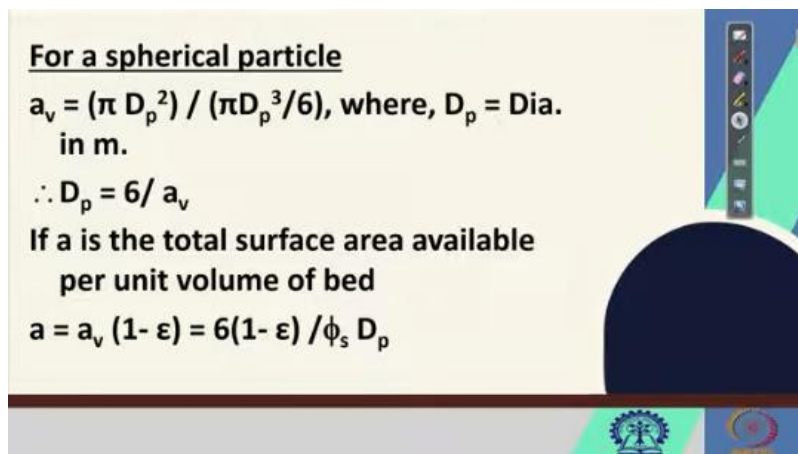


# IMPACT OF FLOW OF FLUIDS IN FOOD PROCESSING AND PRESERVATION

## Lecture37

### LECTURE 37 : HYDRAULIC RADIUS

Good afternoon, my dear boys and girls and students and friends. flow through packed beds right. We have defined some of the parameters according to its definition. It may be required that some more we need to define. For example, we had said for spherical particles



A, these we have defined, but we also defined the sphericity and the sphericity is defined as like this. This is true that the particles will always be spherical, hence let us define a term sphericity which will tell. how the particle is close to the sphere or in other words

sphericity will tell the spherical nature of the non spherical particle. it can be defined as  $\phi_s$  is  $S_s$  over  $V_s$  over or divided by  $S_p$  over  $V_p$   $S_s$  over  $V_s$  divided by  $S_p$  over  $V_p$  where obviously,  $S_s$  is the surface area. of a sphere having a diameter of  $D_p$  in meter square,  $V_s$  is the volume of this same sphere in meter cube,  $S_p$  is the surface area of a particle and  $V_p$  is the volume of the particle same particle in meter cube right.

- Now, it is always true that particles will always be spherical. Hence, let us define a term sphericity, which will tell how the particle is close to the sphere. Or, in other words, sphericity will tell the spherical nature of the non spherical particle. It can be defined as:

$$\phi_s = \frac{\frac{S_s}{V_s}}{\frac{S_p}{V_p}}$$

- Where,
- $S_s$ = surface area of a sphere having diameter  $D_p$ ,  $m^2$ ;  $V_s$ = volume of the same sphere,  $m^3$ ;  $S_p$ = surface area of a particle;  $V_p$ = volume of the same particle,  $m^3$ .

So, if this is the definition then we can say  $D_p$  for large particles is the diameter of a sphere having equal volume as the particle. Large particles means particles whose diameter can be measured or is greater than 1 millimeter you see 1 millimeter is very low right. and more than a millimeter that could be as large as small you see as much. I gave you example of grinding.

$D_p$  for large particles is the diameter of a sphere having equal volume as the particle. Large particles means particles whose diameter can be measured or is greater than 1 mm.

For fine particles,  $D_p$  is the sieve analysis diameter or nominal diameter.

