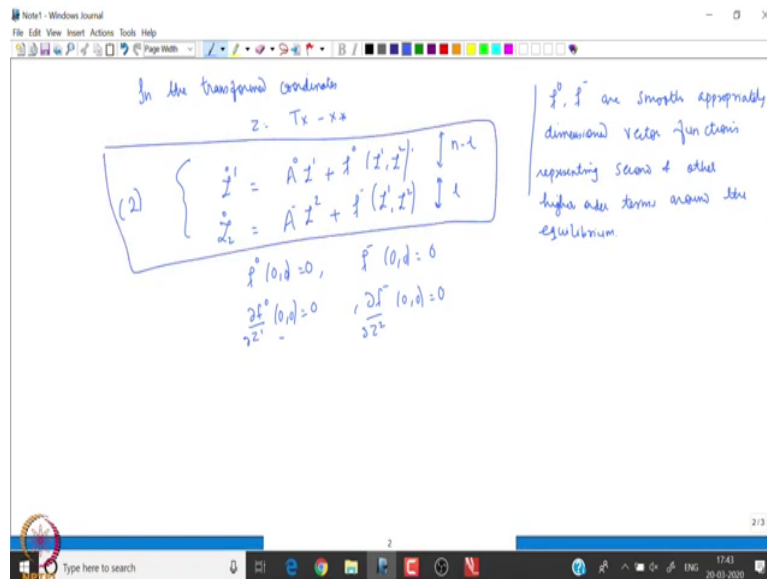
Then it does not really matter if there are otherwise some of the eigen values lie on the origin or say see if I have a system like this which throws up a linear is in this way. I do not really worry about what do I do with this 0 eigenvalues right or maybe there is an eigenvalue here. As because there is already unstable eigenvalue here right, so this is the sigma and the j omega axis ok.

So, we will we will rule out these cases where the system where the linearization throws up a natural instability in terms of an eigenvalue on the right half plane ok. So, therefore you know once I rule out that possibility that A should not have an eigenvalue in the right half plain then the eigenvalues of A can be written as sum disjoint union of all the negative eigenvalues and eigenvalues on the imaginary axis.

So, these are the eigenvalues in the open left half plain then these are the eigenvalues on the imaginary axis ok. So, now, say in \mathbb{R}^n let us assume that there are l eigenvalues in the open left half plane right then there will be $n - l$ eigenvalues which are on the imaginary axis ok.

Now what we know that well we can do some similarity transformation on A of the form $T^{-1} A T$ which will split my A matrix exactly in the form which has I can explicitly write down the eigenvalues on the imaginary axis. And another component which has sorry, eigenvalues in the open left half plane right $z^{-1} A^{-1}$ is $(n - l) \times (n - l)$ matrix and A^{-1} is a $l \times l$ matrix ok.

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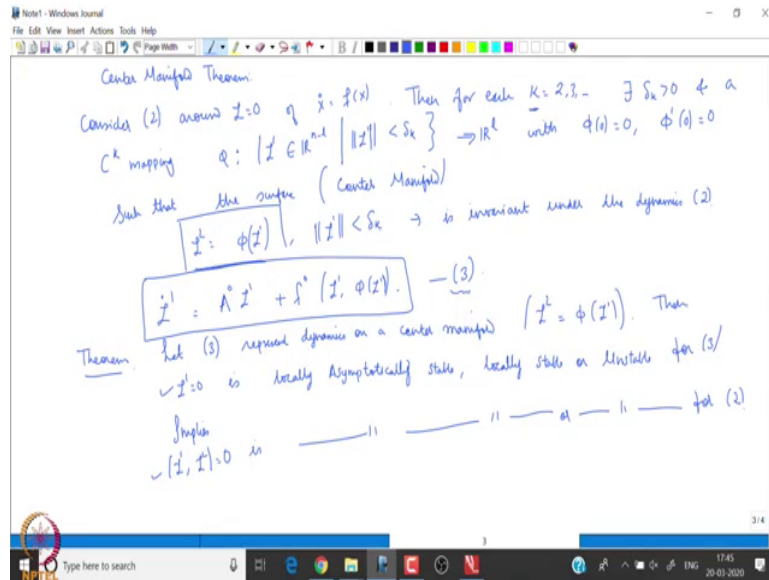
So in the transform coordinates I can write is Z is $T x$ minus x^* being the equilibrium. I can write \dot{x} equal to $A x$ sorry \dot{x} equal to $f x$ of the form say I have I split this into Z^1 dot. I am sorry, this is 1 here Z^1 dot is A naught Z^1 plus f naught Z^1 Z^2 is Z^2 dot is A minus Z^2 plus f with a minus sign.

