

Nonlinear Vibration
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Lecture - 34

Nonlinear dynamics of turning operation with delay and internal resonance

Welcome to today class of Non-linear Vibration. So, we are continuing in module 9. So, where we have already discussed about 2 applications. So, today we are going to study about the third applications that is on a manufacturing process.

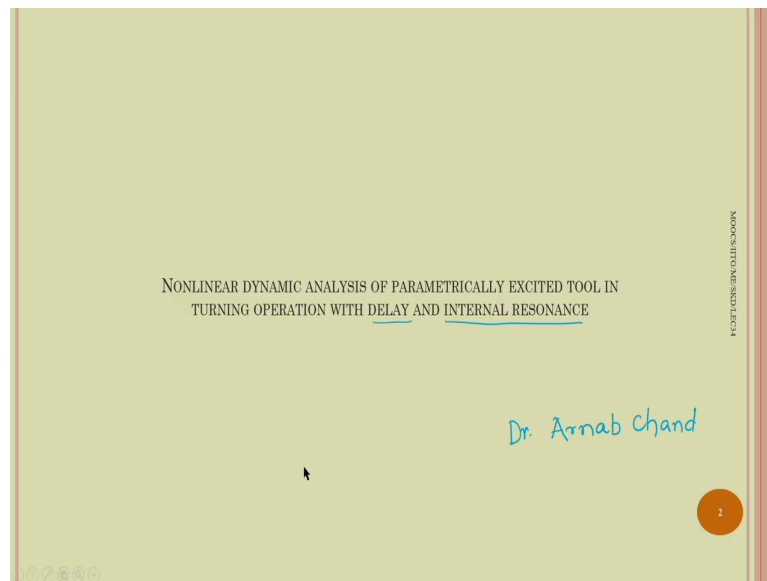
So, we will take the turning operation as the application and we will study. So, if the vibration of the tool or the vibration of the workpiece to be non-linear and we will derive these equation on motion and see how to solve these equation motion.

So, in the previous 2 weeks. So, you have studied regarding these energy harvester and also vibration observer. So, where we have studied regarding the passive and active vibration observer. So, in case of energy harvester. So, we have seen how this piezoelectric excited cantilever beam can be used as a, energy harvester.

So, in that case. So, we have taken both external and internal resonance conditions and we have studied the equation governing equation of motion and then we have solved those governing equation of motion and we have seen how we can obtain this voltage from this vibration or we can harvest the energy from this ambient condition.

So, in case of vibration observer. So, which can find many applications in industry. So, we have seen the passive type of vibration observer and active type of vibration observer. So, today we will extend our analysis to a manufacturing process. So, where we are going to study regarding the vibration of the tool and work piece in case of the turning operation. So, the first work I will show you.

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So, this is part of the PhD work of my PhD student Doctor Arnab Chand and some of these work have been already published in the journal of sound and vibration and so, these work.

So, you just see regarding these non-linear dynamic analysis of parametrically excited tool in turning operation with delay and internal resonance. So, here we are going to consider the delay and also the internal resonance condition and we will find the instability regions to study the conditions in which there will be no vibration.

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EQUATION OF MOTION OF TOOL

➤ The Equations of motion of tool in cutting force and feed force direction are as follows respectively

$$m\ddot{x}(t) + c_x\dot{x}(t) + k_{cx}x(t) + k_{cy}y^2 = -F_x - k_y w \dot{y}^2$$

$$m\ddot{y}(t) + c_y\dot{y}(t) + k_{cy}y(t) + k_{cx}x^2 = -F_y = -k_y w \dot{x}^2$$

where, $h = vt + y(t) - y(t - \tau_c)$

➤ But, because of vibration in cutting force direction $\dot{h} = 2\dot{v}t + x(t) - x(t - \tau)$

➤ Thus, re-writing the chip thickness

So, one can write, $h = vt(x) + y(t) - y(t - \tau(x))$

$$\ddot{x}(t) + \frac{c_x}{m}\dot{x}(t) + \frac{k_{cx}}{m}x(t) + \frac{k_{cy}}{m}y^2 = \frac{k_{cy}w}{m}(vt(x) + y(t) - y(t - \tau(x)))^2$$

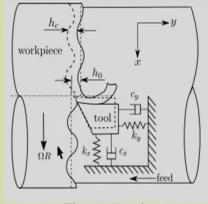
$$\ddot{y}(t) + \frac{c_y}{m}\dot{y}(t) + \frac{k_{cy}}{m}y(t) + \frac{k_{cx}}{m}x^2 = -\frac{k_{cx}w}{m}(vt(x) + y(t) - y(t - \tau(x)))^2$$


Figure : Turning operation modeled as 2-DOF system

Assumption is made that tool is slender and workpiece is rigid.

So, this is the turning operation. So, you can see the turning operation. So, in this turning operation. So, we have a work piece and this is the tool. So, the tool holder can be modeled or tool is supported the supporting condition of the tool can be model by using the spring and dumper system. So, it can be modeled. So, here we are assuming the workpiece to be rigid and the tool to be flexible.

So, as we are considering a flexible tool. So, here these flexibility is represented by using the spring constant or spring and damper in both x and y direction. So, in x direction. So, we have taken these vertical direction. So, this is x direction and the horizontal direction as y direction in this case. So, here. So, the spring in x direction is taken to be k x and the dumper is taken to be c x and in y direction it is k y and c y.

As the work piece is rotating. So, we assume the workpiece is rotating with angular velocity ω . So, there is initial undulation in the cutting. So, you can see. So, initially. So, due to this undulation. So, this is the chip. So, you can see this is the chip and so, the chip thickness is represented by h .

So, we will see the thing. So, here initially it is h_0 and after some time. So, it will be h_c . So, we have to find how much force is acting on the by the tool on the workpiece or the how much force is exerted by the work piece on the tool in a reverse direction.

So, this is the feed direction. So, the equation of motion of the tool can be retained in both x and y direction. So, in the cutting force direction and feed direction. So, this is the feed direction y is the feed direction and this the cutting force direction. So, x is the cutting force direction and y is the feed direction. So, in cutting force direction the equation of motion can be retained equal to $m \ddot{x}$.

So, in x . So, this is the, this is the cutting force direction, x is the cutting force direction. So, if the displacement is x in cutting force direction. So, it can be retained as $m \ddot{x} + c \dot{x} + kx = F$.

So, we have written this is $kx + k_1 x^3$ into x plus k . So, we are taking a cubic. So, you are modeling this spring as a non-linear spring. So, generally in most of the literature you can find they have taken this linear spring. So, here it is taken to be a non-linear spring. So, that is why this kx is represented by $kx + k_1 x^3$ is the linear part kx and the non-linear part is written $k_1 x^3$.

So, we are taking a cubic non-linear spring. So, this way we can model the spring force as, $kx + k_1 x^3$ into x plus $kx + k_1 x^3$. So, this will be equal to the cutting force that is F_c . So, cutting force F_c , these cutting force can be retained by using this equation that is $kx + k_1 x^3 = F_c$.

So, similarly in case of y direction. So, this is the y direction this oriental direction is the y direction. So, in this or the feed direction. So, in the feed direction. So, the equation motion can be written as $m \ddot{y} + c \dot{y} +$ similarly here in y direction also we have taken we have taken these cubic nonlinearity.

So, that is why it can be written $k y + y^3$. So, this will be equal to minus $f y$. So, would be feed direction. So, the work piece will absorb a force on the. So, on the tool. So, that is why it is written as minus $F y$ that is equal to minus $K y w$ into h to the power q .

So, where h . So, that is the chip thickness can be written as. So, this is equal to $b t + y t - y t - \tau$. So, here you just see this is $t - \tau$. So, this is the τ is the time delay. So, delay term is used here. So, before previous cut let the previous cut in y direction, previous cut in y direction what is the displacement and this present cut what is the displacement and into. So, this is the $v t$, but because of the vibration in the cutting force direction. So, $\omega R \tau$ will be equal to $2 \pi R + x t - x t - \tau$.

So, substituting these 2 in this equation and writing that h equal to $v \tau + x t + y t - y t - \tau + x t$. So, we can write these equation in this form that is $\ddot{x} + c \dot{x} + k x + k y^3$ by $m x$ cube. So, this is $k x^3$. So, this is not y. So, this is $k x^3$ by $m x$ cube.

So, this will be equal to $k x w$ by m in to $v \tau + x t + y t - y t - \tau + x t$ to the power q . So, similarly in the y direction equation can be written as $\ddot{y} + c \dot{y} + k y + y^3$ by $m y$ cube equal to minus $k y w$ by m into similarly we can put these equation w and h q in this form.

So, h replacing h in this way. So, it will be equal to $v \tau + x t + y t - y t - \tau + x t$ to the power q . So, here the tool is assumed to be rigid and the tool is assumed to be flexible and the work piece assume to be rigid. So, that is why we got these equations. So, here we have taken the delay also into account to find what is the chip thickness actually due to

undulation. So, it will depend on the previous cut. So, the chip thickness will depend on the previous cut. So, that is why. So, taking that into account. So, it can be written in this form.

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EQUATION OF MOTION OF TOOL

➤ The solution of the equations can be assumed as

$$x = x_0 + \zeta \text{ and } y = y_0 + \eta$$

$$\ddot{x}_0 + \zeta + \frac{c_x}{m} \dot{x}_0 + \frac{c_z}{m} \dot{x}_0 + \frac{k_{x1}}{m} x_0 + \frac{k_{x2}}{m} \zeta + \frac{k_{x3}}{m} (\zeta + \zeta)^3 = \frac{K_x w}{m} (\tau(x_0 + \zeta) + y_0 + \eta - y_0(t - \tau_0) - \eta(t - \tau_0))^q$$

➤ If we substitute, $x = x_0$ in the equation, one will obtain,

$$\ddot{x}_0(t) + \frac{c_x}{m} \dot{x}_0(t) + \frac{k_{x1}}{m} x_0(t) + \frac{k_{x2}}{m} x_0^3 = \frac{K_x w}{m} (\tau(x_0) + y_0(t) - y_0(t - \tau_0))^q$$

As $\ddot{x}_0(t) = \dot{x}_0(t) = y_0(t) - y_0(t - \tau_0) = 0$ because of static deflection,

➤ Due to steady state condition, one will get,

$$\zeta + \frac{c_x}{m} \dot{\zeta} + \frac{c_z}{m} \dot{\zeta} + \frac{3k_{x2}}{m} \zeta^2 + \frac{k_{x1}}{m} \zeta + \frac{3k_{x3}}{m} x_0 \zeta^2 + \frac{k_{x3}}{m} \zeta^3 + \frac{K_x w}{m} (\tau)^q = \frac{K_x w}{m} (\tau(x_0) + \eta - \eta(t - \tau_0))^q$$

$$= \frac{K_x w}{m} (\tau(x_0))^q + \frac{K_x w \eta}{m} (\tau(x_0))^{q-1} \left(\frac{K(t) - \zeta(t - \tau_0))}{R} + \eta - \eta(t - \tau_0) \right)$$

$$\zeta + \frac{c_x}{m} \dot{\zeta} + \left[\frac{c_x}{m} + \frac{3k_{x2}}{m} x_0 \right] \dot{\zeta} + \frac{3k_{x2}}{m} x_0 \zeta^2 + \frac{k_{x1}}{m} \zeta^3 = \frac{K_x w \eta (2\pi R)^{q-1}}{m} \left[\rho \right]^{q-1} \rho \left[\zeta - \zeta(t - \tau_0) \right] + \eta - \eta(t - \tau_0) \quad (34)$$

Where, $\rho = \frac{\tau}{R}$

So, now assuming this x to be equal to x 0 plus zeta and y equal to y 0 plus eta. So, at time t for example, x equal to x 0 and time t y equal to y 0. So, at any time t. So, these x can be written as x 0 plus zeta and y equal to y 0 plus eta.

