

WAVE OPTICS
Prof. Samudra Roy
Department of Physics
Indian Institute of Technology Kharagpur
Lecture - 09: Group and Phase velocity

Hello, students to our course wave optics. Today we have lecture number nine and in this lecture we will discuss the concept of group and phase velocity. Okay, so let me start with the concept that we initiated in the last class. So today we have lecture number 9. So in the last class, we discussed the concept of superposition of different waves. So our basic concept was superposition of waves and in order to study that in the last few classes we superimpose two waves and their frequency the same. So now what we do today is we are going to superimpose two waves or we add two waves whose frequency is slightly different. So we will do it today. The addition of this is under the superposition of waves but with two different frequencies we will add two waves having two different frequencies. So, addition of waves with different frequencies. Suppose, we have two waves, say E_1 which is equal to $E_1 \cos(k_1 x - \omega_1 t)$. Previously, it was simply ω , but now we are going to discuss what happened if we add two waves. So, instead of having one wavelength, now we have two different wavelengths for two different waves. So, next is $E_2 \cos(k_2 x - \omega_2 t)$. So their amplitude may be considered to be the same for simplicity and their wavelength is different. So I can take their amplitude to be the same. So I can write E_1 equal to E_2 is simply say I write E_1 because it is the same. So I write E_2 as simply E_1 . So that is the condition we put that their amplitudes are the same. Now the resultant waves will be $E_1 + E_2$ since their amplitudes are the same and we call it E_1 the amplitude of E_1 wave which is the same as E_2 . So, I write E_1 and then simply write $\cos(\alpha + \cos(\beta))$ where my α is equal to $k_1 x - \omega_1 t$ and β is $k_2 x - \omega_2 t$. (Refer slide time: 07:52)

Lec No - 9
 "Superposition of wave"

- Addition of waves with different frequency

$$\Rightarrow E_1 = E_{10} \cos[k_1 x - \omega_1 t] \quad E_{10} = E_{20} \Rightarrow E_{10}$$

$$\Rightarrow E_2 = E_{20} \cos[k_2 x - \omega_2 t]$$

$$E_R = E_1 + E_2$$

$$= E_{10} [\cos \alpha + \cos \beta]$$

$$= 2 E_{10} \cos\left(\frac{\alpha + \beta}{2}\right) \cos\left(\frac{\alpha - \beta}{2}\right)$$

$$= 2 E_{10} \cos\left[\frac{(k_1 + k_2)x - (\omega_1 + \omega_2)t}{2}\right] \times$$

$$\cos\left[\frac{(k_1 - k_2)x - (\omega_1 - \omega_2)t}{2}\right]$$

$$\left. \begin{aligned} \alpha &= k_1 x - \omega_1 t \\ \beta &= k_2 x - \omega_2 t \end{aligned} \right\}$$

two different waves, their amplitudes are same that is the condition we have but their wavelengths are different and in that case if I add these two waves together then this is the resultant wave we get. Okay, now I can define so whatever the terms we can define like

This is the value of alpha and beta and now we simply try to figure out what this value is. So

simple addition of $\cos \alpha$ and $\cos \beta$ gives us $2 E_0 \cos \frac{\alpha + \beta}{2} \cos \frac{\alpha - \beta}{2}$. So $\cos \frac{\alpha + \beta}{2}$ and $\cos \frac{\alpha - \beta}{2}$ that is the next one. And now $\alpha + \beta$ and $\alpha - \beta$ I can calculate because it is already defined here and if I do we are going to get this $2 E_0$ it will be \cos of $\alpha + \beta$ and simply $k_1 + k_2$ divided by $2x$ minus $\omega_1 + \omega_2$ t divided by 2 that is the first term multiplied by \cos of $k_1 - k_2$ whole divided by $2x$ minus $\omega_1 - \omega_2$ whole divided by $2 t$. So that is the total term we have when I simply add two waves having this form, two different waves, their amplitudes are the same that is the condition we have but their wavelengths are different and in that case if I add these two waves together then this is the resultant wave we get. Okay, now I can define whatever the terms we can define like ω_p that is the sum average ω that we had. Similarly I can define k_p , which is the average of the propagation constant of these two waves divided by 2 , also we can write ω_g that is the difference between these two. So I have $\omega_1 - \omega_2$ divided by 2 and k_g . I define $k_1 - k_2$ divided by 2 . So we just defined a few more variables ω_p , k_p , ω_g , k_g but they are related to the initial frequencies that the wave had that is ω_1 , ω_2 and the initial k vector they have, the propagation constant value they have which is k_1 and k_2 . So if I write all together in this new notation then my ER becomes $2 E_0 \cos [k_p x - \omega_p t] \cos [k_g x - \omega_g t]$ that is amplitude and \cos of $k_p x - \omega_p t$ multiplied by \cos of $k_g x - \omega_g t$. So we can see that two propagating waves with different frequency when we add together then the resultant wave is also propagating but it has the amplitude. If I write this portion as an amplitude then the amplitude is also moving. What is the meaning of amplitude and what is the meaning of propagation we're going to understand here. So, before that we can see what the frequency is, here we can see that this function $\cos [k_p x - \omega_p t]$, this is a propagating wave and this is also propagating. So, initially we have, so let me picturise what we get. So, we have two waves with equal amplitude, but different frequency, so this is my say E_1 and another wave with relatively higher frequency is E_2 .