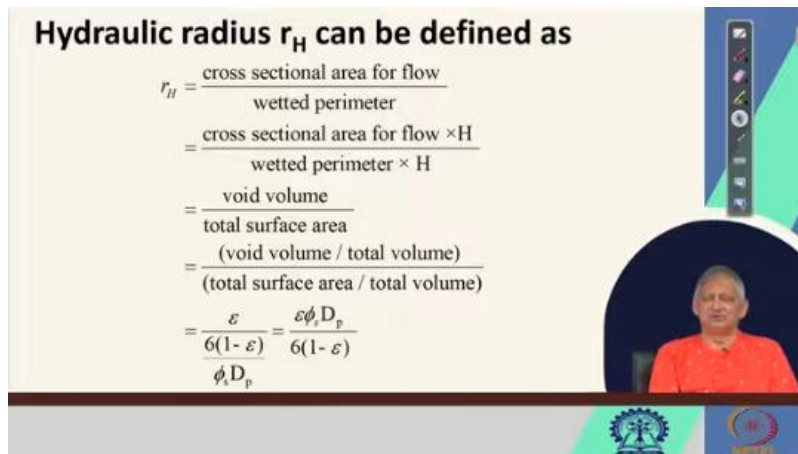
Hence,

$$\phi_s = \frac{6}{D_p} \frac{S_p}{V_p} = \frac{6V_p}{S_p D_p}$$

So, ground particles are not in millimeter they are in micrometer right. For the fine particles  $D_p$  is the sieve analysis diameter. or nominal diameter. Hence, we can write  $\phi_s$  is equal to 6 over  $D_p$  over  $S_p$  over  $V_p$  or that can be equated as 6  $V_p$  over  $S_p$  over  $D_p$  right. Now, we also

Define another term called hydraulic radius. Though you have done it in the previous class, but since these are very new, so a little repeat may not be bad for the new people who are not accustomed with this. That  $R_h$  is the hydraulic radius. is defined as cross sectional area for flow over wetted perimeter that is cross sectional times if we multiply both numerator and denominator with H. So, cross sectional area of our flow times H over wetted perimeter times H, this becomes void volume over total surface

**Hydraulic radius  $r_H$  can be defined as**


$$\begin{aligned}
 r_H &= \frac{\text{cross sectional area for flow}}{\text{wetted perimeter}} \\
 &= \frac{\text{cross sectional area for flow} \times H}{\text{wetted perimeter} \times H} \\
 &= \frac{\text{void volume}}{\text{total surface area}} \\
 &= \frac{(\text{void volume} / \text{total volume})}{(\text{total surface area} / \text{total volume})} \\
 &= \frac{\varepsilon}{\phi_s D_p} = \frac{\varepsilon \phi_s D_p}{6(1 - \varepsilon)}
 \end{aligned}$$


Now, void, if we now introduce another thing that if we divide both numerator and denominator with total volume, then it becomes void volume by total volume divided by total surface area by total volume. Now void volume by total volume by definition we have already given that is nothing but epsilon and this total surface area by total volume just now we have shown it to be for spherical particles it is 6 by  $\phi_s D_p$  into 1 minus epsilon. On rearrangement it can be said that epsilon  $\phi_s D_p$  over 6 into 1 minus epsilon right. So, hydraulic radius is expressed in terms

void volume, particle size and sphericity as epsilon  $\phi_s D_p$  by 6 into 1 minus epsilon. This may be required subsequently. Now, another also important that is called we have said hydraulic radius. So, equivalent of that is equivalent diameter. So, equivalent diameter is capital D that is nothing, but 4 times hydraulic radius that is  $4 R_h$ .

**Equivalent Diameter  $D = 4r_H = 4\epsilon \phi_s D_p / 6(1 - \epsilon)$**   
**Now if  $v'$  is the velocity based on empty cross section of the bed**  
**and  $v$  is the actual velocity through void space.**  
 $v' = \epsilon v$   
**and**

$$N_{Re} = \frac{Dv\rho}{\mu} = \frac{v'(4r_H)\rho}{\epsilon\mu} = \frac{4v'\epsilon\phi_s D_p \rho}{6\mu\epsilon(1-\epsilon)}$$

$$= \frac{4\phi_s D_p v' \rho}{6\mu(1-\epsilon)} \approx \frac{\phi_s D_p v' \rho}{\mu(1-\epsilon)}$$


So, if we substitute  $R_h$  with whatever just in the previous slide we have said that  $4\epsilon \phi_s (\Phi_s) D_p$  over  $6$  into  $1 - \epsilon$ . Now, if  $v'$  is the velocity based on empty cross section which I have shown in the previous class with a water bottle that if it is empty, then the velocity if there will be any flow of air or any fluid, that is  $v'$ . So,  $v'$  is the velocity based on empty cross section of the bed and  $v$  is the actual velocity through void space right. So, therefore, we can write  $v'$  is equal to  $\epsilon v$  right.

So, another term that is the non-dimensional,  $N_{Re}$ . So, on non-dimensional  $N_{Re}$  or Reynolds number we can define as  $D V \rho$  by  $\mu$  that is by definition. So, if we substitute appropriately then we can write  $v$  as  $v'$  and  $D$  as  $4 R_h$   $\rho$  remains same and  $v$  is  $v'$  by  $\epsilon$  right. So,  $\mu$  is there that means,  $Dv \rho$  by  $\mu$  is  $v'$  by  $\epsilon$ .