So, that is why this equation can be now for the written in this form, that is x double dot can be written as x 0 double dot plus zeta double dot and x 0 double dot plus zeta double dot plus c by m, c x by m x 0 dot plus c x by m zeta dot k x 1 by m x 0 plus k x 1 by m zeta plus k x 3 by m into x plus zeta cube equal to k x w by m.

So, $v \tau$ into x_0 plus ζ plus y_0 plus η minus $y_0 t$ minus τ_0 minus η into t minus τ_0 to the power q . So, this q actually the number. So, depending on the conditions of the cutting operations. So, from experimentally this number is can be determined. So, for different materials and cutting conditions. So, q will be different.

So, in the numerical analysis we will see what is the value of q is taken in this case. So, if we will substitute x equal to x_0 in this equation we will obtain. So, for. So, this if x_0 is the steady state condition. So, if we substitute x equal to x_0 then this equation becomes. So, this $x_0 \ddot{x} + e \text{ by } m \times x_0 \dot{x} + k \text{ by } m \times x_0 + k_3 k \times x_0^3 \text{ by } m \times x_0^3 = k \times w \text{ by } m$ into. So, these parameter to the power q .

So, as in case of the static deflection taking x_0 as the static deflection. So, this part will be equal to 0. So, then. So, now, neglecting this part. So, the remaining part can be written as $\zeta \ddot{\zeta}$. So, due to steady state condition. So, one can take this is equal to $k \text{ by } m \times x_0 t + \alpha \frac{1}{3} \text{ by } m \times x_0^3 = k \times w \text{ by } m \text{ v } \tau$ to the power q . So, this is the condition one can take.

So, now. So, the equation will reduce to this form that is $\zeta \ddot{\zeta} + c \text{ by } m \zeta \dot{\zeta} + k \times \text{ by } m \zeta + 3$. So, you just see when you expand these term that is x plus ζ cube. So, here you have put x equal to x_0 plus ζ cube.

So, you just see. So, if you expand these thing then this becomes $x_0^3 + \zeta^3 + 3 x_0^2 \zeta + 3 x_0 \zeta^2$. As ζ is very small the ζ^3 and higher order terms can be neglected or if you want to keep also you can keep these up to cubic non-linear term and so, there will be a term of $3 x_0^2 \zeta$ and another term with $3 \zeta^2 x_0$.

So, the perturbation will contain all these terms. So, that is why. So, this additional term you can see. So, $\zeta \ddot{\zeta} + c \text{ by } m \zeta \dot{\zeta} + k \text{ by } m \zeta + 3 k \times x_0^2 \zeta + k \text{ by } m \zeta^2 + 3 k \times x_0 \zeta^2 + k \times w \text{ by } m \text{ v } \tau$ to the power q .

So, it will be equal to $k \times w$ by m . So, $v \tau \times 0 t$ plus η into ηt minus $\tau 0$ to the power q . So, this is due to the time delay, delay term. So, now, this can be written as, $k \times w$ by m . So, $v \tau$, $v \tau \times 0 t$ to the power q plus $k \times w$ by m . So, this is equal to or $v \tau \times 0 t$ to the power q minus 1 into this the term it can be written.

So, now we can give or rewrite these equation in this form that is $\zeta \ddot{\zeta} + c$ by m $\zeta \dot{\zeta}$ plus. So, we can combine the term with ζ . So, this becomes $k \times w$ by m plus $3 k \times w$ by m into $x 0$ square ζ . So, you just see. So, now, the coefficient of ζ . So, equal to $k \times w$ by m plus $3 k \times w$ by $m \times 0$ square. So, it depend on the initial condition that is initial $x 0$.

So, this will be the coefficient. So, this coefficient is generally the non-linear natural frequency of the system. So, that is ω_n square. So, plus $3 k \times w$ by $m \times 0$ ζ square plus $k \times w$ by m ζ cube. So, we have a quadratic term and a cubic term and other parts we have taken to right hand side. So, which can be acting as a forcing term. So, here we can take this ρ equal to v by ωR .

So, these ωR into ρ equal to v we have taken. So, that is why this ρ equal to v by ωR . So, it is written in terms of now it is written in terms of ρ . So, right hand side is written in terms of ρ .

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EQUATION OF MOTION OF TOOL

$$\ddot{a}_{11} = \frac{k_{11}}{m} + \frac{3k_{13}}{m} \ddot{a}_{21} + \frac{k_{11}}{m} + \frac{3k_{13}}{m} \ddot{a}_{21} \quad \bar{t} = \sqrt{\alpha_{11}} t, \quad \bar{c}_{11} = \frac{\bar{a}_{11}}{\alpha_{11}}, \quad \bar{c}_{21} = \frac{\bar{a}_{21}}{\alpha_{11}}, \quad \bar{\mu}_1 = \frac{c_1}{2\epsilon \sqrt{\alpha_{11} m}}$$

$$\bar{\mu}_2 = \frac{c_2}{2\epsilon \sqrt{\alpha_{21} m}}, \quad \bar{a}_{12} = \frac{3k_{13} y_0 c_1}{\alpha_{11} m \epsilon^2}, \quad \bar{a}_{13} = \frac{k_{13} \rho_0^2 c_1}{\alpha_{11} m \epsilon^2}, \quad \bar{a}_{22} = \frac{3k_{23} y_0 c_2}{\alpha_{11} m \epsilon^2}, \quad \bar{a}_{23} = \frac{k_{23} \rho_0^2 c_2}{\alpha_{11} m \epsilon^2}, \quad \bar{\Omega} = \frac{\Omega}{\sqrt{\alpha_{11}}}, \quad \bar{x} = \alpha x, \quad \bar{y} = \alpha y,$$

□ Removing the bar on the variables and arranging Eqs.

$$\ddot{x} + 2\bar{\mu}_1 \omega_{n1} \dot{x} + \omega_{n1}^2 x + \epsilon^2 \bar{a}_{12} x^2 + \epsilon^2 \bar{a}_{13} x^3 = \bar{c}_{11} \rho^{\epsilon-1}(\eta(t) - \eta(t-t_0)) + \frac{K_1}{K_2} \rho^{\epsilon}(\xi(t) - \xi(t-t_0))$$

$$\ddot{y} + 2\bar{\mu}_2 \omega_{n2} \dot{y} + \omega_{n2}^2 y + \epsilon^2 \bar{a}_{22} y^2 + \epsilon^2 \bar{a}_{23} y^3 = -\bar{c}_{21} \rho^{\epsilon-1}(\eta(t) - \eta(t-t_0)) - \bar{c}_{22} \rho^{\epsilon}(\xi(t) - \xi(t-t_0))$$

Assume, $\rho = \rho_0(1 + r \sin \Omega t)$

$$\ddot{x} + 2\bar{\mu}_1 \omega_{n1} \dot{x} + x + \epsilon^2 \bar{a}_{12} x^2 + \epsilon^2 \bar{a}_{13} x^3 = \bar{c}_{11} k_{11} + \epsilon q_{11} \sin \Omega t (y(t) - y(t-t_0)) + \epsilon^2 k_{12} [1 + \epsilon q_{22} \sin \Omega t] (x(t) - x(t-t_0))$$

$$\ddot{y} + 2\bar{\mu}_2 \omega_{n2} \dot{y} + y + \epsilon^2 \bar{a}_{22} y^2 + \epsilon^2 \bar{a}_{23} y^3 = -\bar{c}_{21} k_{21} + \epsilon q_{21} \sin \Omega t (y(t) - y(t-t_0)) - \epsilon^2 k_{22} [1 + \epsilon q_{22} \sin \Omega t] (y(t) - y(t-t_0))$$

where $q_{11} = r(t-1)/\epsilon; q_{22} = r q_1 / \epsilon; k_1 = k_1 \omega_0 (2\pi R)^{\epsilon-1} / k_2; k_2 = k_2 \omega_0 (2\pi R)^{\epsilon-1} / m; \omega_0^2$

$$k_1 = \frac{k_1}{k_2}; k_{11} = \frac{K_1 \rho_0^{\epsilon-1}}{K_2 \epsilon}; k_{12} = \frac{K_1 \rho_0^{\epsilon}}{K_2 \epsilon^2}; k_{21} = k_{11} k_2; k_{22} = k_{12} k_2; \mu_1 = \bar{\mu}_1;$$

Handwritten notes: $\bar{t} = \omega_{n1} t$, $dde 23$

So, these are the parameters one can use. So, for example, one can write these complicated equation in a simplified form. So, one can use these parameters for example, alpha 11 one can use is equal to k x 1 by m plus 3 k x 3 by m x 0 square. So, it will represent the natural frequency of the systems square of the natural frequency of the system similarly alpha 2 1 can be taken k y one by m plus 3 k y 3 by m into y 0 square.

So, these non-dimensional time can be written as these t bar equal to root over alpha 11. So, that is omega t. So, for example, generally we take these t bar equal to these is non dimensional time equal to omega into t. So, this for the x direction we have taken this is equal to omega 11 or alpha 11.

So, this is radian per second into radian per second into second. So, this becomes radian. So, that is the angle. So, we are taking this non-dimensional time t bar. So, t bar equal to root

over $\alpha_{11} t$ then τ . So, this is the delay time also we can non dimension h that thing. So, delay time also can be written $\sqrt{\alpha_{11}} \tau \omega_{11}^2$ equal to. So, this is the non-dimensional ω_{11} we are taking.

So, ω_{11}^2 equal to α_{11} by α_{11} . So, you just see we have taken it in such way that this term becomes 1. So, then ω_{21} this is α_{21} by α_{11} . So, μ_1 . So, this is to represent the damping we are taking this μ term. So, μ_1 bar equal to c_x by 2ϵ .