(Refer slide time: 14:54)

$$\omega_p = \frac{\omega_1 + \omega_2}{2}$$

$$k_p = \frac{k_1 + k_2}{2}$$

$$\omega_g = \frac{\omega_1 - \omega_2}{2}$$

$$k_g = \frac{k_1 - k_2}{2}$$

$$E_r = 2 E_0 \cos [k_p x - \omega_p t] \cos [k_g x - \omega_g t]$$

Carrier wave moving with a velocity $v_p = \frac{\omega_p}{k_p}$
 Envelope wave moving with a velocity $v_g = \frac{\omega_g}{k_g}$
 $\omega_p \gg \omega_g$

much greater than ω_g . So that means the frequency that the carrier wave is having is much much higher than the envelope wave we have. Now if I picturize that what is these two waves and how it should look. It will be like this.

What we did is we add these two by using the superposition principle that is we simply

linearly add these two and I get two waves with the multiplication of this cos function suggesting that this resulting wave is also propagating because it is having the form kx minus ωt form. Well, this wave is called the carrier wave having a frequency of ω_p and ω_p as I mentioned here it is ω_1 plus, ω_2 divided by. So, carrier wave is wave having a average frequency whatever the frequency we have for ω_1 and ω_2 the frequency of the carrier wave is average of that and it is moving with a velocity like if I write v_p that velocity is ω_p divided by k_p , this is the velocity at which this carrier wave is moving. Also we have another part of this wave and this part which is having relatively lower frequency. If we look at ω_g carefully, ω_g is the difference between the frequency of the two waves. So, this is called the envelope wave. This is the envelope wave and moving with a velocity I write v_g is equal to ω_g by k_g , this is the velocity at which it is moving clearly. It is given that the velocity of these two are not the same. It can be different, it can be the same in a few cases but in general this is not the same and ω_p is much greater than ω_g . So that means the frequency that the carrier wave is having is much much higher than the envelope wave we have. Now if I picture what these two waves are and how it should look. It will be like this. So, let me draw in the same figure. I'm going to draw the two waves. So in one case, let me fix the amplitude and this is a wave having high frequency, this is the one wave and this is my E_1 . Another wave I also draw whose frequency is less and it is something like this. On top of that I draw and this is my E_2 . So, if I add this E_1 and E_2 , the resultant wave will have a form like this. We have an envelope like this. And inside this envelope, we have the carrier waves. The solid line that I am drawing here will be the structure of the resultant wave that one can expect. So this is along x direction and this is moving, this is the resultant one. So resultant one means E_R , which is equal to E_1 plus E_2 and E_R is a moving wave, it is having a modulated amplitude inside this envelope. We have the distribution as well which we called the carrier wave having a frequency ω_p , which is ω_1 plus, ω_2 divided by 2

(Refer slide time: 22:30)

• Carrier wave $v_p = \frac{\omega_p}{k_p} = \frac{\omega_1 + \omega_2}{k_1 + k_2} \rightarrow E_R = E_1 + E_2$
 $\approx \frac{\omega}{k}$ (Phase velocity)