$4 R_h$  into  $\rho$  by  $\mu$ . So, this on rearrangement we can write  $4 v' \epsilon \phi_s D_p \rho$  by  $6 \mu$  into  $\epsilon(1 - \epsilon)$ . On rearrangement, this you can write to be  $4 \phi_s$  into  $D_p$  into  $v' \rho$  over  $6 \mu$  into  $1 - \epsilon$ , that is equal to  $\phi_s D_p$  into  $v' \rho$  over  $\mu$  into  $1 - \epsilon$  right. So, this is roughly equals to because  $4$  by  $6$  is what is almost  $0.7$  right. So, that means, it is close to  $1$ , that is why,

we can say it to be roughly equal to, by neglecting  $4$  by  $6$ , ok. That is  $\phi_s D_p v' \rho$  over  $\mu$  into  $1 - \epsilon$ . This you keep in mind that this is the Reynolds number, that is  $N_{Re}$  for packed bed right. So, there  $\phi_s$  is the sphericity,  $D_p$  is the diameter of the particle,  $v'$  is the velocity through empty cross section,  $\rho$  is the density of the fluid flowing,  $\mu$  is the viscosity of the fluid flowing and  $\epsilon$  is the

void fraction, right. So, this  $N_{Re}$  definition we have taken. Now, if we start from Hagen-Poiseuille equation for laminar flow, what it was? It was,  $\Delta P$  is equal to  $32 \mu v' L$  by  $D$  square, right,  $\Delta P$  is equal to  $32 \mu v' \Delta L$  by  $\epsilon$  into  $4 R_h$  whole square.

**Equivalent Diameter  $D = 4r_H = 4\epsilon \phi_s D_p / 6(1 - \epsilon)$**   
 Now if  $v'$  is the velocity based on empty cross section of the bed  
 and  $v$  is the actual velocity through void space.  
 $v' = \epsilon v$   
 and  

$$N_{Re} = \frac{D v \rho}{\mu} = \frac{v' (4r_H) \rho}{\epsilon \mu} = \frac{4v' \epsilon \phi_s D_p \rho}{6\mu \epsilon (1 - \epsilon)}$$

$$= \frac{4\phi_s D_p v' \rho}{6\mu (1 - \epsilon)} \approx \frac{\phi_s D_p v' \rho}{\mu (1 - \epsilon)}$$

**From Hagen Poiseuille Eq. for laminar flow**  

$$\Delta p = \frac{32\mu v' \Delta L}{D^2} = \frac{32\mu v' \Delta L}{\epsilon (4r_H)^2} = \frac{32\mu v' \Delta L}{\left(4 \frac{\epsilon \phi_s D_p}{6(1 - \epsilon)}\right)^2}$$

$$= \frac{72\mu v' \Delta L (1 - \epsilon)^2}{\epsilon^3 \phi_s^2 D_p^2}$$

$$\approx \frac{150\mu v' \Delta L (1 - \epsilon)^2}{\epsilon^3 \phi_s^2 D_p^2}$$
**This is called Blake – Kozeny equation and is valid for  $N_{Re} < 10$ .**

So, this is equal to  $32 \mu v' \Delta L$  over  $4 \epsilon \phi_s D_p$  over  $6$  into  $1$  minus  $\epsilon$  whole square right. This can be written as  $72 \mu v' \Delta L$  into  $1$  minus  $\epsilon$  whole square over  $\epsilon$  cube  $\phi_s$  square  $D_p$  square. That can roughly be said, instead of  $72$ , it has been found by the scientist that the prediction of  $\Delta P$  is more closer to the actual if instead of  $72$  it is taken  $150$ . That is why it was changed to  $150 \mu v'$  into  $\Delta L$


**From Hagen Poiseuille Eq. for laminar flow**

$$\Delta p = \frac{32\mu v \Delta L}{D^2} = \frac{32\mu v' \Delta L}{\varepsilon (4r_h)^2} = \frac{32\mu v' \Delta L}{\left(4 \frac{\varepsilon \phi_s D_p}{6(1-\varepsilon)}\right)^2}$$

$$= \frac{72\mu v' \Delta L (1-\varepsilon)^2}{\varepsilon^3 \phi_s^2 D_p^2}$$

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**This is called Blake – Kozeny equation and is valid for  $N_{Re} < 10$ .**



L into 1 minus epsilon whole square by epsilon cube phi<sub>s</sub> square D<sub>p</sub> square right. So, we have done very quickly, but we can repeat because a new thing that is the application of the Hagen-Poiseuille equation. So, we start with Hagen-Poiseuille equation. It says that delta P is equal to 32 mu v delta L by D square.