So, here we are keeping or using the bookkeeping parameter ϵ that is why we have this ϵ here. So, μ_1 bar equal to c_x by $2 \epsilon \sqrt{\alpha_{11}}$ bar by m . Similarly, μ_2 bar equal to c_y divided by $2 \epsilon \sqrt{\alpha_{21}}$ m .

Similarly, α_{12} can be written, α_{13} and α_{22} α_{23} ω bar this is the frequency non dimensional frequency of rotation. So, rotational frequency we have taken equal to rotational frequency you have taken equal to capital ω . So, we can non-demisionalize this frequency that is ω bar equal to ω by $\sqrt{\alpha_{11}}$ bar.

So, then this link parameters x by l we have taken z_x bar and y by l also we have taken z_y bar. So, now, removing the bar from these. So, now, using these thing the equations can be written in using a bar form. So, that removing this bar we can have these two equation. So, one in ζ one in ζ and other in η ζ is in x direction and η in y direction.

So, our equation now reduced to $\zeta \ddot{} + 2 \epsilon \mu_1 \omega_{11} \zeta \dot{} + \omega_{11}^2 \zeta + \epsilon^2 \alpha_{11} x^2 + \epsilon^2 \alpha_{13} x^3$ equal to k_1 by $k_r \rho$ to the power q minus 1 into ηt minus ηt minus τ plus k_1 by $k_r \rho$ to the power cube ζ minus ζt minus τ 0.

Similarly, ζ equation in y direction our equation is reduced to $\eta \ddot{} + 2 \epsilon \mu_2 \omega_{21} \eta \dot{} + \omega_{21}^2 \eta + \epsilon^2 \alpha_{22} y^2 + \epsilon^2 \alpha_{23} y^3$ plus

$\epsilon^2 \alpha^3 y^3 = -k_1 \rho^q \int \eta t - \eta t - \tau_0 - k_1 \rho^q \zeta t - \zeta t - \tau_0$.

So, here. So, this ρ is a rotational is a function of the rotational speed. So, that thing can be written as a constant speed plus a variable speed. So, if we are assuming the variation to be sinusoidal, then we can write these ρ equal to $\rho_0 + r \sin \omega t$.

So, here either. So, if it is rotating with a constant speed then ρ will be equal to ρ_0 , but when we are assuming the rotation to be varying then that variation we can write in a sinusoidal form. So, that is why it will reduce to. So, ρ can be written as $\rho_0 + r \sin \omega t$.

So, here you just see. So, we can write now this equation in this form. So, it can be written $x'' + 2\epsilon \mu \omega x' + x = \epsilon^2 \alpha^3 y^3$. So, this ζ is replaced by x and this η is replaced by y for better understanding. So, now, it is written in this form that is $x'' + 2\epsilon \mu \omega x' + x = \epsilon^2 \alpha^3 y^3$. So, these two equations you can check. So, here we got this $\omega^2 = \dots$ as we are writing in this form.

So, $x'' + 2\epsilon \mu \omega x' + x = \epsilon^2 \alpha^3 y^3$ and $x + \epsilon^2 \alpha^3 y^3 = \epsilon k_1 + \epsilon q \sin \omega t$ into $y t - y t - \tau_0 + \epsilon^2 \alpha^3 k_1 + \epsilon q \sin \omega t$ minus $x t - x t - \tau_0$.

Similarly, in this y direction this is the equation can be written. So, here we have taken this q equal to $R^q - 1$ by ϵq R^q by ϵk_1 equal to $k_1 y \omega^2 \pi r$ to the power $q - 1$ divided by k_1 and that is equal to $k_1 y \omega^2 \pi R$ to the power $q - 1$ divided by $m \omega^2$.

Similarly, $k_1 r$ is taken to be $k_1 y$ by $k_1 x$ k_1 equal to $k_1 \rho_0^q$ to the power $q - 1$ by $k_1 r$ ϵk_1^2 equal to $k_1 \rho_0^q$ to the power q by $k_1 r$ $\epsilon^2 k_1^2$ equal to k_1 into $k_1 r k_1^2$ equal to k_1^2 into $k_1 r$ and μ equal to μ equal to. So, here taken μ equal to μ

equal to μ^2 . So, considering symmetric. So, we can write down these two equations in this form.

So now. So, as we got the governing equation of motion. So, to see these equations are not similar to the, what we have studied in case of the damping equation particularly in the right hand side the equation the forcing part in the damping equation is not similar to the what we have written it here.

Here the forcing is slightly different it is written in a slightly different way what. So, the left hand side contains these quadratic and cubic nonlinearity. So, if it is equal to 0. So, we can have a damping equation with both quadratic and cubic non-linearity. Actually you know in a governing equation. So, you can have two-part complementary part and the particular integral.

So, the complementary part also will contribute to the vibration of the system that is the free vibration of the system and these particular integral will be due to the forcing part of the system. So, now, these now we have to solve these equations you just say this equation is written in the form of a delay differential equation. So, here the delay part is considered in this equation. So, this is the delay differential equation.

So, you can solve this delay differential equation using these `dde23`. So, using `dde23` command of the MATLAB or you can write your own equation solver to solve these things also we have discussed previously that by using this Runge Kutta method also you will be able to find the solution or solution of this non-linear governing equation of motion. So, let us take the example or how to solve it using this method of multiple scale.

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PERTURBATION METHOD : METHOD OF MULTIPLE SCALES (MMS)

Following the multiple scale method, the solution can be represented by an expansion having the form

$$x = x_{11}(T_0, T_1, \dots) + \epsilon x_{12}(T_0, T_1, \dots) + \epsilon^2 x_{13}(T_0, T_1, \dots)$$

$$x_2 = x_{21}(T_0 - \tau, T_1 - \epsilon\tau, \dots) + \epsilon x_{22}(T_0 - \tau, T_1 - \epsilon\tau, \dots) + \epsilon^2 x_{23}(T_0 - \tau, T_1 - \epsilon\tau, \dots) \quad \checkmark$$

$$y = y_{11}(T_0, T_1, \dots) + \epsilon y_{12}(T_0, T_1, \dots) + \epsilon^2 y_{13}(T_0, T_1, \dots)$$

$$y_2 = y_{21}(T_0 - \tau, T_1 - \epsilon\tau, \dots) + \epsilon y_{22}(T_0 - \tau, T_1 - \epsilon\tau, \dots) + \epsilon^2 y_{23}(T_0 - \tau, T_1 - \epsilon\tau, \dots) \quad (42)$$

By introducing the new independent variables according to

$$T_n = \epsilon^n t$$

Time derivatives along different time scales lead to the

$$d/dt = D_0 + \epsilon D_1 + \dots \quad \text{and} \quad d^2/dt^2 = D_0^2 + 2\epsilon D_0 D_1 + \dots$$

$$T_n = \epsilon^n t$$

$$T_0 = t \quad \checkmark$$

$$T_1 = \epsilon t$$

$$T_2 = \epsilon^2 t$$

$$t \rightarrow t - \tau$$

$$T_0 = T_0 - \tau = \epsilon T_1 - \epsilon\tau = T_1 - \epsilon\tau$$

PERTURBATION METHOD

So, as you have seen due to the presence of the nonlinearity. So, it will not yield any close form solution. So, as it is not or we are not having any close form solution. So, we can go for some approximate solution. Approximate solution can be obtain by using many different methods such as methods of averaging, method of (Refer Time: 27:59) Poincare method, homotropy method and normal firm methods method of multiple scales also.

So, here method of multiple scales is used to solve these equation. So, here for the x and y we can take the conventional expansion. So, here of two second order expansion is taken. So, x is taken to be $x_{11}(T_0, T_1) + \epsilon x_{12}(T_0, T_1) + \epsilon^2 x_{13}(T_0, T_1)$. Similarly, y is taken to be $y_{11}(T_0, T_1) + \epsilon y_{12}(T_0, T_1) + \epsilon^2 y_{13}(T_0, T_1)$. So, that one T_1 . So, this is T_0, T_1 can be written.

So, here T_n . So, you can recall that a T_n is nothing, but ϵ to the power n t . So, T_n . So, that is why T_0 equal to t and T_1 equal to ϵt and T_2 equal to $\epsilon^2 t$. So, these are the different time scales similar to the time scale we are using in your watch that is for example, this hour hand, this minute hand and the second hand.

So, here we are taking three different time scale that is T_0 equal to t , T_1 equal to ϵt and T_2 equal to $\epsilon^2 t$ where ϵ is a book keeping parameter which is very very less than 1. So, as it is very very less than 1. So, these T_1 and T_2 r further less and less than these time scale t .

So, now, for the delay part you can note for the delay part we can use in a similar way. So, x_d is written as $x_{11} d T_0$ minus τ . So, this t is replaced by t minus τ . So, t is replaced by t minus τ is replaced by t minus τ for the delay differential equation or a delay part. So, delay time can be written by using t minus τ . So, this T_0 can be written as T_0 minus τ .

Similarly, T_1 can be written multiplying ϵ we can write this is T_0 minus τ . So, this ϵT_0 is nothing, but. So, this is equal to ϵT_0 minus $\epsilon \tau$. So, this ϵT_0 is nothing, but our T_1 time scale.

So, this is T_1 minus $\epsilon \tau$. So, that is why it is written T_0 minus τ , T_1 minus $\epsilon \tau$ similarly we can have these ϵ^2 for T_2 it will be $\epsilon^2 T_0$ minus τ . So, it will be $\epsilon^2 T_0$ is nothing but T_2 . So, it will be T_2 minus $\epsilon^2 \tau$. So, that way it can be written.

So, these x_d that is the delay. So, this is the displacement in the previous cut that is x_d can be written equal to $x_{11} d T_0$ minus τ , T_1 minus $\epsilon \tau$ and T_2 minus $\epsilon^2 \tau$. So, this is plus $\epsilon x_{12} d T_0$ minus τ T_1 minus $\epsilon \tau$ T_2 minus $\epsilon^2 \tau$. Similarly, so, we can go up to ϵ^2 . So, $\epsilon^2 x_{13} d$ into T_0 minus τ T_1 minus $\epsilon \tau$ then it will be T_2 minus $\epsilon^2 \tau$ similarly y_d can be written using these $y_{11} d$, $y_{12} d$ and $y_{13} d$.

So, writing these x y and the delay part of it in this form and using this d by d t equal to D 0 plus epsilon D 1 and d square by d t square equal to D 0 square plus 2 epsilon D 0 D 1. So, we can reduced our equation and then separating of the order of epsilon.

