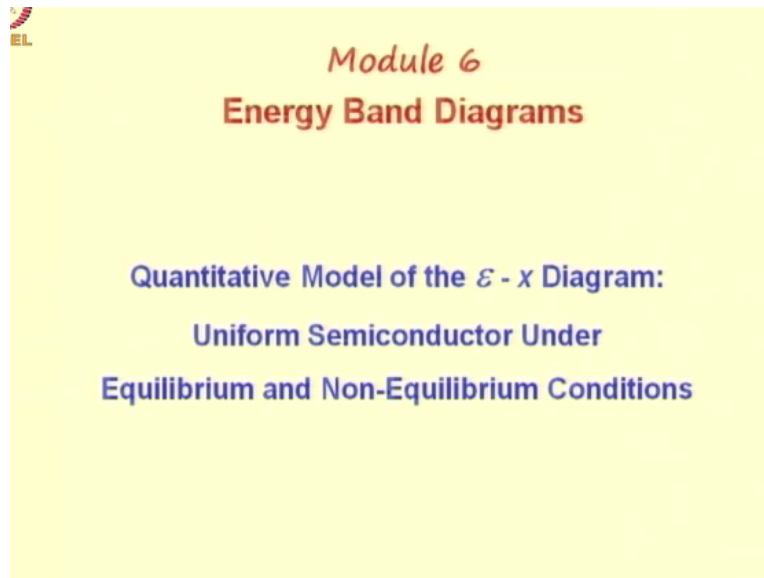


**Semiconductor Device Modeling**  
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**Lecture - 30**  
**Energy Band Diagrams**

In the previous lecture, we have discussed the utility of the E-K relation, we showed how from this relation we can obtain the crystal momentum of the electron by differentiating it we can obtain the electron velocity called the group velocity and by double differentiating it we can obtain the effective mass of the electron. We discussed the concept of crystal momentum in detail. Now in this lecture we are going to start a discussion of the energy versus distance diagram.

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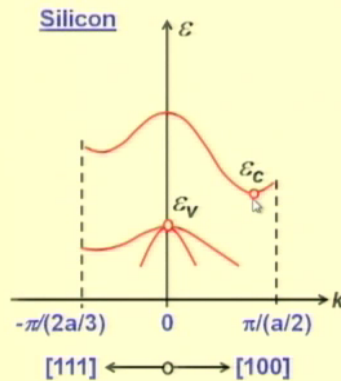


Specifically, we will discuss a quantitative model of a uniform semiconductor under equilibrium and non-equilibrium conditions, so we are going to discuss the E-X diagrams or energy distance diagrams for various conditions. In this lecture we are starting with uniform semiconductor because this is simplest case that one can think of and discussion of the energy levels in this semiconductor will give you a good foundation for drawing the band diagrams for other more complicated cases.

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## $\epsilon$ - $x$ Diagram



### Semiconductor

An  $\epsilon$  -  $x$  diagram reflects the spatial dependence of the following in a semiconductor:

- the  $\epsilon$  -  $k$  relation
- $n$ ,  $p$ ,  $\psi$  of the DD model (this yields  $J_n$ ,  $J_p$ ,  $E$  as well)
- the energy levels associated with defects and impurities

Let us look at some features of the E-X diagram, an E-X diagram reflects the spatial dependence of the following in a semiconductor, so when we are talking of a semiconductor we are talking of semiconductor volume and actually the energy levels E can change as a function of x, y, z, however in this course we will concentrate for most part on the one dimensional band diagram.

We will have approximately introduced you to 2 dimensional band diagrams but that will be only a small part of the course, so an E-X diagram reflects the spatial dependence of the following firstly the E-K relation, in the previously lecture you have come across this E-K diagram okay, now we would like to see how this particular E-K relation shown on this graph varies throughout the volume of the semiconductor as a function of position.

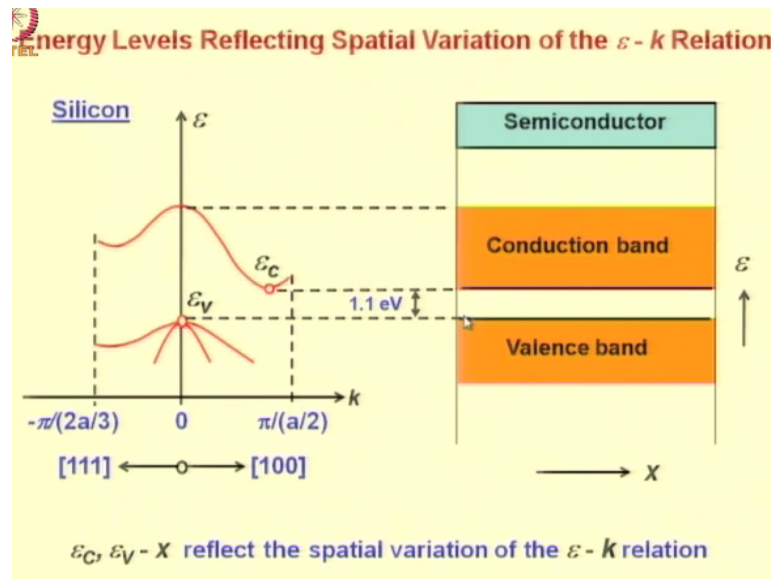
Next the E-X diagram reflects the spatial dependence of electron concentration, hole concentration and electrostatic potential which are the parameters of the drift diffusion model, now you already know that once you know n, p and psi you can also derive information about  $J_n$ ,  $J_p$  and E as well, so although there are 6 equations in the drift diffusion model these can be reduced to 3 coupled equations in n, p and psi.

So  $J_n$ ,  $J_p$  and E are quantities that can be derived from n, p and psi, thirdly the E-X diagram reflects spatial dependence of the energy levels associated with defects and impurities, now when we talk of the spatial dependence of the E-K relation it is sufficient if you focus on these 2

energies namely  $E_c$  and  $E_v$ ,  $E_c$  is the bottom of the conduction band E-K relation and  $E_v$  is the top of the valence band E-K relation.

Once you know the variation of  $E_c$  and  $E_v$  as a function of position in the semiconductor volume we can always build up the complete E-K relation around these points. Let us begin with the E-X diagram of a uniform semiconductor under equilibrium.

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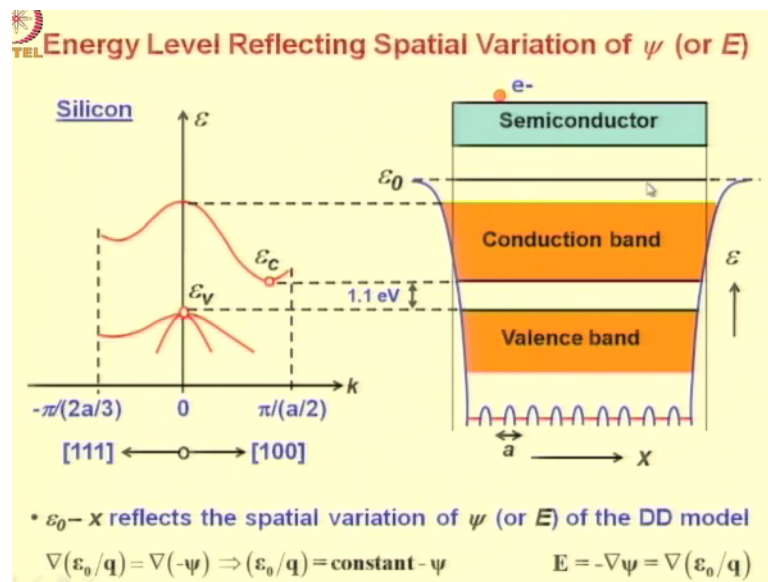
So first we talk about the energy levels, reflecting spatial variation of the E-K relation, here is the E-K relation with the critical energies namely  $E_c$  and  $E_v$  we want to plot  $E_c$  and  $E_v$  as a function of  $X$  in the semiconductor, so within the semiconductor  $E_c$  and  $E_v$  will be constant because the semiconductor is uniform and it is under equilibrium and therefore isolated from all other influences point

So no electric field is applied across the semiconductor and no other fields in fact there is no electromagnetic field applied to the semiconductor and therefore you expect the E-K relation to be uniform and therefore the  $E_c$  and  $E_v$  are constant with  $X$ , now you can also show the entire conduction band energy width and valence band energy width, because the top of the conduction band is obtained from the top of the conduction band E-K relation on this diagram here point

Similarly, you can take the E-K line correspond to the valence band it will have a bottom and you can use that bottom point to get this energy and then get the bottom of the valence band, so  $E_c$  and  $E_v$  as a function of  $X$  reflects the spatial variation of the E-K relation, we do not bother about the top of the conduction band or bottom of the valence band because once you know  $E_c$  and  $E_v$  you can always build up all other quantities.