I mean this is 0 and minus as usual we will denote curve the dynamics corresponding to eigenvalues on the imaginary axis and this will be eigenvalues on the open left half plane ok. So, this will be the first $n-1$ component this will be the last 1 components ok.

And what are this f^1 and f^2 and f^1 and f^2 are certain smooth appropriately dimensioned, vector functions representing the second and other higher order terms around the equilibrium. And also in such a way that $f^1(0,0) = 0$ $f^2(0,0) = 0$

to 0, then df_0 by dz_1 to 0 this like the we are computing the gradients also at the audition ok.

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Now given this first is ok, so let us first write on what is the center manifold theorem. So, let me call this some equation say 2 ok, we start with 2 around the equilibrium point z equal to 0 of an autonomous system $\dot{x} = f(x)$ in the original coordinates. Then for each k is from 2 I will tell you why this restriction on 2 and 3 comes that there exists a number δ_k greater than 0 and a C^k mapping ϕ which goes from z^1 .

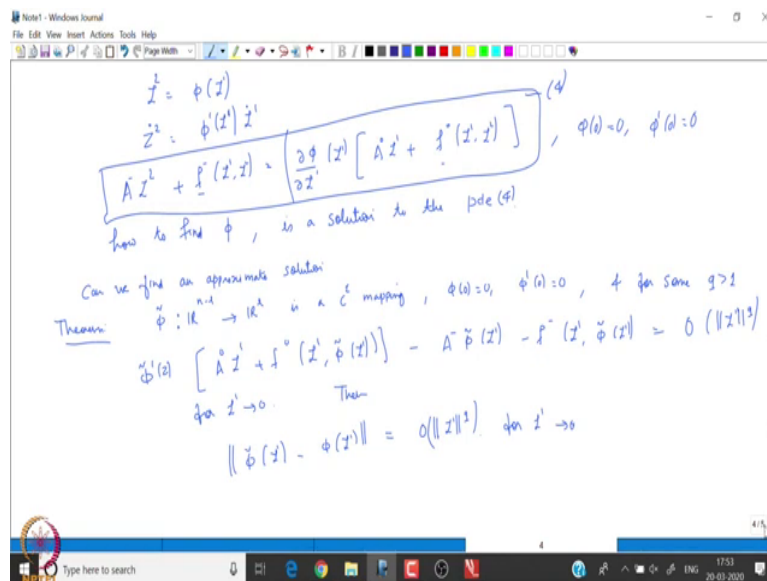
Such that since, the normal z^1 is less than this number δ_k \mathbb{R}^1 with $\phi(0)$ equal to 0 ϕ dash the first derivative is equal to 0. Such that the surface which is called as the center manifold, z^k is ϕ of z^1 . So, I am writing all z^1 z^k in terms of z^1 for all z^1 less than δ_k right and this is such that this is invariant under the dynamics 2 ok.

So, therefore, on this center manifold my dynamics would just now look something like this then \dot{Z}_1 is $A_0 Z_1 + f_0(Z_1, \phi(Z_1))$ ok. Now what is the theorem that then we will tell us about stability ok, so if this say call this \mathcal{F} so you select \mathcal{F} represent the dynamics on the center manifold this how is a center manifold given this is Z_2 is ϕ of Z_1 .