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PERTURBATION METHOD : METHOD OF MULTIPLE SCALES

Equating the coefficient of $\epsilon^0, \epsilon^1, \epsilon^2$ and ϵ^3 in the equations (3) and (4) equal to zero, it can be obtained

ϵ^0

$$D_0^2 x_{11} + \omega_{11}^2 x_{11} = 0$$

$$D_0^2 y_{11} + \omega_{21}^2 y_{11} = 0$$

ϵ^1

$$D_0^2 x_{12} + \omega_{11}^2 x_{12} + 2D_0 D_1 x_{11} + 2D_0 x_{11} \omega_{11} - k_{12} x_{11} + k_{13} y_{11} = 0$$

$$D_0^2 y_{12} + \omega_{21}^2 y_{12} + 2D_0 D_1 y_{11} + 2D_0 y_{11} \omega_{21} + k_{21} x_{11} - k_{22} y_{11} = 0$$

ϵ^2

$$D_0^2 x_{13} + \omega_{11}^2 x_{13} + D_1^2 x_{11} + 2D_0 D_1 x_{12} + 2D_0 x_{12} \omega_{11} + 2D_0 x_{11} \omega_{11} - \frac{1}{2} i \exp(-i \Omega T_0) k_{11} \phi_{11} y_{11} + \frac{1}{2} i \exp(i \Omega T_0) k_{11} \phi_{11} y_{11} + \frac{1}{2} i \exp(-i \Omega T_0) k_{12} \phi_{11} y_{12} - \frac{1}{2} i \exp(i \Omega T_0) k_{21} \phi_{11} x_{12} - k_{12} x_{11} + k_{13} x_{11} - k_{13} y_{11} + x_{11}^2 \omega_{11} + x_{11}^2 \omega_{21} = 0$$

$$D_0^2 y_{13} + \omega_{21}^2 y_{13} + D_1^2 y_{11} + 2D_0 D_1 y_{12} + 2D_0 y_{12} \omega_{21} + 2D_0 y_{11} \omega_{21} - \frac{1}{2} i \exp(-i \Omega T_0) k_{21} \phi_{11} x_{12} - \frac{1}{2} i \exp(i \Omega T_0) k_{21} \phi_{11} x_{12} - \frac{1}{2} i \exp(-i \Omega T_0) k_{22} \phi_{11} y_{12} + \frac{1}{2} i \exp(i \Omega T_0) k_{22} \phi_{11} y_{12} + k_{21} x_{11} - k_{22} x_{11} - k_{22} y_{11} + y_{11}^2 \omega_{21} + y_{11}^2 \omega_{11} = 0$$

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Handwritten notes:
 $D_0^2 x_{12} + \omega_{11}^2 x_{12} = f(x_{11}, y_{11})$
 $D_0^2 y_{12} + \omega_{21}^2 y_{12} = g(x_{11}, y_{11})$

So, we can different order epsilon that is epsilon to power 0, epsilon to the power 1 epsilon to the power 2. So, we have these equations. So, our equation reduces to D 0 square x 11 plus omega 11 square x 11 equal to 0, D 0 square y 11 plus omega 21 square y 11 equal to 0. So, actually while writing this equation you must take care so, that the right hand side forcing term will be of the order epsilon so, that in the first case that is epsilon to the power 0 the right hand side you will get 0.

So, this will happen. So, in case the system is weakly the forcing is weakly non-linear. So, if the forcing is strong then in that case in the right hand side you we have a forcing term. So,

which not will be equal to 0 and your equation of motion the solution of this equation will be different. It will be similar to that what we have studied in case of the hard excitation.

So, how this equation is similar to what we have studied in case of the weak excitation. So, that is why in the right side we have 0 0 term. So, our equation is reduced to this $D^2 x + \omega_1^2 x = 0$ $D^2 y + \omega_2^2 y = 0$.

So, already we know the solution of this thing where x_1 can be written by using. So, this is second order differential equation. So, by taking two constants it can be written similarly now for the after we get the solution for x_1 y_1 .

So, we can substitute it in the second equation and then getting the solution for these x_2 and y_2 we can substitute those things in the third equation that is of epsilon square to get the required expression. So, here you can recall that we have to find the secular terms and eliminating those secular terms will get the reduced equation.

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PERTURBATION METHOD : MMS

$$\begin{aligned}
 & \epsilon^3 \\
 & D_0^2 x_{14} + \omega_{11}^2 x_{14} + 2D_0 D_1 x_{11} + 2D_0 D_2 x_{11} + D_1^2 x_{11} + 2D_0 D_1 x_{12} + 2D_0 D_2 x_{12} + 2D_1 D_1 \mu \omega_{11} + 2D_1 D_2 \mu \omega_{11} \\
 & + 2D_0 D_1 \mu \omega_{11} - \frac{1}{2} \exp(-i\Omega T_0) k_{12} q_{22} x_{11} + \frac{1}{2} \exp(i\Omega T_0) k_{12} q_{22} x_{11} + \frac{1}{2} \exp(-i\Omega T_0) k_{12} q_{22} x_{14} \\
 & - \frac{1}{2} \exp(i\Omega T_0) k_{12} q_{22} x_{14} - \frac{1}{2} \exp(-i\Omega T_0) k_{11} q_{11} y_{12} + \frac{1}{2} \exp(i\Omega T_0) k_{11} q_{11} y_{12} + \frac{1}{2} \exp(-i\Omega T_0) k_{11} q_{11} y_{14} \\
 & - \frac{1}{2} \exp(i\Omega T_0) k_{11} q_{11} y_{14} - k_{12} x_{12} + k_{12} x_{14} - k_{11} y_{13} + k_{11} y_{14} + 2x_{11} \mu \omega_{11} + 3x_{11}^2 \mu \omega_{11} = 0 \\
 \\
 & D_0^2 y_{14} + \omega_{11}^2 y_{14} + 2D_0 D_1 y_{11} + 2D_0 D_2 y_{11} + D_1^2 y_{11} + 2D_0 D_1 y_{12} + 2D_0 D_2 y_{12} + \frac{1}{2} \exp(-i\Omega T_0) k_{22} q_{22} x_{11} \\
 & - \frac{1}{2} \exp(i\Omega T_0) k_{22} q_{22} x_{11} - \frac{1}{2} \exp(-i\Omega T_0) k_{22} q_{22} x_{14} + \frac{1}{2} \exp(i\Omega T_0) k_{22} q_{22} x_{14} + \frac{1}{2} \exp(-i\Omega T_0) k_{21} q_{11} y_{12} \\
 & - \frac{1}{2} \exp(i\Omega T_0) k_{21} q_{11} y_{12} - \frac{1}{2} \exp(-i\Omega T_0) k_{21} q_{11} y_{14} + \frac{1}{2} \exp(i\Omega T_0) k_{21} q_{11} y_{14} + k_{22} x_{12} - k_{22} x_{14} + k_{21} y_{13} \\
 & - k_{21} y_{14} + 2y_{11} \mu \omega_{11} + 3y_{11}^2 \mu \omega_{11} + 2D_1 y_{11} \mu \omega_{11} + 2D_1 y_{12} \mu \omega_{11} + 2D_1 y_{14} \mu \omega_{11} = 0
 \end{aligned}$$

□ The solution of the equations (19) and (20) can be written as

$$\begin{aligned}
 x_{11}(T_0, T_1, T_2) &= A_1(T_1, T_2) \exp(i\omega_{11} T_0) + \bar{A}_1(T_1, T_2) \exp(-i\omega_{11} T_0) \\
 x_{14}(T_0 - \tau, T_1) &= A_1(T_1, T_2) \exp(i\omega_{11}(T_0 - \tau)) + \bar{A}_1 \exp(-i\omega_{11}(T_0 - \tau)) \\
 y_{11} &= A_2(T_1, T_2) \exp(i\omega_{21} T_0) + \bar{A}_2(T_1, T_2) \exp(-i\omega_{21} T_0) \\
 y_{14}(T_0 - \tau, T_1) &= A_{2d}(T_1) \exp(i\omega_{21}(T_0 - \tau)) + \bar{A}_{2d}(T_1) \exp(-i\omega_{21}(T_0 - \tau))
 \end{aligned}$$

So, you just see as we are taking up to epsilon square. So, you can take up to epsilon q. So, these equation is written up to epsilon cube now the solution of this one solution of D 0 square x 11 plus omega 11 square x 1 equal to 0 and this D 0 square y 11 plus omega to 1 square y 11 equal to 0.

So, in this form. So, A it is equal to A 1. So, as this A 1 and A 1 bar will be constant. So, they will not be function of T 0. So, that is why they may be function of T 1 and T 2 that s why it is written A 1 T 1 T 2 it is not written T 0. So, you can note that it is not a function of T 0. So, x 11, T 0, T 1, T 2 can be written as A 1 T 1, T 2 sine e to power i omega 11 T 0 plus A 1 bar T 1 T 2 e to the power minus i omega 11 T 0.

So, here. So, you can note that. So, these i this A 1 bar is the complex conjugate of A1 A 1 bar A 1 and A1 bar are complex number and these A 1 bar is the complex conjugate of A 1

similarly. So, we can write down the second equation y_{11} , y_{11} equal to $A_2 T_1 T_2 e$ to the power $i \omega_2 T_0$ plus $A_2 \bar{d} T_1 T_2 e$ to the power $-i \omega_2 T_0$, Similarly, $y_{11 d}$ can be written $A_2 d$.

So, A_2 will be replaced by $A_2 d T_1 T_2$ can be written A to the power $i \omega_2 T_0$. So, T_0 will be replaced by $T_0 - \tau$ similarly $A_2 \bar{d}$ equal to $A_2 \bar{d} T_1 e$ to the power $-i \omega_2 T_0 - \tau$.

Similarly, this $x_{11 d}$ can be written equal to A_1 . So, it will be $A_1 d$. So, here we are using this $A_1 d T_1, T_2 e$ to the power $i \omega_1 t$. So, these T_0 will be replaced by $T_0 - \tau$ plus similarly $A_1 \bar{d} e$ to the power $-i \omega_1 T_0 - \tau$. So, it will be $T_0 - \tau$. So, this way it can be written.

So, now, writing these $x_{11}, x_{1 d}, y_{11}, y_{1 d}$. So, we can substitute these equation in these equation. So, you just see only we will take these 2 terms in the left hand side. So, these equation can be written D_0^2 . So, if you recall. So, $D_0^2 x_{12} + \omega_1^2 x_{12}$ plus $\omega_1^2 x_{12}$. So, other part will shift it to the right hand side.

Similarly, the second equation also we will write this is equal to $D_0^2 y_{12} + \omega_2^2 y_{12}$. So, other part will shift to the right hand side. So, you can note that this other part is a function of both. So, it is a function of both x_{11}, y_{11} and $x_{1 d}, y_{1 d}$. So, these are the function of both x and the delay part the system. So, that way we can write down these equations. So, now, substitute these solution they are.