Now let us look at the energy level reflecting spatial variation of electrostatic potential  $\psi$  or the electric field  $E$ .

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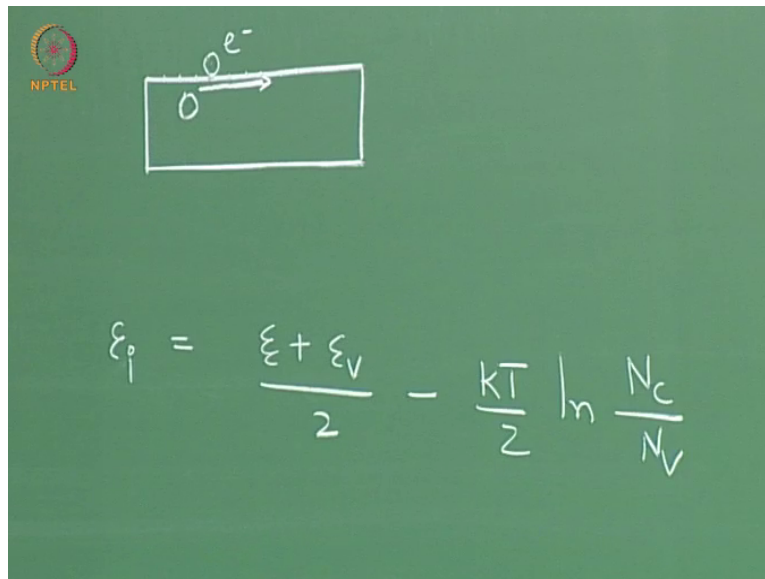


This energy level happens to be  $E_0$  which is present in the crystal potential diagram, this crystal potential diagram actually contributes to these allowed bands and energy gaps, there is also allowed band shown here which is very thin because this band is associated with electrons which are located close to the crystal atom.

So this  $E_0$  you recall from the previous lectures is the vacuum level that means it corresponds to an electron which has been moved from semiconductor outside into the vacuum but it is touching the semiconductor surface, so that is the energy level  $E_0$ , so this  $E_0$  is shown as a constant line and this  $E_0$  as a function of  $X$  reflects the spatial variation of  $\psi$  or  $E$  of the drift diffusion model.

It is shown as a constant line because there is no electric field apply to the semiconductor it is under equilibrium, therefore there is no potential variation and therefore the  $E_0$  does not change, so if I move the electron on the surface along the this line it will not experience any change in the potential because there are no fields within the semiconductor, now why should the energy of an electron touching the semiconductor surface from outside in vacuum reflect the electrostatic potential.

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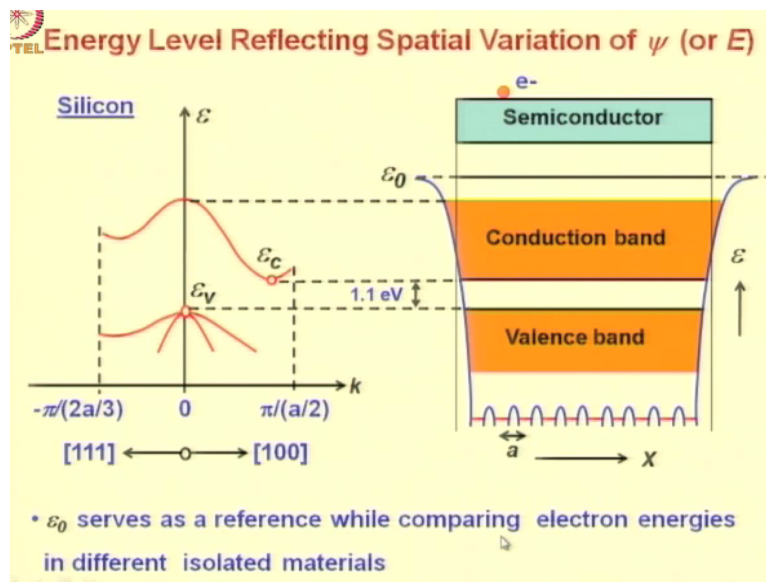
So suppose you have a semiconductor like this and there is an electric field inside then this electric field will cause a potential variation on the surface and this potential variation can be sensed by the electron which is present here, you might ask why not an electron inside will that not experience this electric field yes this electron will also experience this electric field, however this electron will also experience the effect of the atomic potentials.

Therefore, we prefer to use an electron which is touching the semiconductor surface from outside, so that it experiences the electrostatic field alone because of space charges and so on and not any electric field because of the atoms, the relation between the vacuum level  $E_0$  which is an energy and the electrostatic potential  $\psi$  of the drift diffusion model is given here, so gradient of  $E_0/q = \text{gradient of } -\psi$ .

So in fact what this means is  $E_0/q = \text{constant} - \psi$ , the negative sign comes about here because  $\psi$  is the potential associate with positive charge whereas  $E_0/q$  is the potential associate with an electron which has a negative charge, so this relation is obtained just by integrating this, now you can get the electric field as negative gradient of  $\psi$  which is nothing but gradient of  $E_0/q$ .

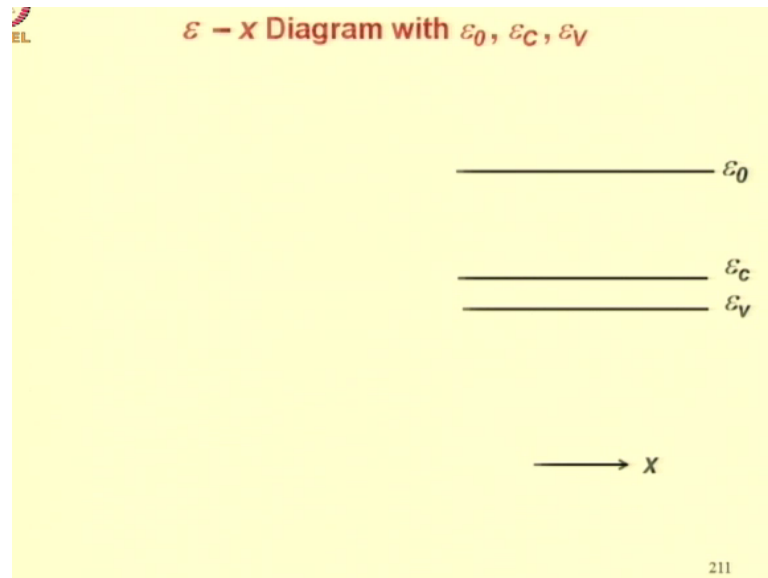
So if you take the slope of this line that should give you the electric field in this case the slope is 0 because there is no electric field because the semiconductor is in equilibrium, so no external fields are applied and since it is uniform there are no built-in fields either.

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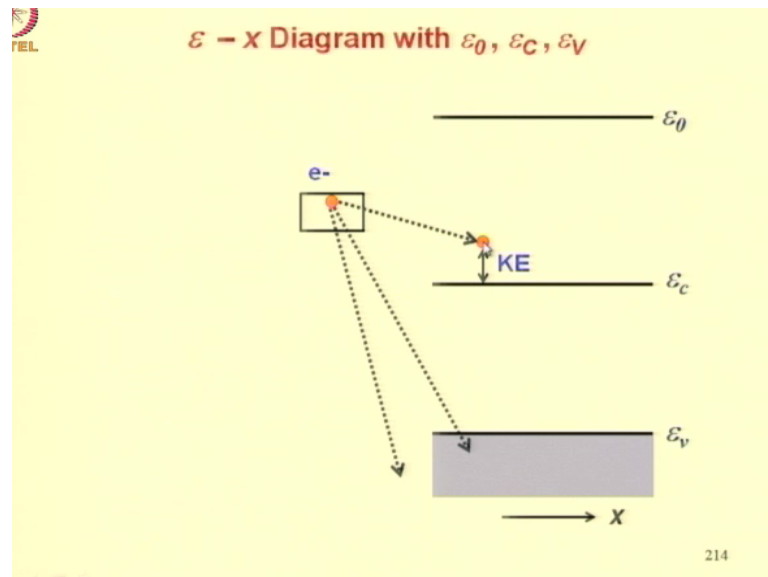
The  $E_0$  is an important energy which also serves as a reference while comparing electron energies in different isolated materials.

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Normally, we do not show the conduction, valence bands and the crystal potential etc. because it complicates the diagram, so what we do is we simply extract these 3 levels  $E_0$ ,  $E_c$  and  $E_v$  in an E-X diagram. Once we know the variation of these 3 levels as a function of X we can always build up the other details about what happens to the energy conduction band, what happens to the valence band, what happens to the crystal potential as a function of X, what happens to the E-K relation those can always be build up around these energy levels.