That was the application of Hagen-Poiseuille's equation. We come for laminar flow as delta P is equal to 32 mu v delta L over D square. This was for Hagen Poiseuille for laminar flow right in pipe flow. So, this is equal to 32 mu instead of v we write v prime delta L, the moment we are writing it is v prime, obviously, epsilon is also coming.

So, v prime by epsilon where we are replacing D with R<sub>h</sub> hydraulic radius and that hydraulic radius is 4 R<sub>h</sub> whole square right. So, 32 mu v prime delta L that hydraulic radius again we are substituting with the value of hydraulic radius, that is, 4 times epsilon, it was already there and phi<sub>s</sub> D<sub>p</sub> by 6 into 1 minus epsilon whole square. So, rearranging it we can write it to be 72 mu v prime delta L into 1 minus epsilon whole square divided by epsilon



cube into  $\phi_s$  square  $D_p$  square right,  $\phi_s$  square  $D_p$  square yes. some scientists found that instead of this 72, if 150 is taken, then the prediction of  $\Delta P$  is more closer to the actual.

So, that is why they substituted it with 150 So,  $\Delta p$  is  $150 \mu v' \Delta L$  into  $1$  minus  $\epsilon$  whole square by  $\epsilon$  cube  $\phi_s$  square  $D_p$  square. This was as I said by some scientist did it. So, this was done by Blake Kozini. and that in their name the equation is named as Blake Kozeny equation right, Blake Kozeny, K O Z E N Y and this is applicable for or valid for  $N_{Re}$  less than 10, right mind it that Blake Kozeny equation

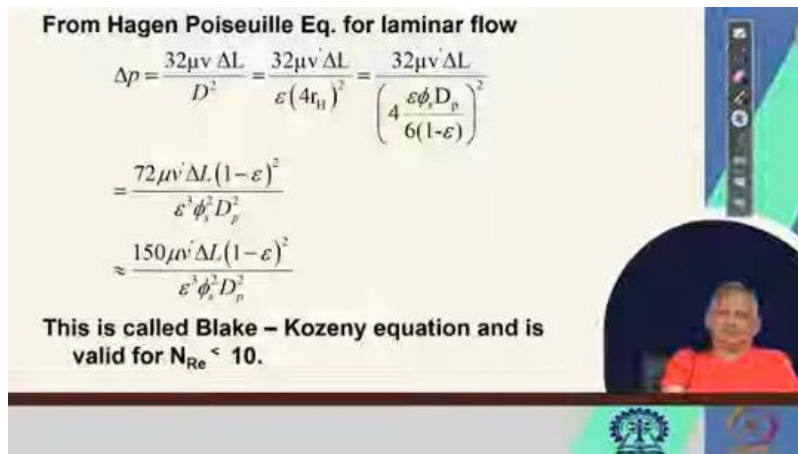
**From Hagen Poiseuille Eq. for laminar flow**

$$\Delta p = \frac{32\mu v' \Delta L}{D^3} = \frac{32\mu v' \Delta L}{\epsilon (4r_H)^3} = \frac{32\mu v' \Delta L}{\left(4 \frac{\epsilon \phi_s D_p}{6(1-\epsilon)}\right)^3}$$

$$= \frac{72\mu v' \Delta L (1-\epsilon)^3}{\epsilon^3 \phi_s^3 D_p^3}$$

$$\approx \frac{150\mu v' \Delta L (1-\epsilon)^3}{\epsilon^3 \phi_s^3 D_p^3}$$

**This is called Blake – Kozeny equation and is valid for  $N_{Re} < 10$ .**



derived from the Hagen Poiseuille's equation for laminar pipe flow is valid for  $N_{Re}$  less than 10 right. Now, the other one for turbulent flow we can use the equation  $\Delta P_f$  is  $4 f$  rho  $\Delta L$  by  $D$  into  $v$  square by 2. Now, if you substitute,  $\Delta P_f$  is  $4 f$  rho into  $\Delta L$ . So,  $D$ , we can substitute with  $4 R_h$ , and in turn it is  $4$  by  $6 \epsilon \phi_s D_p$  by  $1$  minus  $\epsilon$  into, that  $v$  as  $v'$  square by  $2 \epsilon$  square,