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PERTURBATION METHOD : MMS

Eliminating the secular terms, the solution of (21) and (22) can be written as
 Now substituting the above solution in (19) and (20). For the condition of $\omega_{21} \approx \omega_1$ and $\Omega - \omega_{11} \approx \omega_{11}$

$$\left. \begin{aligned} 2i\omega_{11}D_1A_1 + 2i\mu_1\omega_{21}^2A_1 - k_{21}A_2 \exp(i\sigma_1T_1) + k_{11}A_{2d} \exp(-i\tau\omega_{21} + i\sigma_1T_1) &= 0 \\ 2i\omega_{21}D_1A_2 + 2i\mu_2\omega_{21}^2A_2 + k_{21}A_1 - k_{21}A_{2d} \exp(-i\omega_{21}\tau) & \end{aligned} \right\}$$

As the equations do have any particular solution, the non-transient solution of these equation will be

$$x_{12} = x_{2d} = y_{12} = y_{2d} = 0$$

$$\left. \begin{aligned} 2i\omega_{11}D_2A_1 + D_1^2A_1 - k_{21}A_1 + k_{21}A_{2d} \exp(-i\omega_{11}\tau) + \frac{1}{2}ik_{11}q_{11}A_1 \exp(i\sigma_1T_1 - i\sigma_2T_1) \\ - \frac{1}{2}ik_{11}q_{11}A_{2d} \exp(i\omega_{21}\tau + i\sigma_1T_1 - i\sigma_2T_1) + 3\sigma_{13}A_1^2\hat{A}_1 + 2\mu_1\omega_{11}D_1A_1 &= 0 \\ 2i\omega_{21}D_2A_2 + D_2^2A_2 + k_{22}A_1 \exp(-i\sigma_2T_2) - k_{22}A_{2d} \exp(-i\omega_{21}\tau - i\sigma_2T_2) - \frac{1}{2}ik_{21}q_{11}A_1 \exp(i\sigma_1T_1 - 2i\sigma_2T_2) \\ + \frac{1}{2}ik_{21}q_{11}A_{2d} \exp(i\omega_{21}\tau + i\sigma_1T_1 - 2i\sigma_2T_2) + 3\sigma_{23}A_2^2\hat{A}_2 + 2\mu_2\omega_{21}D_1A_2 &= 0 \end{aligned} \right\}$$

Substituting the solution of the solution of x_{12}, x_{2d}, y_{12} and y_{2d} and the secular terms will be obtained following equations

$$\left. \begin{aligned} 2i\omega_{11}D_3A_1 + 2D_3D_1A_1 + \frac{1}{2}ik_{12}q_{12}A_1 \exp(i\sigma_1T_1) - \frac{1}{2}ik_{12}q_{12}A_{2d} \exp(i\omega_{11}\tau + i\sigma_1T_1) + 2\mu_1\omega_{11}D_1A_1 &= 0 \\ 2i\omega_{21}D_3A_2 + 2D_3D_2A_2 + 2\mu_2\omega_{21}D_2A_2 - \frac{1}{2}ik_{22}q_{22}A_1 \exp(-i\sigma_2T_1 + i\sigma_1T_1) + \frac{1}{2}ik_{22}q_{22}A_{2d} \exp(-i\sigma_2T_1 + i\sigma_1T_1 + i\omega_{11}\tau) &= 0 \end{aligned} \right\}$$

So, we can get. So, now, we can eliminate the secular terms. So, by eliminating secular term. So, our equation will reduce to these form and so, as you know that as the equation do not have any particular solution the non-transient solution of these equations will be $x_{12} = x_{2d} = y_{12} = y_{2d} = 0$.

So, by putting these for steady state actually $x_{12} = x_{2d} = y_{12} = y_{2d} = 0$. So, we can have these equation. So, now, we have these equation. So, substituting the solution of the $x_{13} = x_{2d} = y_{13} = y_{2d}$ in the secular term. So, we can obtain a set of equations like this and here also now, what we can write.

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PERTURBATION METHOD : MMS

□ The reduced temporal equation of tool in cutting and feed force direction can be obtained after substituting the expression of secular terms in below expressions and separating the real and imaginary parts as follows

$$\frac{dA_j}{dt} = \frac{d}{dt} \left(\frac{1}{2} a_j e^{i\theta_j} \right) = \epsilon D_1 A_j + \epsilon^2 D_2 A_j + \epsilon^3 D_3 A_j \dots \dots \dots \quad (61)$$

Where, $j=1$ for equation in cutting force direction and $j=2$ for equation in feed direction

□ Letting $A_j = A_{jR} + i A_{jI} = \frac{1}{2} a_j e^{i\theta_j}$ and $A_2 = A_{2R} + i A_{2I}$

$$a_1^R = c_{11} a_1 + c_{12} a_2 \cos[\gamma_2] + c_{122} a_2^2 \cos[\gamma_2] + c_{13} a_2 \sin[\gamma_2] + c_{133} a_2^2 \sin[\gamma_2] + c_{14} a_1 \cos[2\gamma_1] + c_{144} a_1 \sin[2\gamma_1] + c_{144} a_2 \cos[2\gamma_1 - \gamma_2] + c_{144} a_2 \sin[2\gamma_1 - \gamma_2] \quad (62)$$

$$a_1^I = \frac{\epsilon \sigma_1 a_1}{2} + c_{21} a_1 + c_{211} a_1^2 + c_{22} a_2 \cos[\gamma_2] + c_{222} a_2^2 \cos[\gamma_2] + c_{23} a_2 \sin[\gamma_2] + c_{233} a_2^2 \sin[\gamma_2] + c_{24} a_2 \cos[2\gamma_1] + c_{244} a_2 \sin[2\gamma_1] + c_{244} a_2 \cos[2\gamma_1 - \gamma_2] + c_{244} a_2 \sin[2\gamma_1 - \gamma_2] \quad (63)$$

$$a_2^R = c_{31} a_1 + c_{311} a_1^2 + c_{32} a_2 \cos[\gamma_2] + c_{322} a_2^2 \cos[\gamma_2] + c_{33} a_2 \sin[\gamma_2] + c_{333} a_2^2 \sin[\gamma_2] + c_{34} a_2 \cos[2\gamma_1 - 2\gamma_2] + c_{344} a_2 \sin[2\gamma_1 - 2\gamma_2] + c_{344} a_2 \cos[2\gamma_1 - \gamma_2] + c_{344} a_2 \sin[2\gamma_1 - \gamma_2] \quad (64)$$

$$a_2^I = \frac{\epsilon \sigma_2 a_2}{2} - \epsilon \sigma_2 a_2 + a_2 c_{41} + c_{411} a_1^2 + c_{42} a_1 \cos[\gamma_2] + c_{43} a_1 \sin[\gamma_2] + c_{44} a_2 \cos[2\gamma_1 - 2\gamma_2] + c_{444} a_2 \sin[2\gamma_1 - 2\gamma_2] + c_{444} a_2 \cos[2\gamma_1 - \gamma_2] + c_{444} a_2 \sin[2\gamma_1 - \gamma_2] \quad (65)$$

Where $-\theta_1 + \theta_2 + \tau_1 \sigma_2 = \gamma_2$ and $-2\theta_1 + \tau_1 \sigma_1 = 2\gamma_1$
 $-\theta_1 - \theta_2 + \tau_1 \sigma_1 - \tau_1 \sigma_2 = -\gamma_2 + 2\gamma_1$ and $-2\theta_2 + \tau_1 \sigma_1 - 2\tau_1 \sigma_2 = 2\gamma_1 - 2\gamma_2$

So, we can write these taking these A equal to A equal to half a e to the power i theta form. So, we can write A equal to as is a complex number. So, we can write A equal to half a e to the power i theta and this d A i by d t that is d A 1 by d t for example, will be equal to d by d t half A 1 e to the power i theta 1.

So, which will be equal to epsilon D 1 A1 plus epsilon square D 2 A 1 plus epsilon cube D 2 A 1. So, in that way if we can write. So, we can have these four equations. So, these four equations that is A 1 dash A 1 gamma 1 dash A 2 dash and A 2 theta 2 dash. So, this can be written in terms of gamma 2 dash because we are reducing this theta to gamma form to make the equation in non-autonomous form.

So, theta 1 can be written theta 1 minus theta 1 plus theta 2 plus T 1 sigma 2. So, here we are using the detuning parameter sigma 1 and sigma 2. So, minus theta 1 plus theta 2 plus T 1 sigma 2 equal to gamma 2 and minus 2 the theta 1 plus T 1 sigma one equal to 2 gamma 1.

And it can be written this minus theta 1 minus theta 2 plus T 1 sigma one minus T 1 sigma 2 equal to minus gamma 2 plus 2 gamma 1 and minus 2 theta 1 plus T 1 sigma 1 minus 2 T 1 sigma 2 equal to 2 gamma 1 minus gamma 2. So, by writing in this form.

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PERTURBATION METHOD : MMS

With this below assumption,

$$p_1 = a_1 \cos \gamma_1; q_1 = a_1 \sin \gamma_1; p_2 = a_2 \cos(\gamma_1 - \gamma_2); q_2 = a_2 \sin(\gamma_1 - \gamma_2)$$

The reduced temporal equation of tool in cutting and feed force direction will be obtained as

$$\begin{aligned} \dot{p}_1' &= -c_{11}(p_1^2 + q_1^2)p_1 - c_{12}q_1 + c_{14}q_1 - c_{22}q_2 + c_{34}q_2 - c_{22}(p_1^2 + q_1^2)q_2 + c_{11}p_1 - c_{23}p_1 \\ &\quad - c_{23}p_2 - c_{27}p_2 - c_{23}p_1(p_1^2 + q_1^2) - \frac{1}{2}\epsilon_1 p_1 \\ \dot{q}_1' &= c_{21}(p_1^2 + q_1^2)p_1 + c_{21}p_1 + c_{34}p_1 + c_{22}p_2 + c_{34}p_2 + c_{22}(p_1^2 + q_1^2)p_2 + c_{11}q_1 + c_{23}q_1 \\ &\quad - c_{23}q_2 + c_{27}q_2 - c_{23}q_1(p_1^2 + q_1^2) + \frac{1}{2}\epsilon_1 q_1 \\ \dot{p}_2' &= -c_{41}(p_2^2 + q_2^2)p_2 - c_{41}q_2 + c_{44}q_2 - c_{42}q_1 + c_{46}q_1 + c_{211}(p_1^2 + q_1^2)p_2 + c_{21}p_2 - c_{43}p_2 \\ &\quad + c_{43}p_1 - c_{47}p_1 - \frac{1}{2}\epsilon_2 p_2 + \epsilon_2 q_2 \\ \dot{q}_2' &= c_{41}(p_2^2 + q_2^2)p_2 + c_{41}p_2 + c_{44}p_2 + c_{42}p_1 + c_{46}p_1 + c_{211}(p_1^2 + q_1^2)q_2 + c_{21}q_2 + c_{43}q_2 \end{aligned}$$

The above reduced equations can be used for the steady state response in cutting force direction and feed direction

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So, our equation is reduced and it can be observed that due to the present of the term a 1 gamma 1 dash, the trivial state conditions will not contain the perturbation. So, it can be. So, we cannot obtain the stability from the earlier equation. So, now, we have to make the

transformation in such way that. So, we will get distinct equations which can be perturb to get these turbulent equation.