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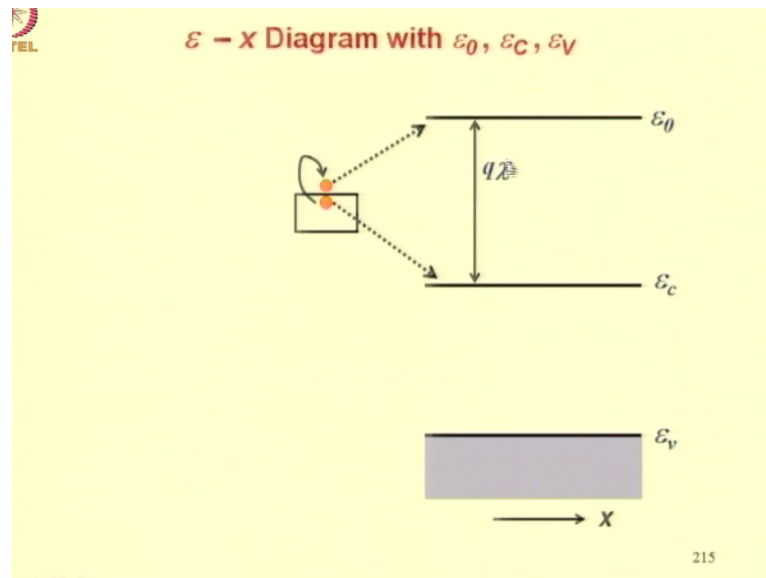


Let us increase the distance between  $E_0$ ,  $E_c$  and  $E_v$  because we want to show many more energy levels within these levels, so now the diagram is not to scale, let us familiarize ourselves with these 3 levels  $E_0$ ,  $E_c$  and  $E_v$  a little more consider an electron inside a semiconductor what

would be the energy associated with this electron, now this electron in terms of its energy could be anywhere as shown here.

For example, it could be in the conduction band, it could be in the valence band or it could be in any of the lower bands which are present, suppose electron is in the conduction band and its total energy is as shown here with this particular location, then the distance between the energy of the electron and  $E_c$  the energy level  $E_c$  is actually the kinetic energy therefore  $E_c$  is the potential energy of a conduction electron any energy higher than  $E_c$  would imply presence of a kinetic energy.

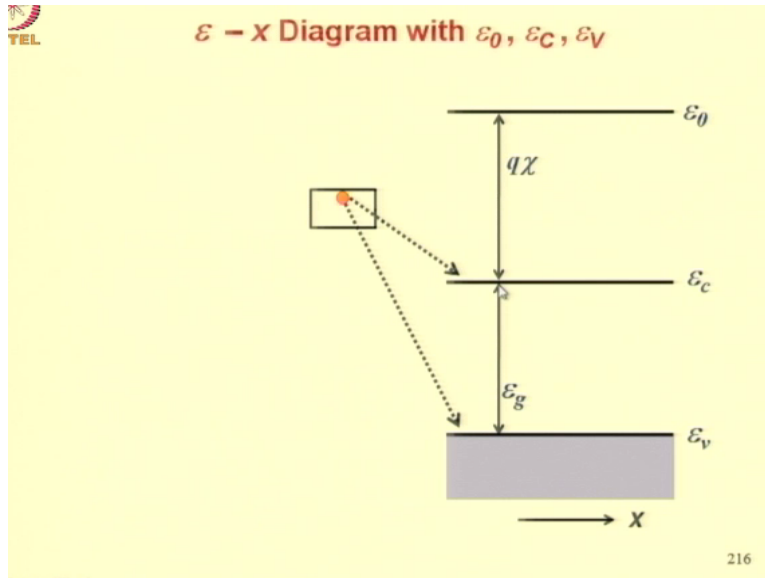
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The difference in  $E_0 - E_c$  is normally represented using the symbol  $q$  times  $\chi$ , so  $\chi$  is an potential dimension  $q \chi$  is an energy and this is called electrons affinity, now what is the physical significance of this difference if there is an electron inside the semiconductor that is at the conduction band edge and it is just moved out into the vacuum touching the semiconductor surface then its energy would change by a magnitude =  $q$  times  $\chi$ .

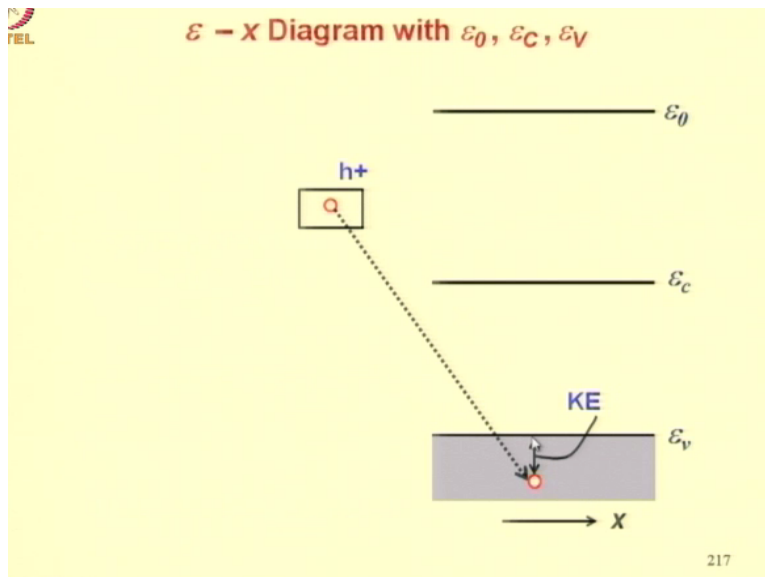
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An electron inside the semiconductor if it changes the energy from the conduction band edge to the valence band edge then it would undergo a change in energy equal to the energy gap, so  $\epsilon_c - \epsilon_v$  referred to as energy gap

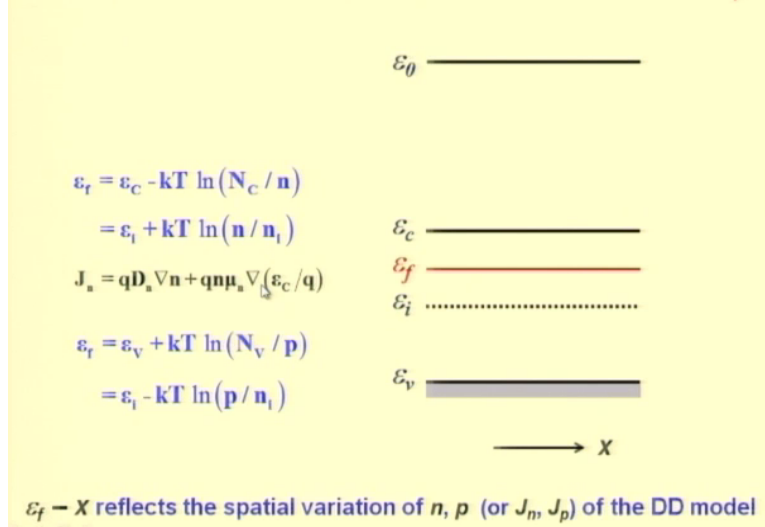
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A hole inside a semiconductor will be located in the valence band, so the kinetic energy of the hole = the energy level  $\epsilon_v$  which is the potential energy of a hole - the energy of the hole, so this difference is the kinetic energy, so as the hole gains more and more kinetic energy it moves down from  $\epsilon_v$ .

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### Energy Level Reflecting Spatial Variation of $n, p$ (or $J_n, J_p$ )



Now once we are familiarized with these levels  $E_0, E_c$  and  $E_v$ , let us now at more energy levels, first we consider energy level reflecting spatial variation of electron and hole concentrations are equivalently electron and hole current densities under equilibrium this energy level is the fermi level,  $E_f$  is the function of  $X$  reflects the spatial variation of  $n, p$  or  $J_n, J_p$  of the drift diffusion model, now let us see how does it reflect this information.

Now these are the relations for fermi level in terms of the electron and hole concentrations, so in terms of electron concentration  $E_f = E_c - kT \ln N_c/n$ , where  $N_c$  is called the effective density of states in the conduction band and similarly the same  $E_f$  can also be written as  $E_v + kT \ln n/n_i$  in terms of effective density of states in the valence band divided by the hole concentration, so here we are reflecting this different  $E_f - E_v$  in terms of the hole concentration.

And we are reflecting the difference  $E_c - E_f$  in terms of the electron concentration, these relations you would have come across already in the first level course, under equilibrium normally the symbol for electron concentration  $n$  suffix 0 and for hole concentration it is  $p$  suffix 0, so these are the 2 relations which connect the electron and hole concentrations to  $E_f$ , note that the same  $E_f$  is representing both electrons and holes.

Because this is the situation of equilibrium we will see later that in general in non-equilibrium the  $E_f$  associated with electron concentration called the  $E_f$  suffix  $n$  are quasi fermi level for

electrons is different from  $E_f$  suffix p that is the quasi-fermi level for holes which reflects the non-equilibrium hole concentration, we also have another energy level  $E_i$ ,  $E_i$  is the energy level corresponding to the fermi level in an intrinsic semiconductor.