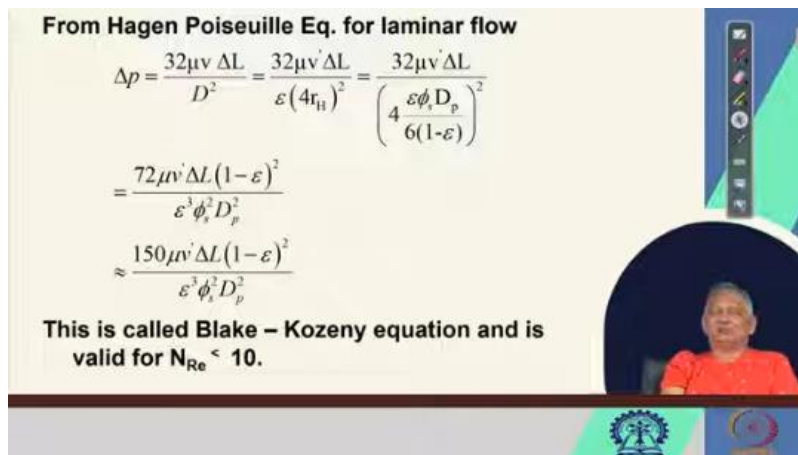
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$$= \frac{72\mu v' \Delta L (1-\epsilon)^3}{\epsilon^3 \phi_s^3 D_p^3}$$

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**This is called Blake – Kozeny equation and is valid for  $N_{Re} < 10$ .**



**For turbulent flow we use the following eq.:-**

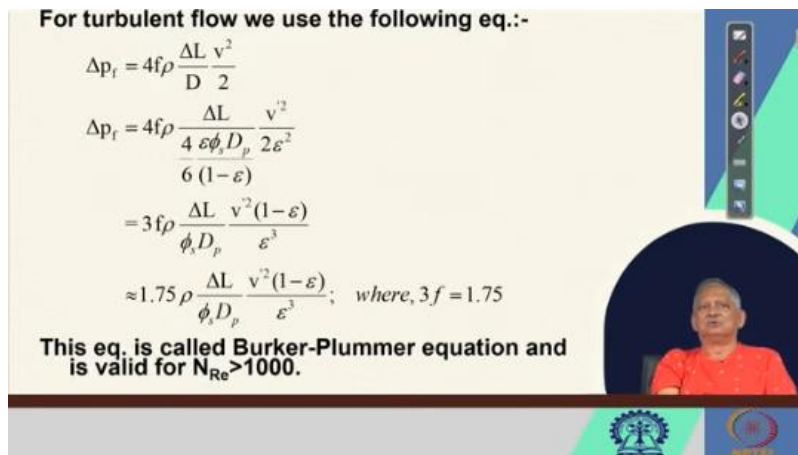
$$\Delta p_f = 4f\rho \frac{\Delta L}{D} \frac{v^2}{2}$$

$$\Delta p_f = 4f\rho \frac{\Delta L}{4 \varepsilon \phi_s D_p} \frac{v'^2}{2 \varepsilon^2}$$

$$= 3f\rho \frac{\Delta L}{\phi_s D_p} \frac{v'^2 (1-\varepsilon)}{\varepsilon^3}$$

$$\approx 1.75 \rho \frac{\Delta L}{\phi_s D_p} \frac{v'^2 (1-\varepsilon)}{\varepsilon^3}; \text{ where, } 3f = 1.75$$

**This eq. is called **Burker-Plummer** equation and is valid for  $N_{Re} > 1000$ .**



right, because this  $v'$  square by  $\varepsilon$  square and there was a 2. So, these on rearrangement we can write  $3f\rho \Delta L$  by  $\phi_s D_p$  into  $v'$  square into  $1 - \varepsilon$  by  $\varepsilon$  cube right. That again, some of the scientist have found that, instead of  $3f\rho$ , if we take, instead of  $3f$ , 1.75, then the prediction is more closer to the actual. So, if this be then they have said that it is 1.75  $\Delta p_f$ , is 1.75 into  $\rho$  into  $\Delta L$  by  $\phi_s D_p$  into  $v'$  square into  $1 - \varepsilon$  by  $\varepsilon$  cube. Obviously, in this case  $3f$  is nothing but equal to 1.75.