So, here we can assume this p_1 equal to $a_1 \cos \gamma_1$ q_1 equal to $a_1 \sin \gamma_1$, p_2 equal to $a_2 \cos \gamma_1 - \gamma_2$ and q_2 equal to $a_2 \cos \gamma_1 - \gamma_2$ q_2 equal to $\sin a$. So, it will be $a \sin$. So, it will be $a \sin \gamma_1 - \gamma_2$.

So, then substituting these equation. So, we can obtain this p_1 dash q_1 dash p_2 dash and q_2 dash equations. So, which we will perturb to get these δp_1 , δp_1 dot δq_1 dot, δp_2 dot and δq_2 dot. So, this thing can be written equal to this Jacobean matrix into δp_1 , δq_1 , δp_2 , δq_2 .

So, in this form can be written and so, this is the Jacobean matrix. So, now, finding these Eigen value of the Jacobean matrix. So, we can find whether the system is stable or not. So, as we are interested for the trivial state stability region or to find the instability region. So, we can find for the system parameter for which the system becomes stable or unstable.

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Physical and Material Properties of Tool

Parameter Description	Values
Tool Length (mm)	150
Tool cross-section (mm ²)	28 x 28
Tool 1st natural frequency (Hz)	988
Tool 2nd natural frequency (Hz)	988
Tool material	Steel
Density (kg/mm ³)	7.85E-06
Tool mass (kg)	0.92
Workpiece radius (mm)	12
Chip width (mm)	0.3
Feed rate (mm/s)	0.2
Workpiece rotational speed (rad/s)	3
Force coefficient in cutting force direction (N/mm ²)	1500
Force coefficient in feed force direction (N/mm ²)	450

Please note: Tool is modeled in Ansys and static stiffness and first fundamental frequency is computed from there.

Using the scaling factor (τ_s) = 0.001 and book-keeping parameter (ϵ) = 0.01, the following **non-dimensional System Parameters** can be obtained.

$\alpha_{13} = 8.33; \alpha_{23} = 8.33; \alpha_{12} = -2.00; \alpha_{22} = 8.00; K_1 = 0.005; k_{11} = 6.55; k_{12} = 3.64; k_{21} = 1.96; k_{22} = 1.09; \omega_{11} = 1; \omega_{12} = 1; \mu_1 = 1; \mu_2 = 1$

The coefficients are mainly function of sine or cosine of $\omega_{11}\tau$ and $\omega_{21}\tau$ where $\tau = \sqrt{\omega_{11}}$ is the angle in radian using the relation $\tau = \sqrt{\omega_{11}}t$, ω_{11} and ω_{21} is non-dimensional and has unit value.

The values of sine or cosine of $\omega_{11}\tau$ and $\omega_{21}\tau$ will be repeated in the range of 0 to 2π based on sine or cosine in case of the steady-state condition.

$m_1 = 0.92 \text{ kg}; k_{11} = 3.6 \times 10^9 \text{ N/mm}; k_{12} = 3.6 \times 10^8 \text{ N/mm}; k_{21} = 30 \text{ N/mm}^2; k_{22} = 30 \text{ N/mm}^2; k_3 = 1500 \text{ N/mm}^2; k_f = 450 \text{ N/mm}^2; R = 12 \text{ mm}; w_c = 0.3 \text{ mm}; v_0 = 0.2 \text{ mm/s}; q = 0.75; \Omega_w = 3 \text{ rad/s}; \epsilon = 0.1$

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So, if the real part of the Eigen values of the Jacobean matrix is positive then the system is unstable and if the real part of the Eigen values of the Eigen values are negative, then the system is stable. So, taking these conditions. So, we can study the stability of the system.

So, here these parameters have been taken for example, you have taken the tool length to be 150 mm. So, tool cross section 20. So, it is square cross section we have taken 28 into 28 mm square. So, tool first natural frequency it can be observe that it is around 988 hertz, tool second natural frequency. So, as we are taking a square cross section. So, it is same that is 988.

So, tool material is taken steel. So, density is taken to be 7.85 2 E minus 6 kg per mm cube. So, this tool mass we have taken only 0.92 kg and work piece radius to be 12 mm. So, chip width that is 0.3, feed rate is taken to be 0.2, work piece rotational speed is 3 radian per

second. So, force coefficient in cutting forces direction Newton per mm square this taken to be 1500 and the force coefficient in speed force direction. So, it is taken to be 450.

So, taking these parameters. So, the other parameters are obtained the coefficient are mainly function of sin and cosine of $\omega_{11} \tau$ and $\omega_{21} \tau$ where τ equal to $\omega_{11} \tau$ bar τ is the angle in radian using the relation τ bar equal to root over $\omega_{11} \tau$ bar t . So, $\omega_{11} \tau$ bar and $\omega_{21} \tau$ bar are non-dimensional and has unity value. So, we have taken it is in such way that. So, they have the unity value. The values of sin and cosine of $\omega_{11} \tau$ and $\omega_{11} \tau$ and $\omega_{21} \tau$ will be repeated in the range of 0 to 2 pi.

So, based on the sin and cosine in case of the steady state condition. So, this value taking these value. So, the τ term have been determined and they have been taken in the analysis. So, here you just see the other parameter what we have taken in this analysis m_1 equal to 0.92. So, here it is written the mass is taken to be 0.92.

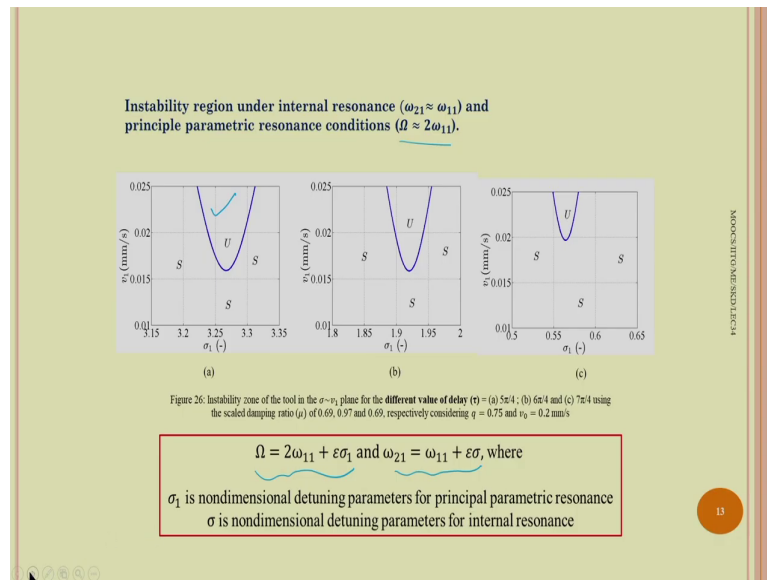
So, that is why m_1 is taken to be 0.92, k_{x1} equal to 3.6 into 10 to the power 4 newton per m, k_{y3} 13.6 into 10 to the power 4 newton per mm, k_{x3} equal to 30 newton per mm cube, k_{y3} 30 newton per mm cube as we are taking these cross section square cross section. So, the stiffness are taken similar in both x and y direction. So, k_c is taken to be 1500 Newton per mm square, k_f equal to 450 newton per mm square, R equal to 12 mm, w_c equal to 0.3 mm and v_0 you just see this is the feed taken to be.

So, feed v_0 equal to 0.2, v_0 equal to 0.2 millimeter per second. So, q that coefficient that power term you have seen q . So, q is taken to be 0.75 and this is taken this rotational speed of the spindle is taken to be 3 radian per second and r is taken to be. So, this the scaling parameter r . So, that is taken to be 0.1. So, taking all these parameters. So, you can absorb that the coefficients are coming to be in the. So, coefficients are coming in the proper order.

So, using scaling factor r_1 equal to 0.001 and book keeping parameter epsilon equal to 0.01, the following non dimensional system parameters are obtained. So, that α_{13} equal to 8.33 α_{23} 8.23 q_{12} minus 2 q_{22} equal to 8 and k_1 equal to 0.005, k_{11} equal to 6.55, k_{12} equal to 3.64, k_{21} 1.96, k_{22} 1.09 ω_{11} you just note. So, it is non-dimensionalize in

such way that this ω_{11} equal to 1 and ω_{12} also equal to 1 and μ_1 equal to 1 and μ_2 also it is coming to be 1.

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So, taking these parameters. So, and from the Jacobean matrix. So, one can find the stability just by finding the Eigen values and checking the Eigen values. So, for different value of b_1 and σ . So, that is the feed rate. So, and σ one is the detuning parameter. So, you can get this instability region.

So, the region outside these curve its stable inside this is unstable and outside it is stable. So, the operator has to take the feed rate and also the rotational speed in such way that it should in these stable region to have a vibration free or there will be no vibration of the tool.

So, if it is within these range in this instability region then the tool will vibrate. So, the obtained region. So, will be very much useful to find the operating conditions for which there will be no chatter in the turning operation. So, here the instability region under internal resonance condition here we have taken. So, you just see we have taken this ω_{21} equal to ω_{11} .

So, that is why. So, we have this as this is 1 is to 1. So, we are considering this 1 is to 1 internal resonance condition in this case and along with these thing. So, we are considering the principle parametric resonance conditions also, that is capital ω equal to $2\omega_{11}$ plus $\epsilon\sigma_1$ and ω_{21} equal to ω_{11} plus $\epsilon\sigma_1$.

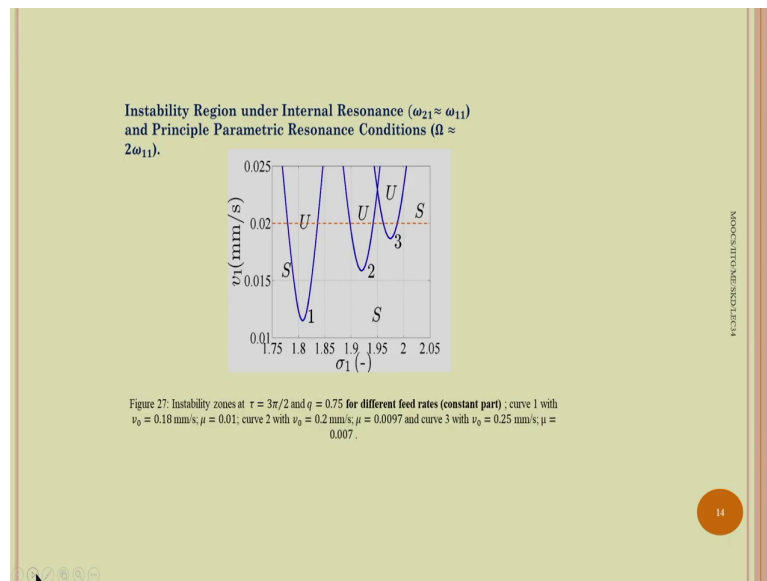
So, this is scattering the internal resonance condition and this is scattering the this is scattering the condition for though we can tell this is external resonance condition that is principle parameter resonance condition, but this is due to the rotational speed the variation in the rotational speed in a periodic wave. So, due to that assumptions we are getting this thing.