So you already know what is the intrinsic semiconductor, it is a semiconductor which has no doping okay, now you can say why not call it pure semiconductor, so the difference between pure semiconductor intrinsic semiconductor is that a pure semiconductor has no defects or doping whereas the intrinsic semiconductor may have some defects and doping but if it has doping then its n type and p type doping will cancel each other.

So that the carrier concentration of this particular semiconductor which is not pure will be equal to the carrier concentration in a pure semiconductor, so an intrinsic semiconductor can be a pure semiconductor or an impure semiconductor whose carrier concentration = the concentration in a pure semiconductor, because the acceptor and donor type doping's cancel each other, so in an intrinsic semiconductor the fermi level will be close to the middle of the energy gap.

The reason why it is not in the middle of energy gap is because  $N_c$  and  $N_v$  are different in general, so from your first level course you would have known that  $E_i$  is given by  $E_c + E_v$  upon  $2 - kT \ln(N_c/N_v)$ , for silicon since  $N_c$  is more than  $N_v$  it means a fermi level in intrinsic semiconductor is a little below the middle of the energy gap  $E_c + E_v/2$  is the middle of energy gap okay.

Now the  $E_i$  has the following utility you can express the fermi level as fermi level with reference to  $E_i$ , so in n type semiconductor as for example the semiconductor shown here  $E_f$  is above  $E_i$  and therefore  $E_f$  is given by  $E_i + kT \ln(n/n_i)$  and the same  $E_f$  can be written in terms of hole concentration as  $E_i - kT \ln(p/n_i)$ , now how does the fermi level reflect the current density.

Now the formula for current density drift diffusion model is  $J_n = q D_n \text{grad } n + q n \mu_n \text{ into } E$ , now since it is a uniform semiconductor  $E_f$  is shown as flat line and the semiconductor is under equilibrium therefore there is no electric field, so  $E_0$ ,  $E_c$  and  $E_v$  are all flat there no slope

because no electric field and therefore uniform difference between  $E_c$  and  $E_f$  as a function of  $X$  reflects a constant electron concentration and therefore  $\nabla n$  is 0.

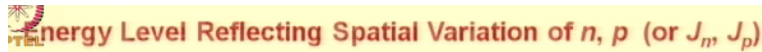
Similarly, we already remarked that the electric field is 0 because  $E_0$  is flat,  $E_c$  is flat and so on, therefore  $J_n = 0$  so that is how flat  $E_c$  and flat  $E_f$  together imply a 0 electron current, now the electric field acting on the electrons is given by gradient of  $E_c/q$ , so it is the gradient of  $E_c/q$  that reflects the electric field, now you might say why not gradient of a  $E_0$  just now we had mentioned that the gradient of  $E_0$  is related to the electric field electrostatic field. Now the point is the gradient of  $E_c$  and gradient of  $E_0$  are same in this particular case.

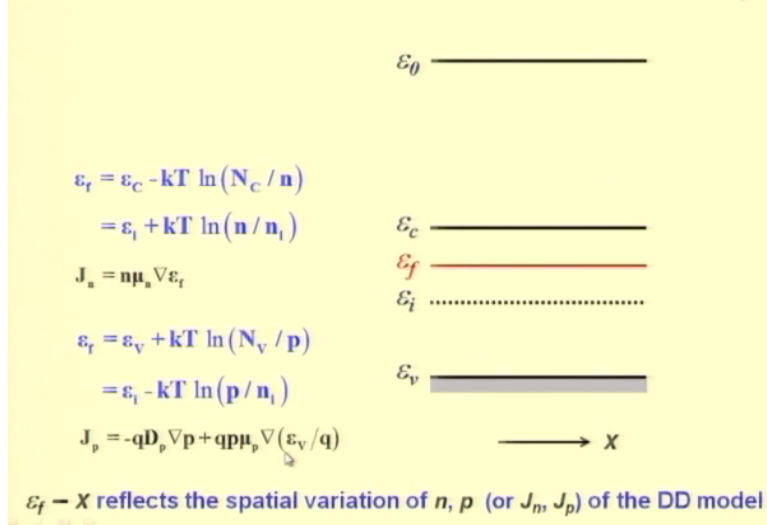
Because this is the semiconductor whose atomic structure is not changing with distance and therefore even if you were to replace  $E_c/q/E_0/q$  your electric field would have been the same. However strictly speaking the electric field that comes into this location here for the expression for the electron current is the electric field experienced by the conduction electron it is not the electric field experienced by an electron which is touching the semiconductor from outside right?

So  $E_0$  corresponds to the energy of this electron which is touching semiconductor from outside. Whereas the electric field in the  $J_n$  expression is the electric field acting on the electron in the conduction band that is the electron inside the semiconductor but in the conduction band, therefore the correct expression is, accurate expression is in terms of gradient of the  $E_c/q$  which is the potential energy which is the potential of a conduction electrons, when we go to the heterostructures where the composition changes with distance.

The atomic structure changes with distance then  $E_0$  and  $E_c$  in general would not be parallel and therefore a conflict will arise has to should you take the slope of  $E_0$  or should you take the slope of  $E_c$  for the electric field on an electron in a semiconductor, so for a conduction electron you must use  $E_c/q$ ,

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 Energy Level Reflecting Spatial Variation of  $n, p$  (or  $J_n, J_p$ )



Now we are shortly going to give you an assignment in which you can show that drift+diffusion current both the terms can be compressed in the compact form as shown in this relation  $n$  into  $\mu_n$  into gradient of  $E_f$ , so the gradient of  $E_f$  reflects both the drift as well as diffusion currents together, this compact form is very useful because it directly tells you how, it tells you how directly  $E_f$  or gradient of  $E_f$  reflects the electron current.

Assignment will be given in a short time okay, now let us discuss the current of holes current density for holes following the same approach this is the expression for hole current density which you can write as shown here where the electric field acting on the holes is represented by gradient of  $E_v/q$ , because of arguments which are similar to the arguments given for the case of electron the electric field acting on holes should be related to the gradient of the valence band edge.

In heterostructures the valence band edge may not be parallel to the conduction band edge and may not be parallel to the vacuum level  $E_0$  and therefore in such a situation it is important to note that the electric field in the hole current expression is given by the gradient of  $E_v/q$  and you cannot use  $E_0$  here nor can you use  $E_c$ .

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### Energy Levels Reflecting $\psi, n, p$ (or $E, J_n, J_p$ )

$$(\epsilon_0/q) = \text{constant} - \psi$$

$$\epsilon_0 \text{ —————}$$

$$\mathbf{E} = -\nabla\psi = \nabla(\epsilon_0/q)$$

$$\epsilon_f = \epsilon_c - kT \ln(N_c/n)$$

$$= \epsilon_i + kT \ln(n/n_i)$$

$$\epsilon_c \text{ —————}$$

$$\mathbf{J}_n = n\mu_n \nabla \epsilon_f$$

$$\epsilon_f \text{ —————}$$

$$\epsilon_i \text{ .....}$$

$$\epsilon_f = \epsilon_v + kT \ln(N_v/p)$$

$$= \epsilon_i - kT \ln(p/n_i)$$

$$\epsilon_v \text{ —————}$$

$$\mathbf{J}_p = p\mu_p \nabla \epsilon_f$$

$$\longrightarrow x$$

$\epsilon_0$  reflects  $\psi$  (or  $E$ ) and  $\epsilon_f$  reflects  $n, p$  (or  $J_n, J_p$ ) of the DD model

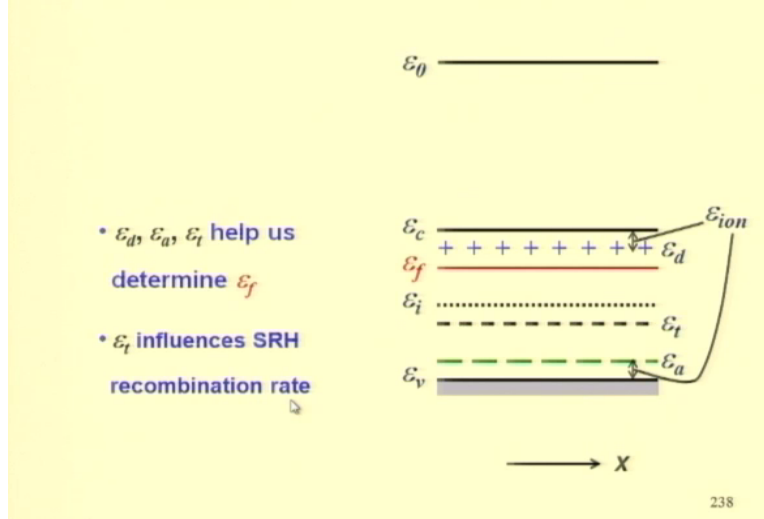
Once again in analogy to the case of electron current density, the hole current density can be expressed in terms of gradient of the fermi level, since it is in equilibrium case the same fermi level is entering into the expression for holes as well as for electrons, otherwise as the assignment will show what you will have in the expression for current density of electrons would be the quasi-fermi level electrons  $E_{fn}$ .