This was done by Barker Plummer, B-U-R-K-E-R Plummer, P-L-U-M-M-E-R, these two scientists and according to their name this equation is named as Barker Plummer, right? This equation and this is valid, which is very, very important that this is also valid for  $N_{Re}$  greater than 1000. that means, from the Hagen Poiseuille and the from the turbulent, for both laminar and turbulent flow. We have studied from the very beginning that is Hagen Poiseuille's equation and the friction factor equation and then they came to two equations, two relations of  $\Delta p$ .

**For turbulent flow we use the following eq.:-**

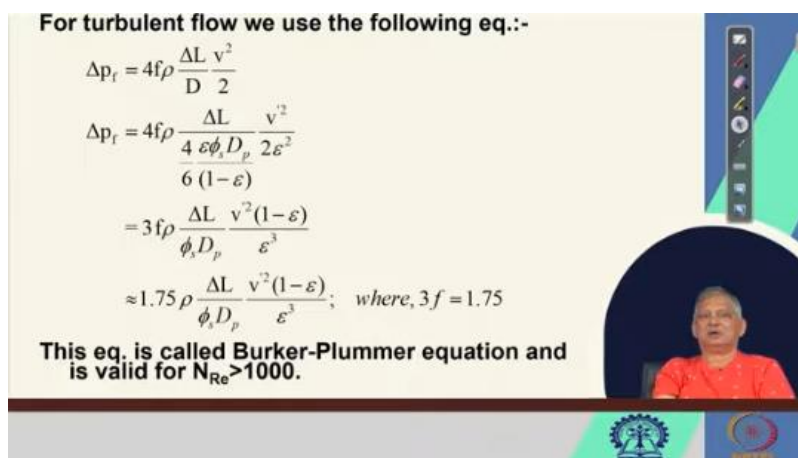
$$\Delta p_f = 4f\rho \frac{\Delta L}{D} \frac{v^2}{2}$$

$$\Delta p_f = 4f\rho \frac{\Delta L}{4 \varepsilon \phi_s D_p} \frac{v'^2}{2 \varepsilon^2}$$

$$= 3f\rho \frac{\Delta L}{\phi_s D_p} \frac{v'^2 (1-\varepsilon)}{\varepsilon^3}$$

$$\approx 1.75 \rho \frac{\Delta L}{\phi_s D_p} \frac{v'^2 (1-\varepsilon)}{\varepsilon^3}; \text{ where, } 3f = 1.75$$

**This eq. is called **Burker-Plummer** equation and is valid for  $N_{Re} > 1000$ .**

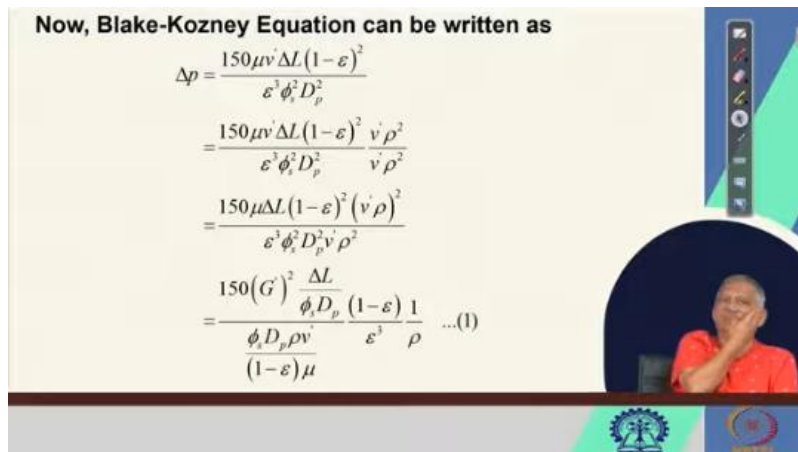




that is finding out the pressure drop theoretically, but both has the limitation that for the Hagen Poiseuille equation that is Blake Kozeny equation. It is valid up to  $N_{Re}$  less than 10 and for Barker Plummer for the turbulent flow. it has been shown that the relation delta P is valid for  $N_{Re}$  greater than 1000. Then the question comes obviously that what about, between 10 to 1000.