So, σ_1 is non dimensional detuning parameter for principle parameter resonance and σ_2 is non dimensional detuning parameter for internal resonance condition. So, you can observe that. So, in stability zone for tool in σ_1 v 1 plane for the different value of delay. So, three different values of delay has been taken.

So, in the first case the delay is taken to be $5\pi/4$, in the second case it is $6\pi/4$ and in this third case it is $7\pi/4$ for μ equal to 0.69, 0.97 and 0.69 respectively considering q equal to 0.75 and v_0 equal to 0.2 millimeter per second.

So, you can observe how the instability region varies. So, in this case in the last case the instability region when the delay is taken to be $7\pi/4$. So, it is to be. So, you can have a very small instability region and for the most to cases most of the cases the system remain in stable region. So, taking different delay conditions.

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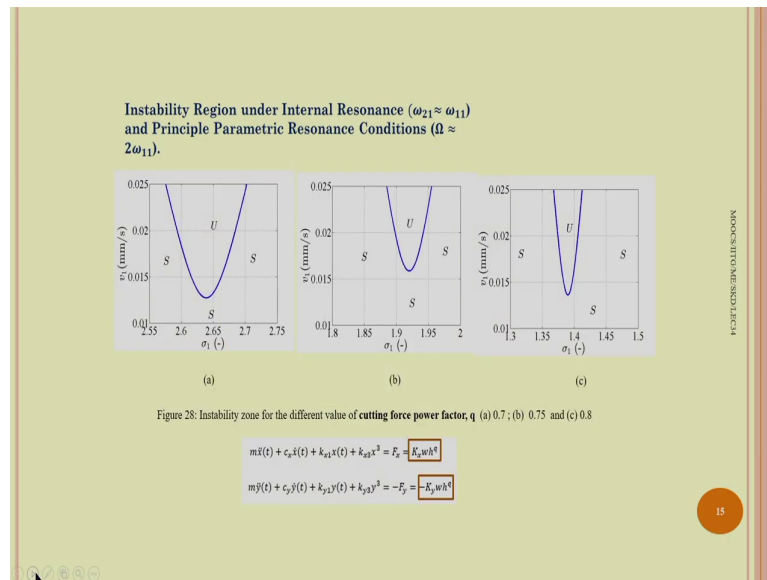


So, one can actually obtain this instability region and the operator will get the sense to operate it at a system parameter where the system will not have any vibration or there will be no chatter in the system. So, the instability region and under these internal resonance condition for principle parameter resonance it is studied here.

So, here tau is taken to be $3\pi/2$ and q equal to 0.75 for different feed rates. So, 3 different speed rates have been taken. So, first curve is for feed rate v_0 equal to 0.18, second mu or first case v_0 equal to 0.18 mu equal to 0.01 for the second curve. So, you can see with v_0 equal to point 0.2 millimeter per second mu equal to 0.0097 and for the third curve you can see v_0 equal to point 0.25 millimeter per second and mu equal to 0.007.

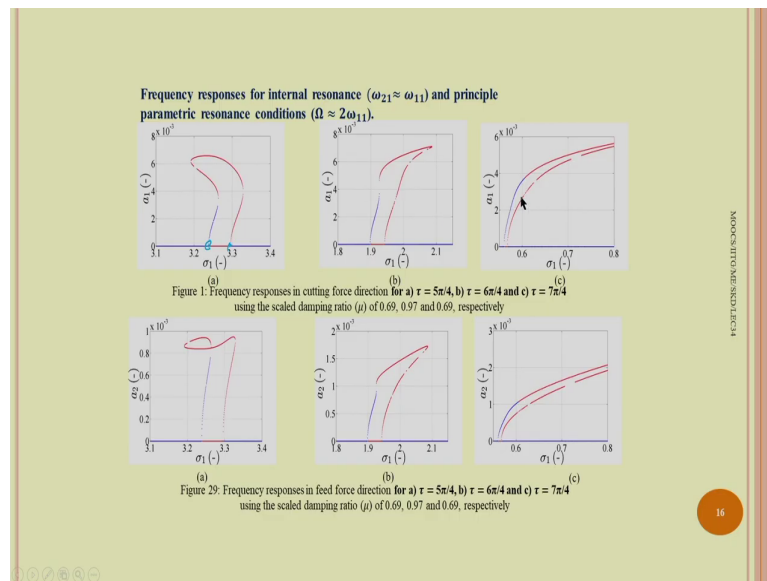
So, here also you can see. So, in this third case we have the we have more stable zone and the operator can operate it freely in this or for a wide range of system parameter.

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Similarly, so, we can vary different parameters and study the instability region.

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So, in addition to the parametric instability region. So, one can plot the frequency response plot also. So, in these unstable range. So, if someone operate this thing in unstable range for example, 1 to have the texture in the turning operation. So, in that case. So, he can operate it in the unstable region. So, when it is operated in this unstable region. So, it is one should know the frequency response plot and one can find so, how much response or how much vibration will be there in the system or in the tool.

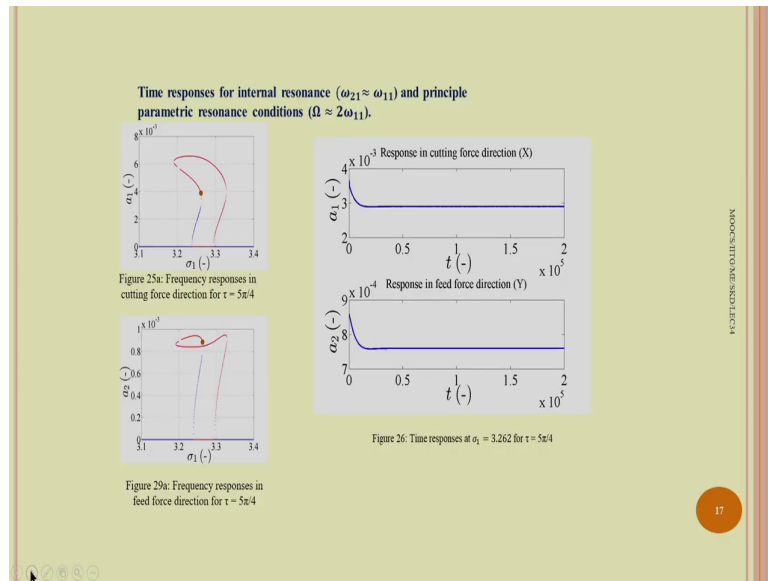
So, if one can plot the frequency response plot so, one can see that these point correspond to these point correspond to the super critical pitch for bifurcation point and these correspond to the sub critical pitch for bifurcation point. So, here. So, in the super critical pitch for bifurcation point. So, after the supercritical pitch for bifurcation point with increase in the sigma.

So, that is with increase in the speed of the speed rotational speed. So, it will vibrate and these response can be obtained. So, from this a 1 one can see. So, what will be the response similarly for a 2 a 1. So, for a different value of tau it has been plotted. So, in the first case tau equal to $5\pi/4$, second case it is $6\pi/4$ and third case it is $7\pi/4$.

Similarly, for frequency response can be plotted for different feed rate. So, this is the first case it is plotted for a different tau value that is delayed and in the second case different feed rate for tau equal to $5\pi/4$, $6\pi/4$ and tau equal to $7\pi/4$. So, this is in the cutting force direction and the second set of figures are in the feed force direction that is a 2. So, this is a 1 versus the sigma 1, this is a 2 verses versus sigma 2.

So, you can see. So, in the first case it is 8 into 10 to the power minus 3, 8 into and in the second case in y direction that is in the feed direction. So, it is 1 into 10 to the power minus 3. So, in feed direction. So, it is less and in the cutting force direction the response is found to be more. So, this way one can study different frequency and force response plot for a different system parameters.

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Time responses for internal resonance ($\omega_{21} \approx \omega_{11}$) and principle parametric resonance conditions ($\Omega \approx 2\omega_{11}$).

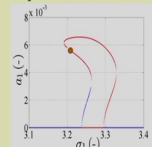


Figure 25a: Frequency responses in cutting force direction for $\tau = 5\pi/4$

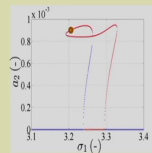


Figure 29b: Frequency responses in feed force direction for $\tau = 5\pi/4$

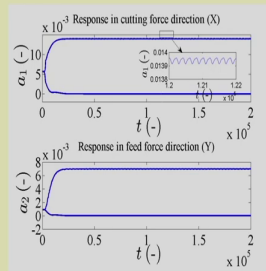
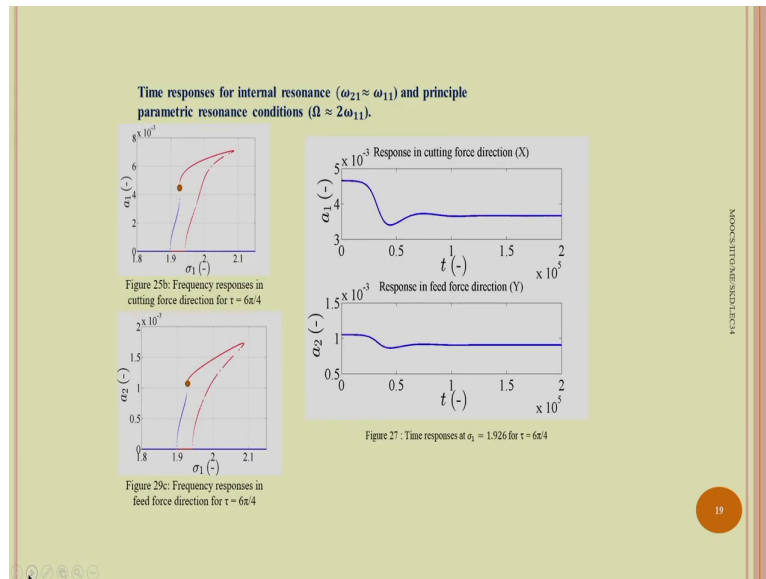


Figure 26: Time responses at $\sigma_1 = 3.207$ for $\tau = 5\pi/4$

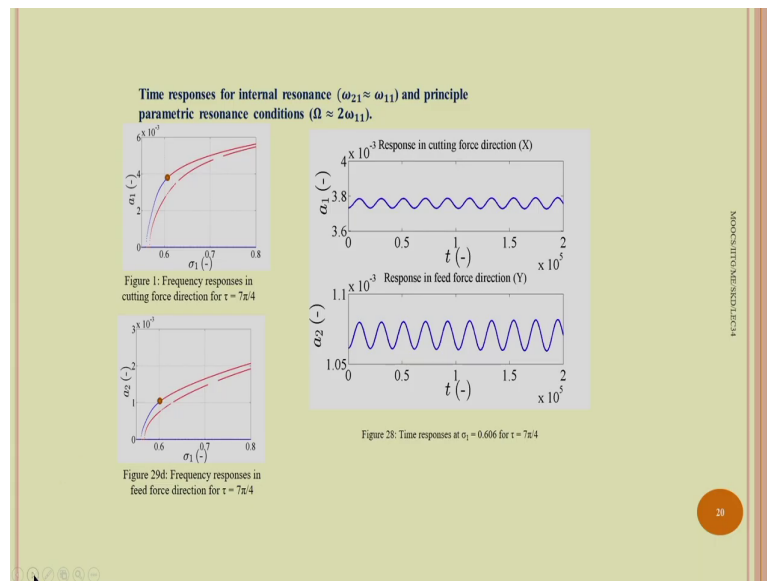
So, here the time response is also plotted to show the response in the cutting force direction and in the feed direction. So, two points have been taken. So, in this unstable zone. So, while taking in this unstable zone two points. So, you can see there is a there is a vibration or so, due to these vibration there will be chatter mark on the workpiece.