And what you will have in the expression for  $J_p$  would be quasi-fermi level for holes that is  $E_{fp}$ , you can see here that seems the fermi level is flat, the hole current is 0 and similarly electron current is 0, so this is very straight forward reasoning based on the gradient of the fermi level for deriving the information about the electron and hole current densities, now let us complete the picture let us show all the energy levels which reflects  $\psi$   $n$  and  $p$  or  $E, J_n, J_p$ .

So reproducing from our previous discussion the  $E_0$  reflects the electrostatic potential  $\psi$ , it also reflects the electric field  $E$  according to this relation, so we summarize by saying  $E_0$  reflects  $\psi$  or  $E$  and  $E_f$  reflects  $n$  or  $p$  or equivalently  $J_n$  or  $J_p$  of the drift diffusion model, so this is how the levels  $E_0, E_c, E_f$  and  $E_v$  together represent the conditions in a semiconductor device those conditions which are used in a drift diffusion model and these levels are also useful to give you the spatial variation of the E-K relation.

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## Energy Levels Reflecting Crystal Defects and Impurities



Let us add now to this diagram energy levels reflecting crystal defects and impurities, the donor level corresponding to a shallow donor would be somewhere here close to the in energy gap close to the conduction band edge, now it is not necessary at a donor level should always be close to the conduction band edge only a shallow donor that is that kind of a donor which gives you a large number of electrons in a semiconductor.

And converts the semiconductor into n type would be a shallow donor its energy level will be close to the conduction band edge, example of this is phosphorus or antimony or arsenic in silicon, the positive sign here indicates donors because it means that when this impurity loses its electron it is positively ionized so donor is one which gives an electron right and in process becomes positively charged so that is why positive signs are used to indicate the donor level.

Similar way one can show the acceptor level with negative signs this is the shallow acceptor such as boron or aluminium in silicon, so it is located close to the valence band edge, once again in acceptor level would lie anywhere in the energy gap a shallow acceptor would be close to the valence band edge, a trap level apart from donors and acceptors you can also have defects.

And you can also have apart from shallow donors and acceptors other types of impurities which create energies near the middle of energy gap, so either defects or impurities both can give us to the energies near the middle of the energy gap and these are reflected here by this energy level

call the trap level it is called a trap because these impurities can actually physically trap electrons or holes.

The distance between the conduction band edge and the donor level is referred to as ionization energy of the donor and the distance between the valence band edge and the acceptor level is referred to as ionization energy of the acceptor, for example if you take this trap level if it is acceptor type as it is shown here with negative signs then its ionization energy would be the difference between  $E_t$  and  $E_v$ .

If the same trap level was donor type the ionization energy would be  $E_c - E_t$ , the donor level the acceptor level and the trap level help us determine  $E_f$ , so this levels are useful in determination of  $E_f$ , the trap level  $E_t$  influences the Shockley Read Hall recombination rate, in our early lectures we have discussed about Shockley Read Hall recombination and there you would recall we had used an energy level called this trap level.

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**Determination of  $\epsilon_f$  and  $\epsilon_g$  Including Heavy Doping Effects**

**Charge balance eqn.**  

$$p_0 + N_d^+ = n_0 + N_a^-$$
 links  $n, p, \epsilon_f$  to  $N_d, N_a, \epsilon_d, \epsilon_a$

**Joyce-Dixon Approximation**

$$p_0 = \frac{N_v \exp[(\epsilon_v - \epsilon_f)/kT]}{1 + 0.25 \exp[(\epsilon_v - \epsilon_f)/kT]}$$

$$n_0 = \frac{N_c \exp[(\epsilon_f - \epsilon_c)/kT]}{1 + 0.25 \exp[(\epsilon_f - \epsilon_c)/kT]}$$

$$N_d^+ = \frac{N_d}{1 + 2 \exp[(\epsilon_f - \epsilon_d)/kT]}$$

$$N_a^- = \frac{N_a}{1 + 4 \exp[(\epsilon_a - \epsilon_f)/kT]}$$

$$\epsilon_g = \epsilon_{g0} - \Delta \epsilon_g(T) - \Delta \epsilon_g(N)$$

Now let us discuss how you determine the  $E_f$  and  $E_g$  in a uniform semiconductor under equilibrium including heavy doping effects, in the first level course you would have determined  $E_f$  in a dope semiconductor, what we want to do here is to so slightly improve upon the formula and include heavy dropping effects when the energy gap can decrease as compared to low doping.

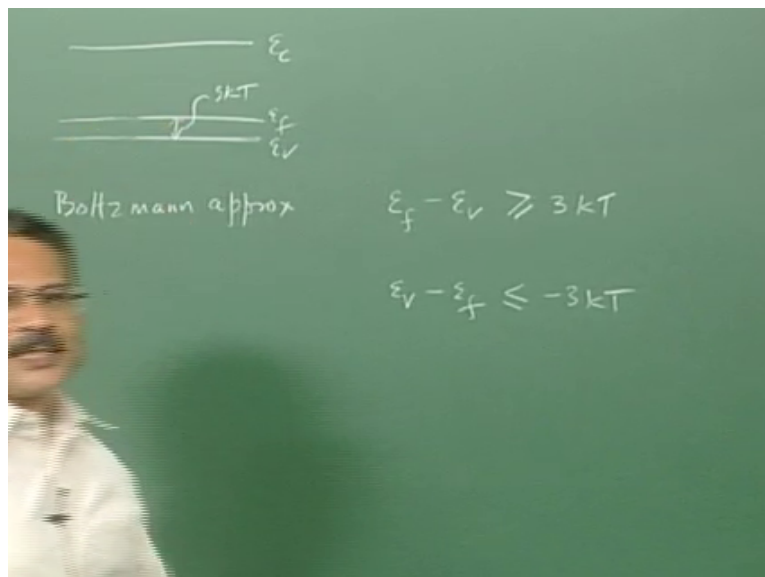


And Fermi Dirac statistics sorry the Boltzmann statistics for electrons may not be valid and you will have to use Fermi Dirac statistics, so this is your energy level diagram relevant to this problem, the determination of  $E_f$  starts from the charge balance equation that is  $p_0 + N_d^+ = n_0 + N_a^-$ , you have come across this equation in the first level course, the suffix 0 indicates the fact that we are considering equilibrium condition  $N_d^+$  is a concentration of ionized donors

And  $N_a^-$  is the concentration of ionized acceptors, now this charge balance equation links electron concentration, hole concentration and fermi level to the donor doping, acceptor doping the donor level and the acceptor level, let us see how now we discuss the Joyce-Dixon approximation to the Fermi Dirac statistics.

According to this approximation the hole concentration is given by  $N_v$  which is the effective density of states in the valence band into exponential of  $E_v - E_f/kT$  divided by  $1 + 0.25$  exponential  $E_v - E_f/kT$ , the Boltzmann approximation is  $p_0 =$  the numerator part of this expression, so this expression tells you that as your doping becomes heavy acceptor doping's becomes heavy the  $E_f$  will start moving closer and closer to the valence band and in such a case this  $E_v - E_f$  can become smaller in magnitude.

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Now Boltzmann approximation is valid when  $E_v - E_f$  is  $\geq 3$  times  $kT$  or  $E_v - E_f \leq -3$  times  $kT$ , so this is your band diagram and this is your  $E_f$ , this is  $E_v$ , this  $E_c$ , so as your doping becomes heavy this starts coming closer and closer if the  $E_f$  gets closer than 3 times  $kT$  then you know that the Boltzmann approximation fails now that is where the Joyce-Dixon approximation becomes important okay.

So the Joyce-Dixon is such that for low doping it will be reduced to Boltzmann approximation you can see from when the doping is low  $E_v - E_f$  is really a large negative quantity and therefore this denominator term simply becomes about 1 and you have only the numerator term, whereas for example when  $E_f$  becomes  $= E_v$  for heavy doping in such a case  $p_0$  becomes  $= N_v/1.25$  whereas the Boltzmann approximation will only give you  $N_v$  okay.

So there is 25% error almost due to expression based on Boltzmann approximation when the doping is a heavy and  $E_f$  is at  $E_v$ , a similar approximation exists for electron concentration and it is given by  $n_0 = N_c \exp(E_f - E_c/kT)$  divided by  $1 + 0.25 \exp(E_f - E_c/kT)$ , so when  $E_f = E_c$  the Joyce-Dixon approximation predicts  $n_0 = N_c/1.25$  whereas the Boltzmann approximation which is simply the numerator part of this expression predicts only  $N_c$ .