Is there anything? Yes. Subsequently, some scientists came forward and they started by taking both the equations, right. So, from the black cousin equation, they had taken, and I do not say, it is renovated obviously, they did something with that and the delta P that became equal to  $150 \mu v' \Delta L (1 - \epsilon)^2$  by  $\epsilon^3 \phi_s^2 D_p^2$  square right. So, this is equals to  $150 \mu v' \Delta L (1 - \epsilon)^2$  by  $\epsilon^3 \phi_s^2 D_p^2$  square into  $v'$  prime into  $\rho$  square by  $v'$  prime into  $\rho$  square. That means, they have introduced both in the numerator and denominator a new term or not new term some parameters like  $v'$  prime  $\rho$  square both in numerator and denominator, right and then reshuffled the equation.

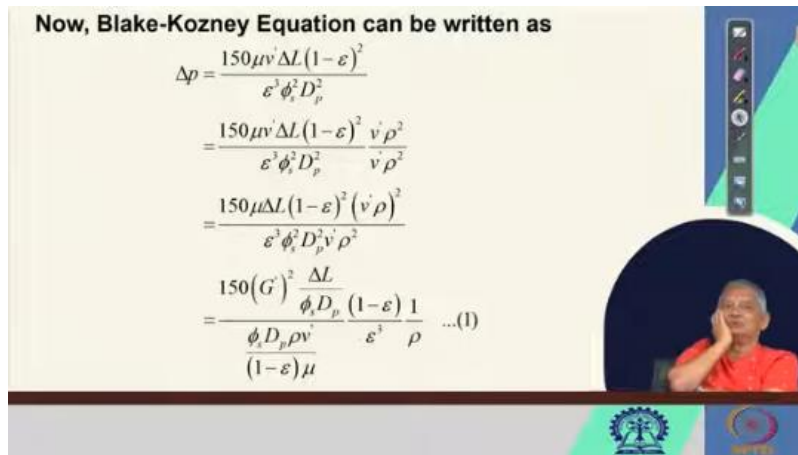
**Now, Blake-Kozeny Equation can be written as**

$$\begin{aligned} \Delta p &= \frac{150 \mu v' \Delta L (1 - \epsilon)^2}{\epsilon^3 \phi_s^2 D_p^2} \\ &= \frac{150 \mu v' \Delta L (1 - \epsilon)^2}{\epsilon^3 \phi_s^2 D_p^2} \frac{v' \rho^2}{v' \rho^2} \\ &= \frac{150 \mu \Delta L (1 - \epsilon)^2 (v' \rho)^2}{\epsilon^3 \phi_s^2 D_p^2 v' \rho^2} \\ &= \frac{150 (G')^2 \Delta L}{\phi_s D_p \rho v' (1 - \epsilon)^2} \frac{1}{\epsilon^3 \rho} \dots (1) \end{aligned}$$


as  $150 \mu \Delta L (1 - \epsilon)^2$  into  $v' \rho$  whole square because one  $v'$  prime was here inside. So, that they have taken inside and it has become  $v' \rho$  whole square. and the denominator has become  $\epsilon^3 \phi_s^2 D_p^2$  square then  $v'$  prime into  $\rho$  square right  $\epsilon^3 \phi_s^2 D_p^2$  square  $v'$  prime into  $\rho$  square. Now, they rearranged this equation as  $150$  instead of  $\mu$  and  $v'$  instead of  $\rho$   $v'$ , right they have taken as  $\rho v'$  equal to  $G'$  prime, right,  $150 G'$  prime whole square into  $\Delta L$  by  $\phi_s D_p$  into  $(1 - \epsilon)^2$  right divided by  $\phi_s D_p \rho v'$  into  $(1 - \epsilon)^2$  into  $\mu$  by  $\epsilon^3$  and divided by  $\rho$ . So, this equation some people have given a numbering of the equation as equation 1. I repeat from the Blake-Kozeny equation delta p is equal to  $150 \mu$  by  $\mu v'$  prime into  $\Delta L$  into  $(1 - \epsilon)^2$  by  $\epsilon^3$

cube  $\phi_s$  square  $D_p$  square by introducing both numerator and denominator with  $v$  prime  $\rho$  square. They have introduced and rearranged to the level that  $\Delta p$  becomes equal to