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So, for different point similarly it has been obtained.

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So, you can see this undulation is more in this case. So, time response for sigma 1 equal to 0.606 and tau equal to 7 phi by 4. So, one can study different marking or one can study different response of the turning operation in this way. So, in today class. So, we have discussed regarding the vibration of the tool. So, one can similarly study the vibration of the workpiece also.

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EQUATION OF MOTION OF WORKPIECE AND TOOL

➤ The Equations of motion of workpiece and tool as 2-DOF system with regenerative effect are as follows respectively

$$m_1 \ddot{x}_1 + c_1 \dot{x}_1 + \sum_{j=1}^3 k_{1j} x_1^j$$

$$m_2 \ddot{x}_2 + c_2 \dot{x}_2 + \sum_{j=1}^3 k_{2j} x_2^j = -k_c w_c h(t)$$

➤ Varied chip thickness $h(t)$ around the nominal chip thickness

$$h(t) = h_0 + x_1(t - \tau) - x_1(t) - x_2(t - \tau) + x_2(t).$$

➤ Introducing the external force on workpiece the equation of motion becomes

$$m_1 \ddot{x}_1 + c_1 \dot{x}_1 + \sum_{j=1}^3 k_{1j} x_1^j = k_c w_c h(t) + F \cos(\Omega t) \quad (73)$$

$$m_2 \ddot{x}_2 + c_2 \dot{x}_2 + \sum_{j=1}^3 k_{2j} x_2^j = -k_c w_c h(t) \quad (74)$$

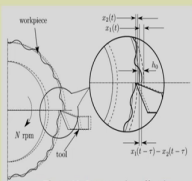


Figure 33: Regenerative effect in Turning operation

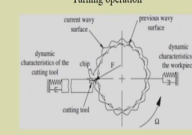
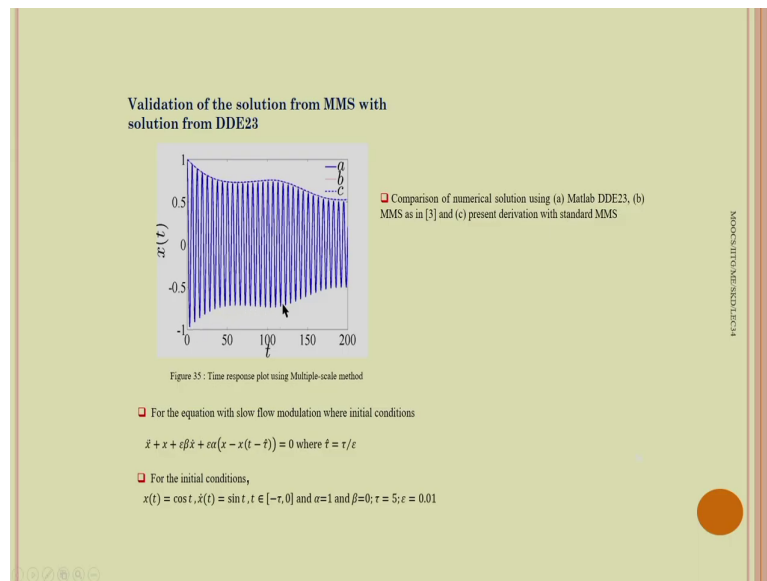


Figure 34: Turning operation modeled as 2-DOF system [2]

So, here. So, this part is given as an assignment to you to study the equation of motion of both work piece and tool when they are flexible. So, in that case this equation can be written in these form and following in the previous way. So, you can write down this governing equation.

So, after write down writing the governing equation of motion so, you can reduce it to this delay differential equations and then in by solving these solving using method of multiple scale. So, you can find the response of the system.

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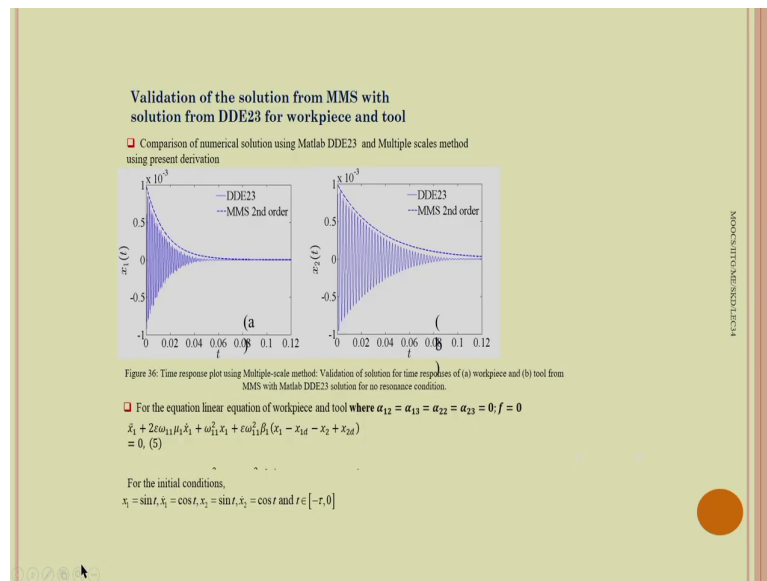


So, this part I am not going to discuss in detail but you can see that using this method of multiple scale on this DDE 23. So, you can verify that the results are coming to be similar or same.

So, the advantage of using this method of multiple scale is that in addition to getting the time response accurate time response you can get the frequency and force response plot of the system and the analysis can be extended to know the effects of different system parameter on the turning operation.

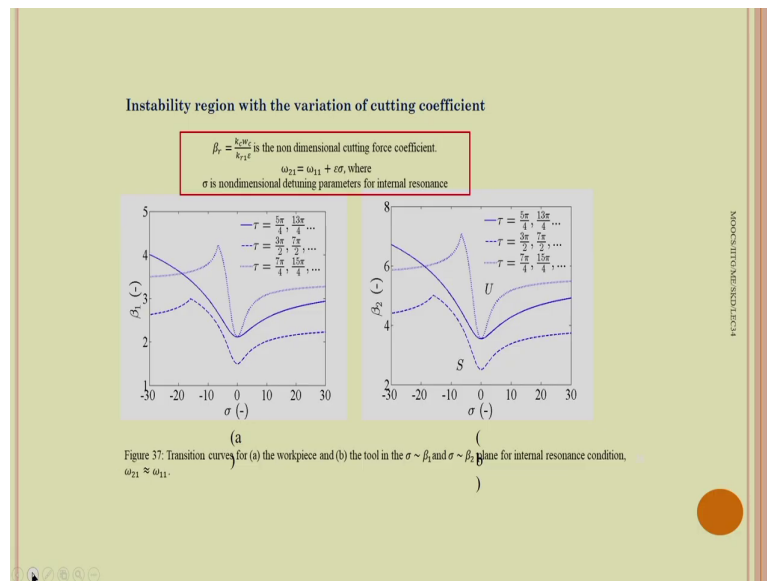
So, initially we have plotted this instability region. So, from the instability region we know the zone at which the system will not vibrate or the tool will not vibrate. So, in this case there will be no vibration of the tool and the workpiece.

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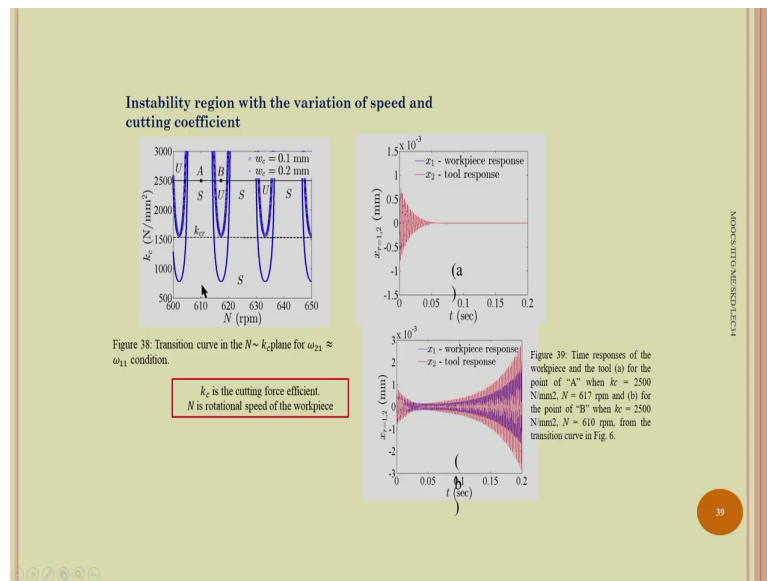
So, knowing those regions actually we can operate the system in that range. So, as to get the response of the system here more comparisons have been given between DDE to 23 and method of multiple scale second order method of multiple scale and these are the instability regions plotted.

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So, when both the tool and workpiece are flexible. So, this is the instability region. So, previous case when only the tool was flexible. So, you have seen we had a wide range of stability region, but in this case. So, the stability region is reduced because both the tools and the work piece are vibrating.

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So, these are the instability region you can plot and you can see between these the response will in the unstable range the response will grow are the response is growing. So, there will be mark on the tool, a mark will be there will be mark in the workpiece.

So, there will be chatter mark in the workpiece. So, we have to avoid those to avoid those chatter mark. So, the system has to be operated in this stable range. So, this way you can study different manufacturing system. So, this is the turning operations I have told similarly you can extend this work for the milling operation also you can derive this equation for the forming operation like this rolling operation.

So, in many manufacturing system as well the process is non-linear. So, you can extend your analysis your linear analysis to the field of non-linear analysis and you can find the operating

system parameter for which the systems will be stable and you can operate the system in the stable zone two half chatter free surfaces in the workpiece. So, thank you very much.