So there is a 25% error due to Boltzmann approximation when  $E_f = E_c$ , the expression for  $N_d$  in this equation is given by  $N_d/1 + 2 \exp(E_f - E_d/kT)$ .

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$$N_d^+ = [1 - f'(E_d)] N_d$$

$$f'(E_d) = \frac{1}{1 + 4 \exp\left(\frac{E_d - E_f}{kT}\right)}$$

This you would have come across in the first level course  $N_d^+ = 1 -$  the fraction of occupied states at  $E_d$  into the total concentration of donors, now this fraction is not given by the Fermi Dirac function exactly so that is why you put a dash so Fermi Dirac function would be  $1/1 + \text{exponential of } E_d - E_f/kT$  but when you are applying the function to impurity level there is always a constant coming here it is called degeneracy factor.

And that is why you use a different function than the Fermi Dirac function and because of this factor that you have 2 coming in here  $N_a^-$  is similarly given by  $N_a / (1 + 4 \text{ times exponential } E_a - E_f/kT)$ , again this 4 is because of the degeneracy factor, so  $N_a^-$  is nothing but  $N_a$  into the fraction of occupied states at  $E_a$ .

Finally, the energy gap is given by the energy gap for low doping and low temperature - a change as a function of temperature - a change as a function of doping this is really referred to as the band gap narrowing because of heavy doping effects.

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## Determination of $\epsilon_f$ and $\epsilon_g$ Including Heavy Doping Effects

Charge balance eqn.

$$p_0 + N_d^+ = n_0 + N_a^-$$

links  $n, p, \epsilon_f$  to  $N_d, N_a, \epsilon_d, \epsilon_a$

Steps for solution

1)  $\epsilon_f$

2)  $p, n$

Joyce-Dixon Approximation

$$p_0 = \frac{N_v \exp[(\epsilon_v - \epsilon_f)/kT]}{1 + 0.25 \exp[(\epsilon_v - \epsilon_f)/kT]}$$

$$N_d^+ = \frac{N_d}{1 + 2 \exp[(\epsilon_f - \epsilon_d)/kT]}$$

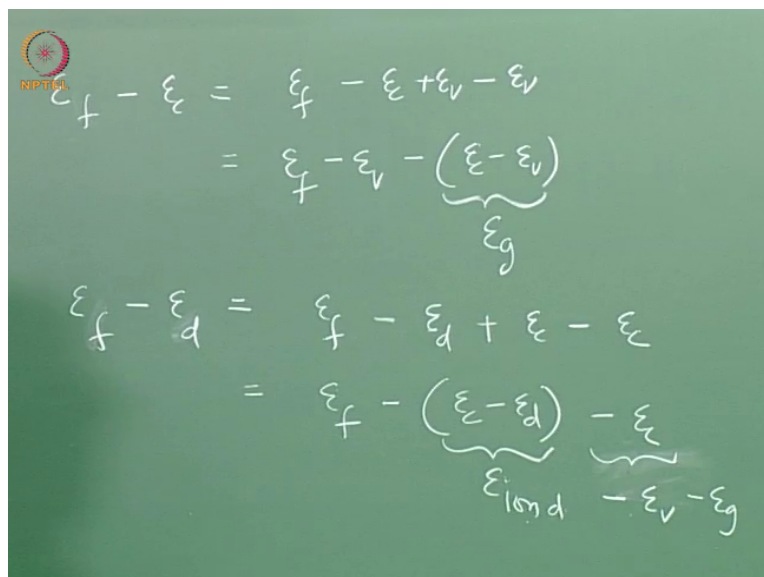
$$n_0 = \frac{N_c \exp[(\epsilon_f - \epsilon_c)/kT]}{1 + 0.25 \exp[(\epsilon_f - \epsilon_c)/kT]}$$

$$N_a^- = \frac{N_a}{1 + 4 \exp[(\epsilon_a - \epsilon_f)/kT]}$$

$$\epsilon_g = \epsilon_{g0} - \Delta\epsilon_g(T) - \Delta\epsilon_g(N)$$

Now the way this equation is solved is that first you solve for  $E_f$  and then you solve for electron and hole concentrations okay, so what you do is you take these expressions substitute for  $p_0$  take this expression substitute for  $n_0$ , similarly substitute these 2 things here for  $N_d^+$  and  $N_a^-$  and then solve for  $E_f$ , now you might wonder how do you solve for  $E_f$  when you have  $E_v$  and  $E_c$ ,  $E_d$  and  $E_a$  also present there.

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$$\begin{aligned} \epsilon_f - \epsilon &= \epsilon_f - \epsilon + \epsilon_v - \epsilon_v \\ &= \epsilon_f - \epsilon_v - \underbrace{(\epsilon - \epsilon_v)}_{\epsilon_g} \\ \epsilon_f - \epsilon_d &= \epsilon_f - \epsilon_d + \epsilon - \epsilon \\ &= \epsilon_f - \underbrace{(\epsilon - \epsilon_d)}_{\epsilon_{Ion d}} - \underbrace{\epsilon}_{-\epsilon_v - \epsilon_g} \end{aligned}$$

So here what you need to do is you take any energy as a reference let us see if I choose to take  $E_v$  as reference, so in other words all expressions which are in terms of other energy levels like  $E_c$ ,  $E_a$ ,  $E_d$  should be converted to energy level in terms of  $E_v$ , let us see how

supposing you have for example this expression  $E_f - E_c$ , so you can write  $E_f - E_c = E_f - E_c + E_v - E_v$ , so add and subtract  $E_v$  because this is your reference.

Now what will happen is this can be written as  $E_f - E_v - E_c + E_v$  and this is nothing but energy gap and you have an expression for energy gap in fact we are shortly going to give you expression for energy gap, so that is how  $E_f - E_c$  can be converted into  $E_f - E_v$ , similarly supposing you want to convert  $E_d - E_c$  sorry you want to convert  $E_f - E_d$  into an expression consisting of the valence band.

So  $E_f - E_d$  so you can write this as  $E_f - E_d + E_c - E_c$  that =  $E_f - E_c - E_d$  which you can identify as ionization energy of the donor so that is why I put a suffix d because ionization energy of the donor can be different from ionization energy of the acceptor, so  $- E_c$  now this  $- E_c$  you can always convert it to  $E_v$  as shown here, so you can see  $E_f - E_c$  was  $E_f - E_v$  - this, so I can write this part as  $- E_v - E_g$  so your  $E_f - E_d$  is  $E_f$  - ionization energy of the donor -  $E_v - E_g$ .

So like that you can convert each expression into an expression containing  $E_f$  and  $E_v$ , then you can solve for  $E_v - E_f$ .

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**Determination of  $\epsilon_f$  and  $\epsilon_g$  Including Heavy Doping Effects**

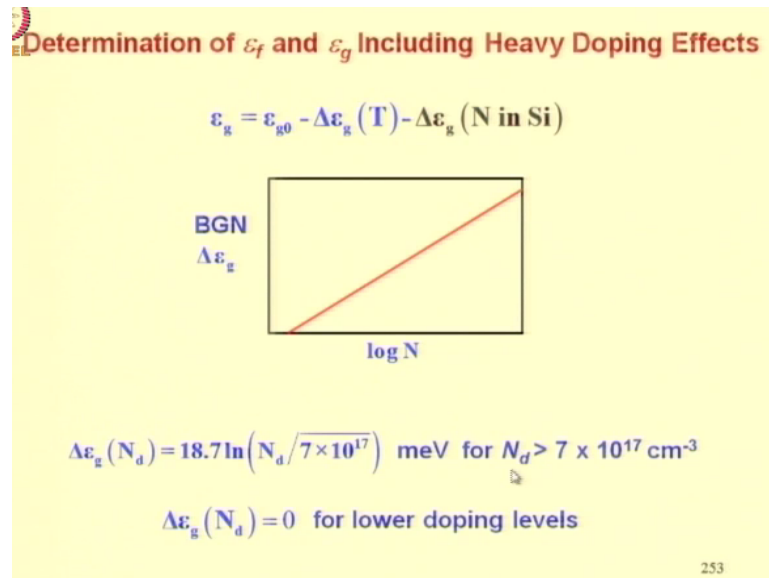
$$\epsilon_g = \epsilon_{g0} - \underbrace{\Delta\epsilon_g(T) - \Delta\epsilon_g(N)}_{\alpha T^2 / (\beta + T)}$$

Material	$\epsilon_{g0}$ (eV)	$\alpha$ (eV/K)	$\beta$ (K)
Si	1.169	$4.9 \times 10^{-4}$	55
GaAs	1.519	$5.4 \times 10^{-4}$	204

Now the expression for energy gap is as follows, let us look at the term that depends on the temperature is given by  $\alpha T^2 / (\beta + T)$  were for different materials the constants

have different values, so for silicon the  $E_{g0}$  that is this term is 1.169 electron volts alpha its units are electron volt per kelvin it is  $4.9 \times 10^{-4}$  and beta has the unit kelvin at 655, for gallium arsenide the values are 1.519 for  $E_{g0}$ ,  $5.4 \times 10^{-4}$  for alpha and 204 for beta, so alpha is of the order of  $10^{-4}$  electron volt per kelvin.