**Now, Blake-Kozeny Equation can be written as**

$$\begin{aligned}\Delta p &= \frac{150 \mu v' \Delta L (1-\varepsilon)^2}{\varepsilon^3 \phi_s^2 D_p^2} \\ &= \frac{150 \mu v' \Delta L (1-\varepsilon)^2}{\varepsilon^3 \phi_s^2 D_p^2} \frac{v' \rho^2}{v' \rho^2} \\ &= \frac{150 \mu \Delta L (1-\varepsilon)^2 (v' \rho)^2}{\varepsilon^3 \phi_s^2 D_p^2 v' \rho^2} \\ &= \frac{150 (G')^2 \Delta L}{\phi_s D_p \rho v' \varepsilon^3} \frac{(1-\varepsilon)}{\rho} \dots (1)\end{aligned}$$


150  $G'$  prime which is nothing by it, but  $\rho$   $\rho$   $v$  prime right into  $\Delta L$  by  $\phi_s$   $D_p$  into 1 minus  $\varepsilon$  divided by  $\phi_s$   $D_p$   $\rho$   $v$  prime over 1 minus  $\varepsilon$  into  $\mu$  by  $\varepsilon$  cube and also divided by  $\rho$  right. I should say the other way around that 150  $G'$  prime  $\Delta L$  by  $\phi_s$   $D_p$  divided by  $\phi_s$   $D_p$  into  $\rho$   $v$  prime over 1 minus  $\varepsilon$  into  $\mu$  times 1 minus  $\varepsilon$  by  $\phi_s$   $\varepsilon$  cube. times 1 by  $\rho$  is the right way of reading it, right. Then from the Barker Plummer equation, they have done that  $\Delta P$  is equal to 1.75

$\rho$  into  $\Delta L$  by  $\phi_s$   $D_p$  into  $v$  prime ( $v'$ ) square into 1 minus  $\varepsilon$  by  $\varepsilon$  cube is equal to 1.75  $\rho$   $\Delta L$  by  $\phi_s$   $D_p$  times  $v$  prime square into 1 minus  $\varepsilon$  by  $\varepsilon$  cube into  $\rho$  by  $\rho$ . they multiplied  $\rho$  both in the numerator and denominator. Now, they rearranged in the form like 1.75  $V$   $\rho$  rather  $V$  prime  $\rho$  whole square into  $\Delta L$  by  $\phi_s$  into  $D_p$  into 1 minus  $\varepsilon$  by  $\varepsilon$  cube into 1 by  $\rho$ . And this is equals to 1.75  $G'$  prime square  $\Delta L$  by  $\phi_s$   $D_p$  into 1 minus  $\varepsilon$  by  $\varepsilon$  cube into 1 by  $\varepsilon$  right this on rearrangement.

That means, from the Berker Plummer equation they have multiplied both numerator and denominator with  $\rho$  and then rearranged to get the final of  $\Delta p$  as  $\Delta p$  is equal to 1.75  $G'$  prime whole square  $\Delta L$  by  $\phi_s$   $D_p$  into 1 minus  $\varepsilon$  by  $\varepsilon$  cube into 1 by  $\rho$  right. Then they have added these two equation both the outcome of the Blake-Kozeny as well as the Berker-Plummer equation.

They have added these two and found out that  $\Delta P$  is equal to 150  $G'$  prime ( $G'$ ) whole square into  $\Delta L$  by  $\phi_s$   $D_p$  into 1 minus  $\varepsilon$  by  $\varepsilon$  cube into 1 by  $\rho$  divided

by  $\phi_s D_p \rho v' (1 - \epsilon) \mu + 1.75 \mu G'^2 \Delta l$  by  $\phi_s D_p (1 - \epsilon)$  by  $\epsilon^3$  by  $1/\rho$ . These on rearrangement they have done  $\Delta p \rho$  by  $G'^2$  into  $\phi_s D_p \Delta L$  into  $\epsilon^3$  by  $1 - \epsilon$  that is equal to  $150 \text{ by } N_{Re} + 1.75$ . So, by adding some people have both by nurturing equation by nurturing Blake Kozini and Barker

plumber equation some people have made adding them and rearranging making a final expression like  $\Delta p$  is equal to  $\Delta p \rho$  by  $G'^2$  into  $\phi_s D_p \Delta L$  into  $\epsilon^3$  by  $1 - \epsilon$  that is equals to  $150 \text{ by } N_{Re} + 1.75$  and this equation is known as very famous equation called Ergun's equation. It is valid very much in pack bed right. Now, the time is up we will come and discuss in the next class.

Thank you very much.

Thank you.