Then Z_1 equal to 0 is locally asymptotically stable or just stable or unstable for \mathcal{F} implies that Z_1, Z_2 equal to 0 the equilibrium of the original system is not locally asymptotically stable, locally stable or unstable right for \mathcal{F} ok. So, I am just looking now at the stability of this equation right and if this is stable then the original system in this original transform system in this coordinates is also stable asymptotically stable or unstable.

Whatever is the property of the system will be translated to the property of the system ok. Now first is so what do I do with this result and how do I actually conclude whether a given system is stable or not. So, first what we have to look at is so I am essentially looking at this dynamics right.

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So, let us write down this here Z_2 is ϕ of Z_1 let us differentiate this I have Z_2 dot is no this is ϕ dash of Z_1 comes Z_1 dot ok. So, what is my Z_2 dot if I just substitute here I have A minus with Z_2 plus f minus Z_1 Z_2 is $d\phi$ by dZ_1 Z_1 times Z_1 dot what is Z_1 dot, this is a naught Z_1 plus f naught Z_1 Z_2 like this right.

So, to arrive at system 3 right, so why do I want to do this because if I just evaluate stability of the system 3 then it will tell me stability or otherwise of the system 3 it will tell me the properties of the original system which is given by the system with equation 2 right.

So, right so this will again together be with conditions of $\phi(0)$ is 0 ϕ dash 0 is also equal to 0 ok. And then the reason we want this function to be of you know for each k is we are excluding like all this linear functions right. So, which where were Z_1 Z_2 could be like linearly re related we see that also when we when we do an example ok. So, first is now how

to find ϕ right so if I label this as equation 4 so this ϕ is a solution to the partial differential equation 4.

I just have equation here, so I know what is $A? B$ minus A naught I know possibly know these functions. So, I just could solve a p d e which may sometimes be analytically difficult to solve p d e. So, can we just look at some approximate methods which can give us some ϕ which is close to the solution or also equal to the solution right ok.

So, can we find an approximate solution ok, so (Refer Time: 26:20) theorem and we will learn how to use this ok. I am skipping the proofs, but we will just as I said we will just know learn how to learn this process ok.

Suppose, I have a instead of ϕ some $\tilde{\phi}$ which goes from \mathbb{R}^{n-1} to \mathbb{R}^1 right and it is a C^2 mapping ok. Which satisfies $\tilde{\phi}$ naught equal to $\tilde{\phi}$ equal to 0 $\tilde{\phi}$ dash at 0 also equal to 0 right and for some q greater than 1 suppose that this $\tilde{\phi}$ of Z times plus.

So, it 1 comma this I am writing everything in the new in the approximated ϕ denoted by $\tilde{\phi}$ minus A minus $\tilde{\phi}$ Z^{-1} , this is of some order or Z^{-1} to the power q for Z^{-1} turn into 0. Then the approximated $\tilde{\phi}$ and the additional $\tilde{\phi}$ sorry additional ϕ are related this way.

The order of this one for is that 1 going to 0 ok, so this is so what we are trying to do is try trying to find a way to approximate the file which may sometimes be analytically difficult to compute ok. So, I just quickly do with a help of an example and then well post a couple of examples online so that will help you practice a little more problems.

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$$\frac{d}{dt} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} x_1 x_2 + x_1^3 + x_1 x_2^2 \\ -x_2 - 2x_1^2 + x_1^2 x_2 \end{pmatrix}$$
 Linearization around $(0,0)$

$$x_2 = \phi(x_1), \quad \phi(0) = 0, \quad \phi'(0) = 0$$

$$\ddot{\phi}(x_1) = \ddot{\phi}_2 x_1^2 + \ddot{\phi}_3 x_1^3 + \dots$$

$$(\ddot{\phi}_2 + 2)x_1^2 + \ddot{\phi}_3 x_1^3 + 0(x_1^4) = \ddot{\phi}'(x_1) \left[x_1 \ddot{\phi}(x_1) + x_1^3 + x_1 \ddot{\phi}^2(x_1) + \ddot{\phi}(x_1) + 2x_1^2 + x_1^2 \ddot{\phi}(x_1) \right]$$

$$x_2 = \phi(x_1) = -2x_1^2 + 0(x_1^3) \quad \text{Sub into the dynamics of the center manifold}$$

$$\frac{d}{dt} x_1 = x_1 x_2 + x_1^3 + x_1 x_2^2 = x_1(-2x_1^2) + x_1^3 + 0(x_1^4) = -x_1^3 + 0(x_1^4)$$

for this $x_1 = 0$ is a locally asymptotically stable eq. point. $\therefore (0,0)$ is also a locally A.S. eq. point