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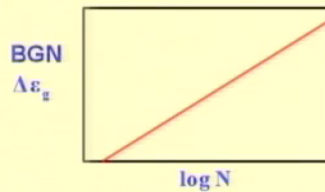
Let us look at the band gap narrowing term for silicon, so N stands for doping it could be donor or acceptor you have different expression for donor and acceptor it turns out that for both donor and acceptor  $\Delta E_g$  as a function of  $\log N$  is close to a straight line. And therefore the form of expression is as follows  $\Delta E_g$  as a function of  $N_d$  if you have donors then it =  $18.7 \times \log$  of  $N_d$  divided by  $7 \times 10^{17}$  where this quantity or  $N_d$  is in per centimeter cube so milli-electron volts for  $N_d > 7 \times 10^{17}$ , now for  $N_d < 7 \times 10^{17}$ .

This logarithmic term will be negative it might give you a feel that the  $\Delta E_g$  becomes negative which is not correct. So for lower doping levels we assume  $\Delta E_g$  is 0, so this formula is valid only for  $N_d > 7 \times 10^{17}$ , so this means in silicon for donor doping  $> 7 \times 10^{17}$  band gap will start narrowing.

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**Determination of  $\epsilon_f$  and  $\epsilon_g$  Including Heavy Doping Effects**

$$\epsilon_g = \epsilon_{g0} - \Delta\epsilon_g(T) - \Delta\epsilon_g(N \text{ in Si})$$



$$\Delta\epsilon_g(N_a) = 9 \left( F + \sqrt{F^2 + 0.5} \right) \text{ meV}$$

$$F = \ln \left( N_a / 10^{17} \right) \text{ for } N_a > 1 \times 10^{17} \text{ cm}^{-3}$$

$$\Delta\epsilon_g(N_a) = 0 \text{ for lower doping levels}$$

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The formula for acceptor type case has given here  $\Delta E_g$  as the function of  $N_a = 9$  times  $F + \sqrt{F^2 + 0.5}$  milli-electron volt where  $F$  is logarithm of  $N_a$  divided by  $10^{17}$  where  $N_a$  is in per centimeter cube, so evidently this formula is valid for  $N_a > 1 \times 10^{17}$  centimeter cube and for lower doping levels  $\Delta E_g$  is assumed to be 0, because then logarithmic becomes negative which is not physical correct.

**(Refer Slide Time: 42:46)**

**Determination of  $\epsilon_f$  and  $\epsilon_g$  Including Heavy Doping Effects**

*Assignment-6.2*

*Derive closed-form relations in terms of  $N_c, N_v, N_d, \epsilon_{ion}, T$  for  $n, p, \epsilon_f$  in an n-type semiconductor.*

*The relations should work even for high doping levels for which the impurity ionization may be partial and Boltzmann approximation for the Fermi-Dirac statistics may breakdown.*

*State the approximations made.*

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So here is an assignment for you derive closed form relations in terms of  $N_c, N_v, N_d$  and  $E_{ion}$  and temperature for electron concentration, hole concentration and fermi level in an n-type semiconductor, so you have to do the charge balance equation the Joyce-Dixon approximation and the equation for concentration of ionized impurities.

The simplification that will result here is that this is an n-type semiconductor therefore you do not have to bother about hole concentration and ionized acceptor impurities that simplifies the expression a lot and it allows you to derive an analytical expression if you want to do a calculation including  $N_d$  and  $N_a$  in a general case then you will have to resolve to numerical calculations.

The relations should work even for high doping levels for which impurity ionization maybe partial and Boltzmann approximation for the Fermi Dirac statistics may breakdown, state the approximations made.

**(Refer Slide Time: 43:54)**

**Determination of  $\epsilon_f$  and  $\epsilon_g$  Including Heavy Doping Effects**

*Assignment-6.3*

*Using the above relations, estimate  $n$ ,  $p$  and  $\epsilon_f$  at 300 K in n-type AlGaAs and Si samples having  $N_d = 5 \times 10^{18} \text{ cm}^{-3}$  and  $\epsilon_{ion}(\text{donor}) = 0.045 \text{ eV}$ . Compare these results with those obtained using complete ionization and Boltzmann approximation.*

	$N_c (\text{cm}^{-3})$	$N_v (\text{cm}^{-3})$
AlGaAs	$4.7 \times 10^{17}$	$7.1 \times 10^{18}$
Si	$2.8 \times 10^{19}$	$1.8 \times 10^{19}$

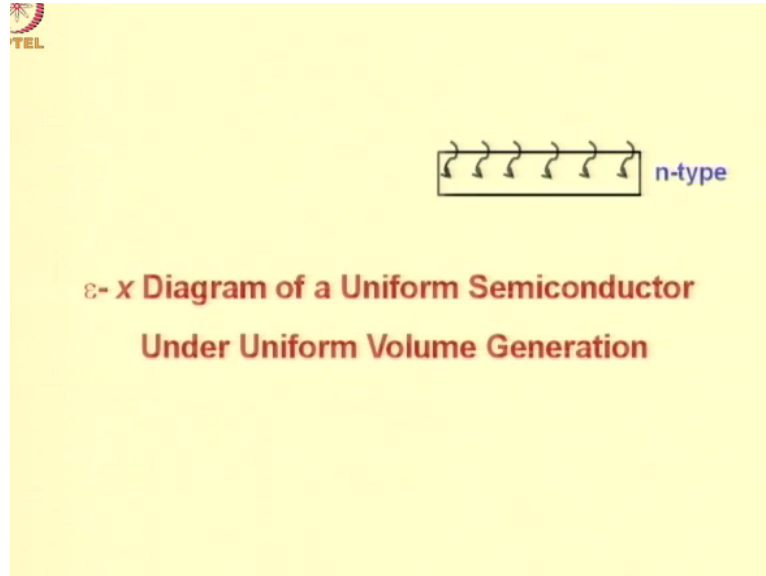
There is the another assignment where you use the derived the expression for calculations, using the above relations estimate  $n$   $p$  and  $E_f$  at 300 kelvin in n-type aluminium gallium arsenide and silicon samples having  $N_d = 5$  into  $10$  power  $18$  per centimeter cube and ionization energy of donor =  $0.045$  electron volts, compare these results with those obtained using complete ionization and Boltzmann approximation.

Here are the values of  $N_c$  and  $N_v$  for aluminium gallium arsenide and silicon, for aluminium gallium arsenide  $4.7$  into  $10$  power  $17$  per centimeter cube and for silicon  $2.8$  into  $10$  power  $19$  per centimeter cube, for  $N_c$  and similarly for  $N_v$  the values are  $7.1$  into  $10$  power  $18$  per



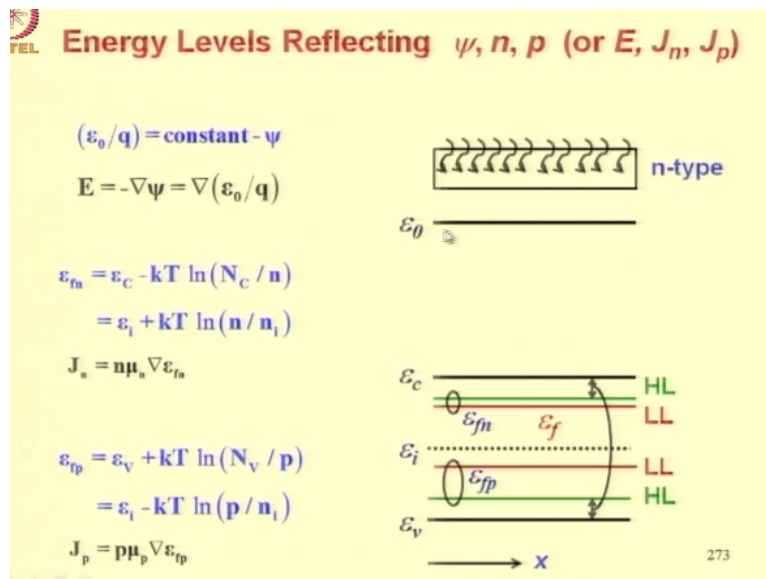
centimeter cube for aluminium gallium arsenide and  $1.8 \times 10^{19}$  per centimeter cube for silicon.

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Finally, let us discuss the E-X diagram of a uniform semiconductor under uniform volume generation.