• Envelope wave $v_g = \frac{\omega_g}{k_g} = \frac{\omega_1 - \omega_2}{k_1 - k_2} = \frac{\Delta\omega}{\Delta k} \approx \frac{d\omega}{dk}$
 (Group velocity)

Now in this wave whatever the drawing we have we can see that the envelope that is the tip of this thing is moving with the velocity v_g and also the carrier wave is also moving with the velocity, which is v_p . Now from the expression we can write that v_p

and also there is a modulation in amplitude and it is periodically modulated. We call this envelope function through which it is modulated and this envelope function on envelope

waves has a frequency which is ω_1 minus, ω_2 divided by 2. So this is also moving in one direction. Now the point is what should be the velocity of this envelope that is moving. And what should be the velocity of this carrier that is moving. Already it is mentioned that the carrier wave will travel. So with the velocity, let me write here, S for carrier wave the velocity if I define v_p , this is ω_p divided by k_p and this is essentially ω_1 plus, ω_2 divided by k_1 plus, k_2 . Now this value is nearly equal to ω by k , where ω , I mean this is where ω is the average one and k is some way if I just write the average of ω and ω_2 is ω_p . So I write ω and k the average of that thing is k . So these values will be simply ω divided by k . Well on the other hand the envelope wave have a velocity v_g , which is ω_g divided by k_g and that is essentially ω_1 minus, ω_2 divided by k_1 minus, k_2 which is $\Delta\omega$ divided by Δk , where $\Delta\omega$ is a difference between these two frequency and Δk is a difference between the wave number they have. So this is essentially $d\omega$ if they are small then I can write it as $d\omega/dk$. So, this velocity has a special name it is called the Group velocity and the carrier wave which is propagating with a velocity ω divided by k is called the phase velocity. So, this is called the phase velocity. So, these two velocities we define in one case it is ω by k and another case it is $d\omega/dk$. Now in this wave whatever the drawing we have we can see that the envelope that is the tip of this thing is moving with the velocity v_g and also the carrier wave is also moving with the velocity, which is v_p . Now from the expression we can write that v_p is simply ω by k and v_g is $d\omega/dk$. So v_g I can write in terms of v_p . So v_g which is $d\omega/dk$ I can write dk multiplied by v_p into K . So, this essentially gives us v_p plus K into dv_p/dk . So this is the relation between the group velocity and the phase velocity. So group velocity is equal to phase velocity plus k multiplied by the derivative of the phase velocity with respect to k . So if we had a situation when the change of phase velocity with respect to k vanishes then one can simply have then we can have v_g is equal to v_p that means the phase and group velocity are same.

(Refer slide time: 30:12)

$$v_p = \frac{\omega}{k}, \quad v_g = \frac{d\omega}{dk}$$

$$v_g = \frac{d}{dk} (v_p k) = v_p + k \frac{dv_p}{dk}$$
 If $\frac{dv_p}{dk} = 0$ Then $v_g = v_p$. Optical wave is moving through a vacuum or free space

$$v_p = v_g = c$$

- For a dispersive medium

$$v_p = \frac{c}{n(k)} \quad v_p < c$$

$$v_g = v_p + k \frac{dv_p}{dk}$$

$$= v_p + k \frac{d}{dn} (v_p) \frac{dn}{dk}$$

$$= v_p + k \left(-\frac{c}{n^2} \right) \frac{dn}{dk} = v_p \left[1 - \frac{k}{n} \frac{dn}{dk} \right]$$

So, normally this case happens when the light or the optical wave is moving through a vacuum or a free space vacuum or free space, then v_p is equal to v_g and they are equal to c .