So start with a simple system which again in \mathbb{R}^2 I have $x_1 x_2$ plus x_1 cube plus $x_1 x_2$ square I have minus x_2 here minus twice x_1 square plus x_1 square x_2 ok. So, I am just looking at linearization around 0 comma 0 and then now you see that it is it has a 0 eigenvalue and I am looking for the center manifold of the form ϕ_2 is ϕ of x_1 ok, again with $\phi(0)$ is 0 ϕ' of 0 also equal to 0 ok. So, let me just say that ϕ tilde as a function of x_1 is this C^2 function twice differentiable function ϕ tilde 2 let me call this coefficient as ϕ tilde 2 x_1 square.

Some ϕ tilde 3 x_1 cube and so on, so what I do is now take this ϕ and then write down with this ϕ tilde this is my approximate ϕ tilde and write an equation like this ok. So, should be simple right so I will have ϕ tilde. So, this expression on the right hand side so let

us write the expression on the left first, this is $\tilde{\phi}^2 + 2x_1^2 + 5\tilde{\phi}x_1^3 + \dots$ plus some other terms of order 4 ok.

So, what I am doing is I am just taking this $\tilde{\phi}$ of x_1 and this substituting it here right and what are the exact dynamics here those are given by this x_2 terms here. And you are substituting this $\tilde{\phi}$ approximate ϕ over here what does this equate to? This equates to this expression on the right here.

So, that will look something like $\tilde{\phi} \dot{x}_1$ times x_1 right I just substitute it here x_1 then x_2 is ϕ of x_1 what is ϕ of x_1 ? Is to write down this approximation here $\tilde{\phi}x_1 + \dots$ plus I have $x_1^3 + \dots$ plus $\tilde{\phi}^2x_1 + \dots$ ok.

So, this is just so very simple thing of substituting this into this one into or into this equation number 4 as this is a little lengthy process, but it is purely mechanical. So, if you do all once we do all the manipulations what we find is the following at ϕ of x_1 is simply $-2x_1^2 + o(x_1^3)$ ok. Now if I substitute this into the dynamics the dynamics of the center manifold ok.

So, what do I get is \dot{x}_1 from this expression, so $\frac{d}{dt}x_1$ is $x_1x_2 + x_1^3 + \dots$ right ok. So, what is x_2 ? Is ϕ of x_1 and from here so this ϕ of x_1 with x_2 . So, I will have x_1x_2 is ϕ of x_1 that will be $-2x_1^2 + x_1^3 + \dots$ plus other terms of order 4 that will give me $-x_1^3 + \dots$ plus other terms of order 4 and above ok.

Now, for this right $x_1 = 0$ is because of this $-x_1^3$ is locally nothing to declare stable equilibrium point. And therefore, for the original system $(0,0)$ is also locally asymptotically stable equilibrium point ok. So, this is a little illustration of how to analyse systems for which the linearization throws up a 0 eigenvalue or eigenvalues on the imaginary axis y are the center manifold theory.

We do not really go deep into the details of what if the systems are of order 4 and higher, but just to give you a little glimpse of how will you handle systems for which there are eigenvalues

on the imaginary axis on of during Jacobi could be a linearization right. So, well put up some problems for you to solve might be there might be helpful for you to answer more solve more problems or questions by yourself and also help you answer the assignment problems.

So you can always post questions online if you have. So, this is about the theory on center manifold a very brief introduction thanks for listening.