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So we want to show energy levels reflecting  $\psi, n$  and  $p$  of the drift diffusion model which will also reflect  $E, J_n$  and  $J_p$ , let us begin from the equilibrium condition for which we already know the band diagram it contains  $E_0, E_c, E_f, E_i$  and  $E_v$ ,  $E_f$  is above  $E_i$  because the semiconductor is n

type, now when you shine light and there is a uniform volume generation in this semiconductor then everywhere there is excess electrons and holes.

And situation is depicted in an energy diagram by splitting of fermi level into electron quasi fermi level and hole quasi fermi level it turns out that when the injection level is low LL stands for low level when the injection level is low the quasi fermi level for electrons is same as or almost the same as the equilibrium fermi level the fermi level under equilibrium, whereas the quasi fermi level for holes moves significantly away from the equilibrium fermi level, now let us understand why?

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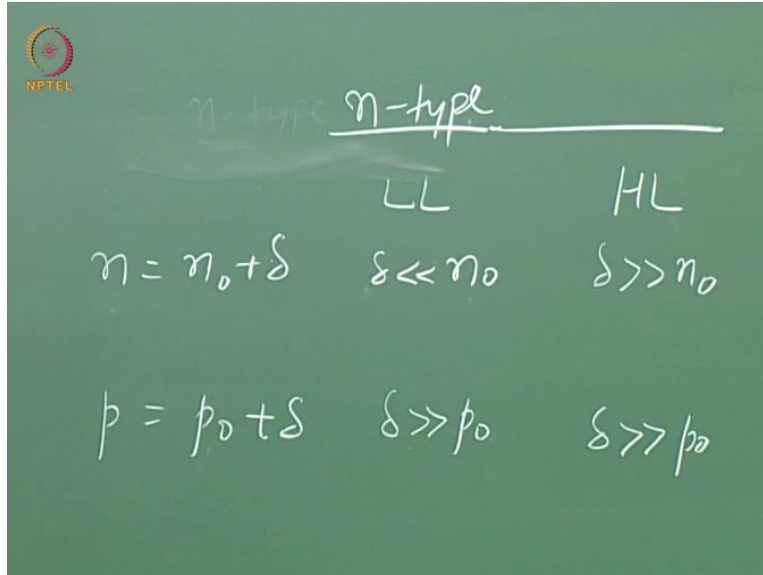
$$E_{fn} = E_c - kT \ln \frac{N_c}{n} \quad n = n_0 + \delta$$

$$E_{fp} = E_v + kT \ln \frac{N_v}{p} \quad p = p_0 + \delta$$

So using the concept of quasi fermi level  $E_{fn} = E_c - kT \ln \frac{N_c}{n}$  upon  $n$  and  $n = n_0 + \delta$  and similarly for holes it is  $E_v + kT \ln \frac{N_v}{p}$  upon  $p$ ,  $p = p_0 + \delta$  excess carrier concentration is the same for electrons and holes, now you can see from here if  $n$  increases this term reduces which means the fermi level goes closer to  $E_c$ , on the other hand if  $p$  increase compared to  $p_0$  this term decreases and according to this formula the quasi fermi level for holes goes close to  $E_v$ .

So the same fermi level now cannot represent both electrons and holes because quasi fermi level for electron should go to closer to  $E_c$  whereas the fermi level for holes should go closer to  $E_v$ .

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Low level conditions delta is much  $< n_0$  however delta is always  $> p_0$  because  $p_0$  is the minority carrier concentration in n type semiconductor, so this very, very small so generally delta will be much  $> p_0$  and therefore this quantity will significantly deviate from  $E_v$  because p changes significantly from  $p_0$ , whereas n remains close to  $n_0$  because if the delta is much  $< n_0$ , n is  $n_0$  itself and therefore this remains close to the equilibrium fermi level.

Now, the carrier concentration are therefore represented these 2 formulae which we just now wrote down on the board they could also be represented using an equation in which the reference is  $E_i$  for electrons and holes rather than  $E_c$  for electrons and  $E_v$  for holes, now if you increase the intensity of light you can have high level conditions, high level conditions there is a small change in the electron fermi level so it moves closer to the  $E_c$  as compared to equilibrium.

And the hole quasi fermi level moves closer to  $E_v$  also in a similar way, now what happens in high level is that if I write expression for high level for the same n type then here delta is becoming much more than  $n_0$  and anyway it is much more than  $p_0$ , so therefore in this case you have  $n$  approximately = delta and  $p$  approximately = delta both n and p are almost the same and = delta.

And therefore if you put delta here and here the difference between  $E_c$  and  $E_{fn}$  and difference between  $E_v$  and  $E_{fp}$  would be almost the same they are not exactly the same because  $N_c$  is

slightly different from  $N_v$  okay otherwise they are almost the same, so therefore the high-level condition is depicted by this situation so this difference equal to this difference approximately, now let us complete the picture because we want to show  $\psi$  also.

So the electrostatic potential  $\psi$  is reflected by  $E_0/q$  and from these levels  $E_0$ ,  $E_f$  and  $E_{fp}$  you can also get electric field, the electron current density and hole current density, so these expressions are repeated from the previous slide, so in this case for example  $E_{fn}$  and  $E_{fp}$  both are flat which means that electron current is 0 and hole current is 0 it is expected because electron and hole concentration are uniform throughout the volume.

Therefore, there cannot be any diffusion current and there are there is no electric field between the semiconductor because it is uniform and no external electric field is applied therefore drift current is also 0, so that is how  $J_n$  and  $J_p$  are 0 we can derive the information directly from slope of a  $E_{fn}$  and  $E_{fp}$ , the electric field is obtained from the slope of  $E_0$ .

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**Current Density in Terms of Quasi-fermi Level**

*Assignment-6.4*

*Show that, the current density of the DD model*

$$J_n = qD_n \nabla n + qn\mu_n E$$

*can be expressed in the compact form*

$$J_n = n\mu_n \nabla \epsilon_{fn}$$

*using the relations*

$$n = N_c e^{-(\epsilon_c - \epsilon_{fn})/kT} \quad E = (1/q) \nabla \epsilon_c \quad D_n = (kT/q) \mu_n$$

*Derive an analogous relation for  $J_p$  in terms of  $\epsilon_{fp}$ .*

Here is an assignment current density in terms of quasi fermi level, show that the current density of the drift diffusion model  $J_n = q D_n \text{grad } n + q n \mu_n E$  can be expressed in the compact form  $J_n = n \mu_n \text{grad } E_{fn}$  using the relations  $n = N_c \exp(-E_c - E_{fn}/kT)$  and  $E = 1/q \text{ gradient of } E_c$ ,  $D_n = kT/q \mu_n$  this is an Einstein relation.

So you use these 3 facts and substitute this here and you should show by algebraic manipulation that this equation containing 2 terms drift and diffusion reduce as to a single term in terms of gradient of  $E_{fn}$ , derive an analogous expression for  $J_p$  in terms of  $E_{fp}$ , when you derive the expression for  $J_p$  you must use gradient of  $E_v$  here instead of gradient of  $E_c$ , we have already explained why you should not use gradient of  $E_0$ .

Because for using for the electric field term in the current density expressions because in general in heterostructures  $E_0$ ,  $E_c$  and  $E_v$  all may not be parallel, so for conduction electrons you must use gradient of  $E_c$  as electric field acting on electrons, for holes you must use gradient of  $E_v$  acting as the electric field acting on holes, with that we have come to the end of the lecture.

Let us make a summary of the important points, so in this lecture we have begun a discussion of the energy versus distance diagram we explained that this energy versus distance diagram contains following critical energy levels these  $E_0$  the vacuum level,  $E_c$  the conduction band edge,  $E_v$  the valence band edge and quasi fermi level for electrons and holes  $E_{fn}$  and  $E_{fp}$  which reduced to a single fermi level  $E_f$  in equilibrium.

We also show the intrinsic fermi level or fermi level in intrinsic semiconductor in the energy band diagram because it proves as a convenient reference, we also showed how the energies associated with impurities can be shown on the energy band diagram basically impurities and defects add additional energy levels within the energy gap the  $E_0$  reflects the electrostatic potential.

And the slope of  $E_0$  reflects the electrostatic field experienced by an electron touching the semiconductor surface from outside, the fermi level the quasi fermi level for electrons reflects the electron concentration as well as the electron current density, the quasi fermi level for holes reflects the hole concentration and current density for holes, this is how the levels  $E_0$ ,  $E_{fn}$  and  $E_{fp}$  together represent  $\psi$ ,  $n$  and  $p$  or  $E$ ,  $J_n$  and  $J_p$  of the drift diffusion model,  $E_c$  and  $E_v$  represents the variation of the E-K relation throughout the semiconductor volume.

Finally, we have mentioned that the electric field in the expression for current density for electrons should be expressed in terms of the gradient of the conduction band edge because conduction band edge which is the potential of electrons in the conduction band and the electric field term in the expression for hole current density should be expressed in terms of gradient of the valence band edge because valence band edge is the potential energy of holes.