So then what happened was that v_p is equal to v_g and they are equal to the value of the speed of the light c but the problem is completely changed. The problem is different when we have a dispersive medium. So what is the meaning of dispersive medium? That means we have a medium with refractive index more than one. So for a dispersive medium what we have V_p is no longer C , it is basically C divided by n_k because refractive index is a function of the propagation constant or wavelength. So that means in dispersive medium v_p is less than c . So the refractive index as I mentioned is a function of wavelength or k then v_g from this expression we had. So v_g was how much it is v_p plus k into $d v_p$ divided by dk . So v_p I know what the functional form is. So that I am going to put. Okay let me do that first. So I have v_p plus k and then I need to make a derivative. So n is a function of k . So what do I do? I will write that d of $n v_p$ and d of $n dk$, so if I execute this quantity it should be v_p and then plus $k dv_p/dn$. From this expression I simply have minus of c/n^2 and then we have dn/dk . So from here I can also find that c/n , I can write it is v_p and this expression simply gives us v_p , if I take v_p common then $1 - k/n dn/dk$. But normally what do we do? We don't put dn/dk because n is a function of λ . So normally we try to write down everything in terms of λ . So this expression we slightly change as per our convenience. So let me write it down. So we have v_g is equal to $v_p [1 - k/n dn/dk]$ instead of writing what we do in this way. So k , I write as $2\pi/\lambda$. So $dk/d\lambda$ will be minus of $2\pi/\lambda^2$ and if I replace this derivative v_g will be $v_p [1 - 2k/n dn/dk]$, I replace $2\pi/\lambda$ then it is $1/n$ and this derivative I write $dn/d\lambda$ because n is essentially function of λ . We know the explicit form of n as a function of λ . We will show in the next class maybe then I write $dk/d\lambda$ which is $d/d\lambda$ rather, $d\lambda/dk$, here I write. So let me do that $d\lambda/dk$, so $d\lambda/dk$ I replace it from here. So we have v_g equal to $v_p [1 - 2\pi/\lambda \cdot 1/n \cdot dn/d\lambda]$ and this quantity is minus of λ^2 divided by 2π . Here should be a λ^2 so it should be a λ^2 by 2π and essentially we get v_p into $1 - \lambda^2/n dn/d\lambda$ and this minus minus will be plus. So plus $\lambda^2/n dn/d\lambda$, so we have $dn/d\lambda$, so this is the expression we have which shows how the quantity v_g and v_p are related to each other. That is, the group velocity and phase velocity are related to each other and we find a very important term that determines whether the group velocity is equal to phase velocity. Under that condition what happens the $dn/d\lambda$ has to be 0 that means n should not be a function of λ and it only happens when we have the value of n equal to 1 that is in free space or in air. So where the refractive index does not change with λ in that case we have v_g equal to v_p that we know but what happened when we have a component like this into the equation, that means when we deal with the velocity of the light in a medium then that medium should have a characteristic form in terms of his refractive index and if you know the refractive index as a function of λ then from this expression we can calculate the group velocity of the light. So when the light with different frequency is passing through that medium the phase and group velocity will differ and also the group velocity is a function of λ . So different wavelength will travel at different speed and that is the phenomena we called as the dispersion which is very important when we discuss the aspects of light as a wave and if it is passing through a dispersive medium then the velocity of the light for different frequency component will be different and that's why what happened we will going to get a very important phenomena called dispersion. So today we

will not have much time to discuss the phenomena of this dispersion. In the next class we are going to discuss in detail the phenomena dispersion and then we will discuss more about the material aspects of the dispersion, how the refractive index is related to lambda, and then go forward with other topics. With that note I would like to conclude here. Thank you very much for your attention and hopefully we are going to see in the next class, where we are going to discuss about the material dispersion in detail. Thank you very much and see you in the next class.

(Refer slide time: 35:11)

The image shows a video recording of a lecture slide with handwritten mathematical derivations. The equations are as follows:

$$v_g = v_p \left[1 - \frac{k}{n} \frac{dn}{dk} \right]$$
$$k = \frac{2\pi}{\lambda} \quad \frac{dk}{d\lambda} = -\frac{2\pi}{\lambda^2}$$
$$v_g = v_p \left[1 - \frac{2\pi}{\lambda} \cdot \frac{1}{n} \frac{dn}{d\lambda} \cdot \left(\frac{d\lambda}{dk} \right) \right]$$
$$v_g = v_p \left[1 - \frac{2\pi}{\lambda} \cdot \frac{1}{n} \frac{dn}{d\lambda} \cdot \left(-\frac{\lambda^2}{2\pi} \right) \right]$$
$$v_g = v_p \left[1 + \frac{\lambda}{n} \frac{dn}{d\lambda} \right]$$

The final equation is underlined, and an arrow points to the $\frac{dn}{d\lambda}$ term. In the bottom right corner, there is a small video inset of a man in a patterned shirt.